

Chiral Perturbation for Large Momentum Effective Field Theory

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- 1 Large Momentum Effective Field Theory (LaMET)
- 2 Finite Volume Extrapolation
- 3 The framework to apply Chiral Perturbation Theory (ChPT)
- 4 Results
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Lattice QCD provides a tool to study nonperturbative properties in QCD, e.g. parton physics.

- partonic observables: nonlocal in light front direction, e.g. parton distribution function
- light cone condition only gives a trivial solution in Euclidean Space

$$x^2 = 0 \tag{1}$$

- LaMET: quasi observables introduced

- Factorization between finite momentum and infinite momentum.

$$\langle \tilde{\mathcal{O}} \rangle(\Lambda, P_z, \Lambda_{QCD}) = Z(\Lambda, P_z, \mu) \otimes \langle \mathcal{O} \rangle(\mu, \Lambda_{QCD}) + \mathcal{O}(1/P_z^2) \quad (2)$$

$$\Lambda_{QCD} \ll \mu \ll P_z \ll \Lambda \quad (3)$$

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 - light quark masses m_q
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- Extrapolation of m_q and L requires large amount of computing resources
 - ⇒ Model independence of extrapolation needed to improve efficiency of computing.
- Effective field theory (EFT): Chiral Perturbation Theory

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To find the framework to apply ChPT to LaMET

- ChPT: EFT under Λ_χ , while partonic processes go with a large momentum transfer.

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- Operator product expansion (OPE) of non-local quark bilinear operators in the basis of local operators.

$$\lambda_\mu \bar{\psi}(z) \Gamma^\mu W[z, 0] \psi(0) \approx \sum_{n=0}^{\infty} \frac{(iz)^n}{n!} \lambda_{\mu_1} \lambda_{\mu_2} \cdots \lambda_{\mu_n} \mathcal{O}^{\mu\mu_1 \cdots \mu_n} \quad (4)$$

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- Only leading twist operators considered \Rightarrow traces suppressed by higher power of momentum $\mathcal{O}(M^{2n}/P_z^{2n})$

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- Finite volume effect between local operator basis.

For example, $\Gamma^\mu = \gamma^\mu$ in isosinglet case

$$\begin{aligned} \mathcal{O}_{u-d}^{\mu\mu_1\mu_2\cdots\mu_n} &= c_1 \bar{N} V^{(\mu} v^{\mu_1} v^{\mu_2} \dots v^{\mu_n)} (u\tau^3 u^\dagger + u^\dagger \tau^3 u) N \\ &+ \tilde{c}_1 \bar{N} S^{(\mu} v^{\mu_1} v^{\mu_2} \dots v^{\mu_n)} (u\tau^3 u^\dagger - u^\dagger \tau^3 u) N + \dots \end{aligned}$$

- Matching quark level operators to hadronic level operators with symmetries preserved.
- The theory further separated by the chiral symmetry breaking scale Λ_χ

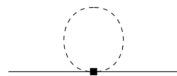
$$\mathcal{L} = \frac{F^2}{4} \text{tr}(\partial_\mu \Sigma \partial^\mu \Sigma^\dagger) + \eta \text{tr}(M \Sigma^\dagger + M^\dagger \Sigma) + \bar{N} i v \cdot D N + 2g_A \bar{N} S \cdot A N + \dots \quad (6)$$



(a)



(b)



(c)



(d)



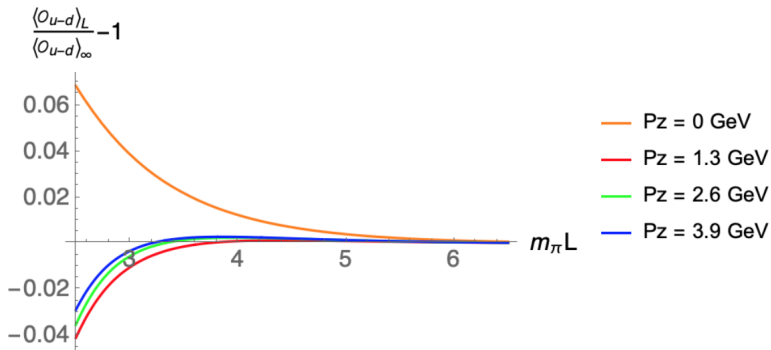
(e)



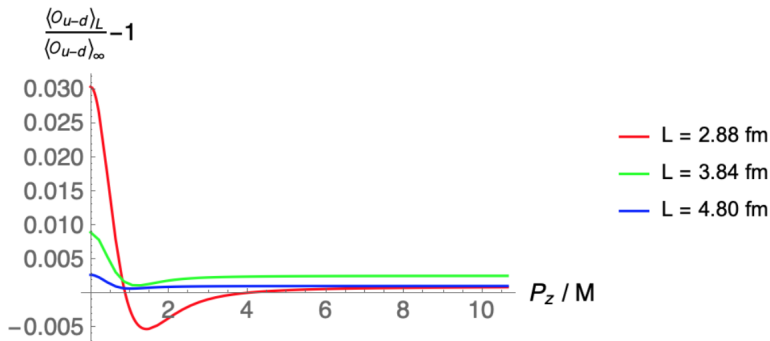
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$$\frac{\langle O_{u-d}^{\mu\mu_1\cdots\mu_n} \rangle_L}{\langle O_{u-d}^{\mu\mu_1\cdots\mu_n} \rangle_\infty} - 1 = -(1-\delta_{0n}) \frac{m_\pi^2}{4\pi^2 F_\pi^2} \sum_{\vec{n} \neq 0} \left[\frac{K_1(nm_\pi L)}{nm_\pi L} + 3g_A^2 J(nm_\pi L, \frac{\vec{n} \cdot \vec{P}}{nM}) \right] \quad (7)$$



(b)



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- Model independent formula for extrapolations
- Lorentz contraction of the nucleon size

Thanks for Your Listening