



# CALCULATION OF PION DISTRIBUTION AMPLITUDE I: *RENORMALIZATION AND POWER ACCURACY*

Rui Zhang

University of Maryland

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# Outline

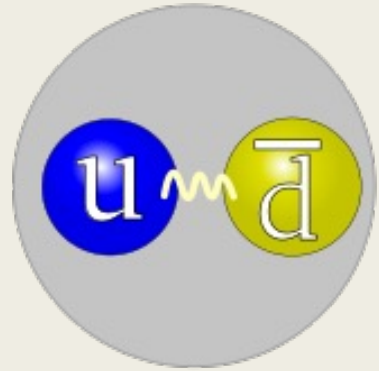
Introduction to Pion DA

Renormalization of DA

Power Accuracy in LaMET

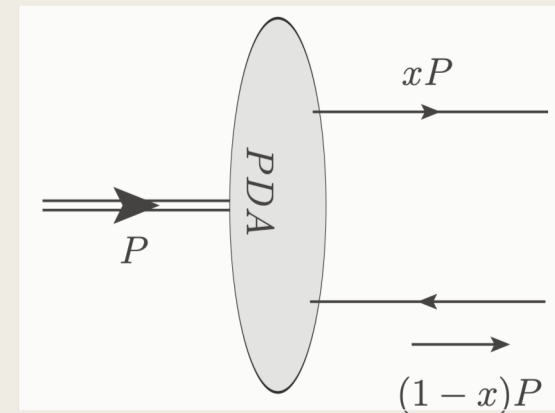
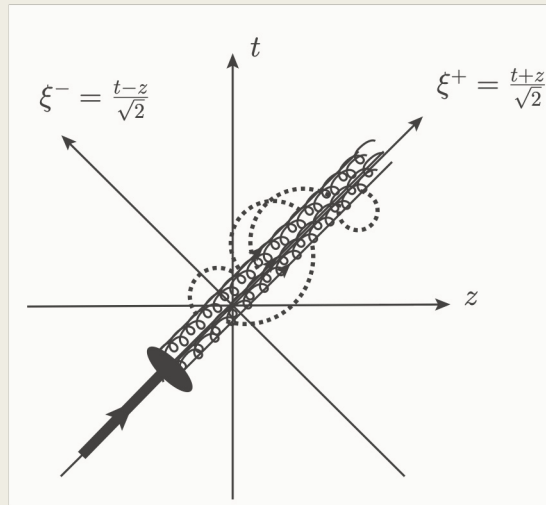
Application to DA Analysis

# Pion Distribution Amplitude (DA)



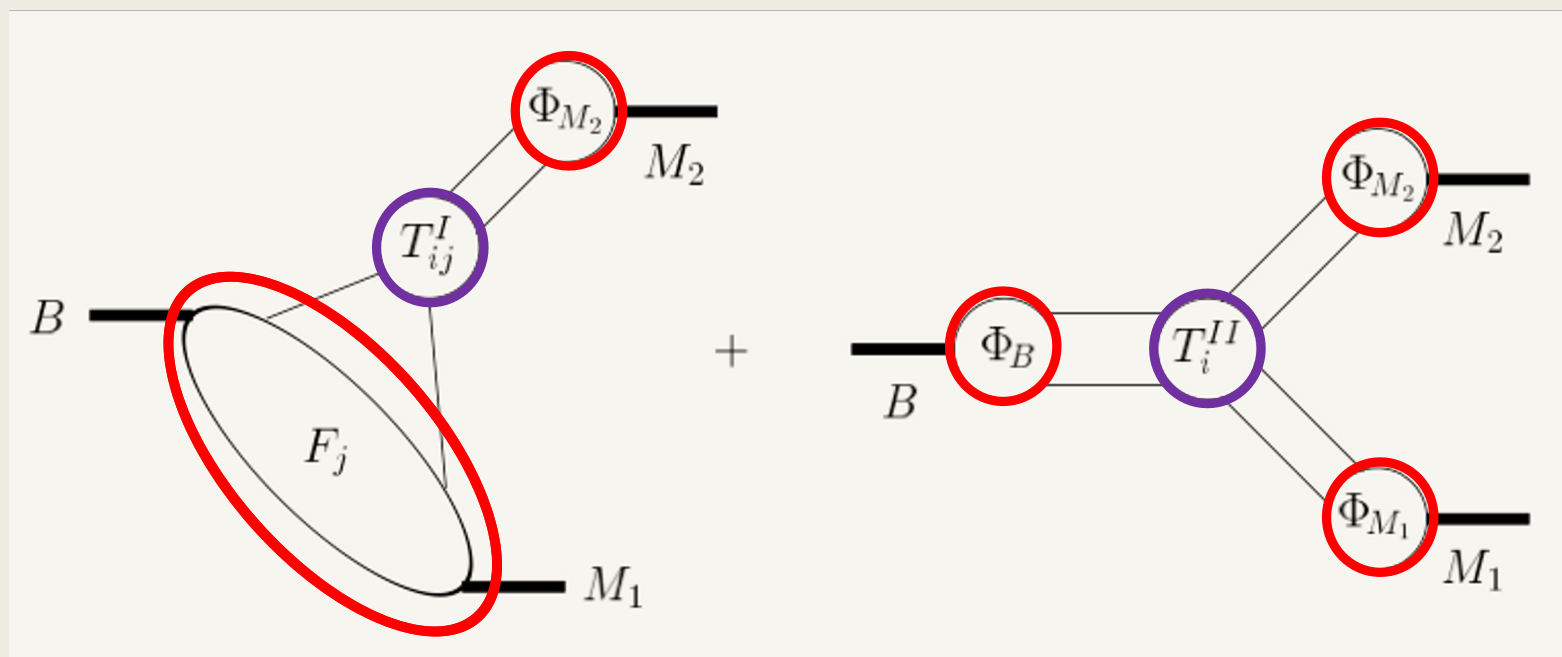
Pion lightfront DA  $\phi(x)$ : **probability amplitude** of pion in the bound state's minimal Fock component  $|q\bar{q}\rangle$


$$\phi(x, \mu) = \frac{1}{if_\pi} \int \frac{d\xi^-}{2\pi} e^{i(\frac{1}{2}-x)\xi^- p^+} \langle 0 | \bar{q} \left( \frac{\xi^-}{2} \right) \gamma^- \gamma_5 U \left( \frac{\xi^-}{2}, -\frac{\xi^-}{2} \right) q \left( -\frac{\xi^-}{2} \right) | \pi(p) \rangle$$




# Factorization of hard exclusive process

DA as important input to hard exclusive process at  $Q^2 \gg \Lambda_{\text{QCD}}^2$ : Beneke, et al. NPB(2001)



 : Nonperturbative, IR

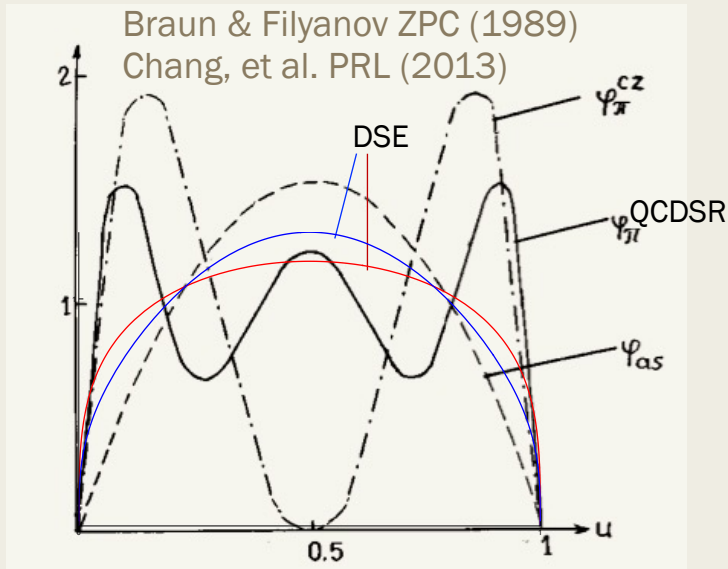
 : Perturbative, UV

$$\begin{aligned} \langle \pi K | Q_i | B \rangle = & F_0^{B \rightarrow \pi} T_{K,i}^I * f_K \Phi_K + F_0^{B \rightarrow K} T_{\pi,i}^I * f_\pi \Phi_\pi \\ & + T_i^{II} * f_B \Phi_B * f_K \Phi_K * f_\pi \Phi_\pi, \end{aligned}$$

# Previous DA Calculations

➤ Theoretical calculations

1. QCD sum rule Chernyak, et al., NPB (1982)
2. Dyson-Schwinger Equation



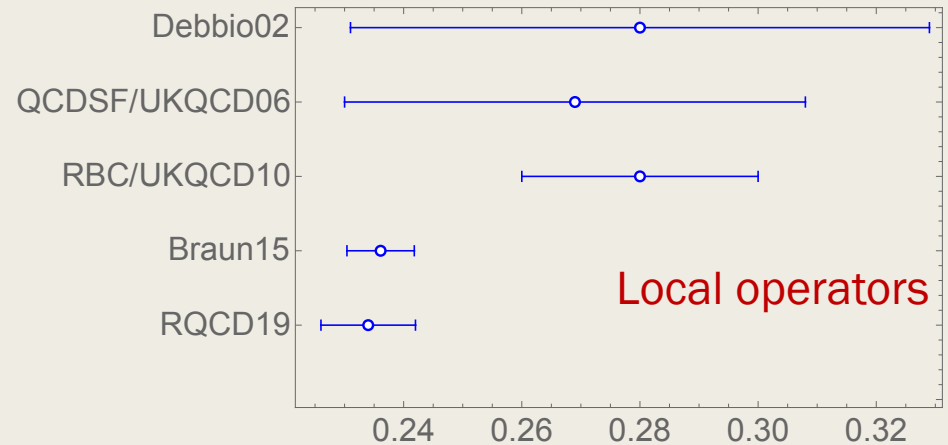
Model-dependent

➤ Moment calculations on Lattice

$$\langle \xi^n \rangle = \int_0^1 dx \phi(x) (2x - 1)^n$$

HOPE PRD (2022)  
RQCD PRD (2019)  
Braun, et al., EPJC (2007)

1. Current-current correlators
2. Local twist-2 operators



Pros: Precise lowest moment

Cons: Unable to extract x-dependence

# X-dependence calculation

- Large Momentum Effective Theory
- quasi-DA: **Same IR behavior/** different UV behavior

$$\tilde{\phi}(x, P_z) = \frac{1}{if_\pi} \int \frac{dz}{2\pi} e^{i(\frac{1}{2}-x)zP_z} \langle 0 | \bar{q}(z) \gamma_z \gamma_5 U(z, -z) q(-z) | \pi(P_z) \rangle$$

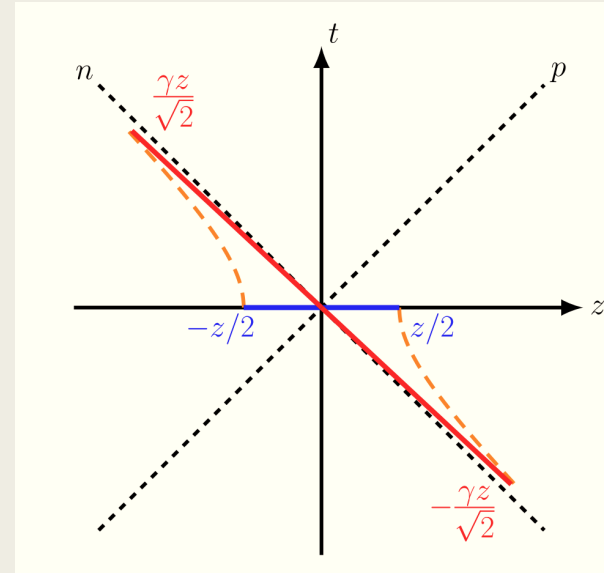
- Approach  $P \rightarrow \infty$  limit through large  $P_z$  expansion
- Matching to lightcone distribution

$$\tilde{\phi}(x, P_z) = \int_0^1 dy C(x, y, \mu, P_z) \phi(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{QCD}^2}{x^2 P_z^2}, \frac{\Lambda_{QCD}^2}{(1-x)^2 P_z^2}\right)$$

Pros: Direct x-dependence calculation, works well in mid-x region

Cons: Large  $P_z$  expansion breaks down near endpoints  
More complicated renormalization

Ji, et al., RMP (2021)



Xiong, et al., PRD (2014)

Ma, et al., PRD (2018)

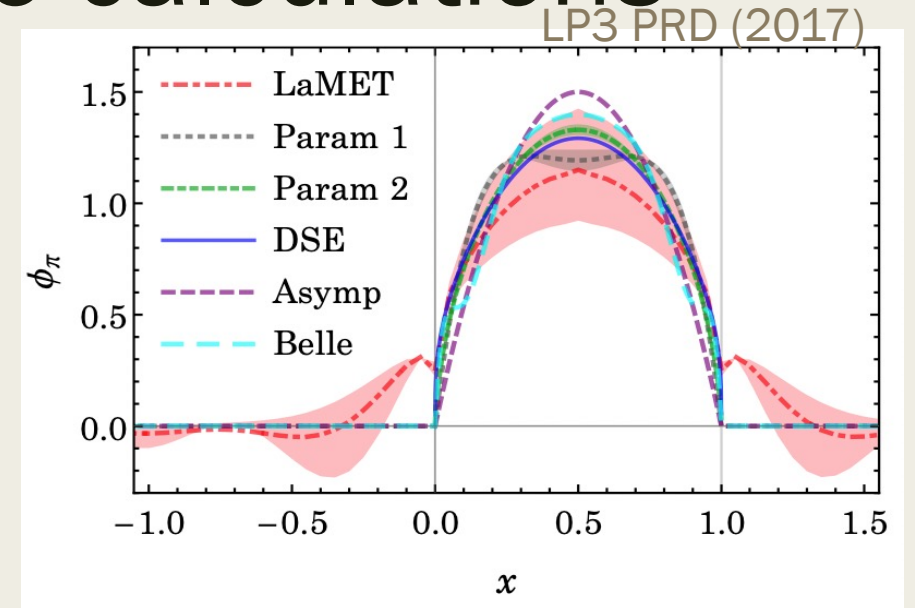
Izubuchi, et al., PRD (2018)

Liu, et al., PRD, (2019)

Ji, et al., RMP (2021)

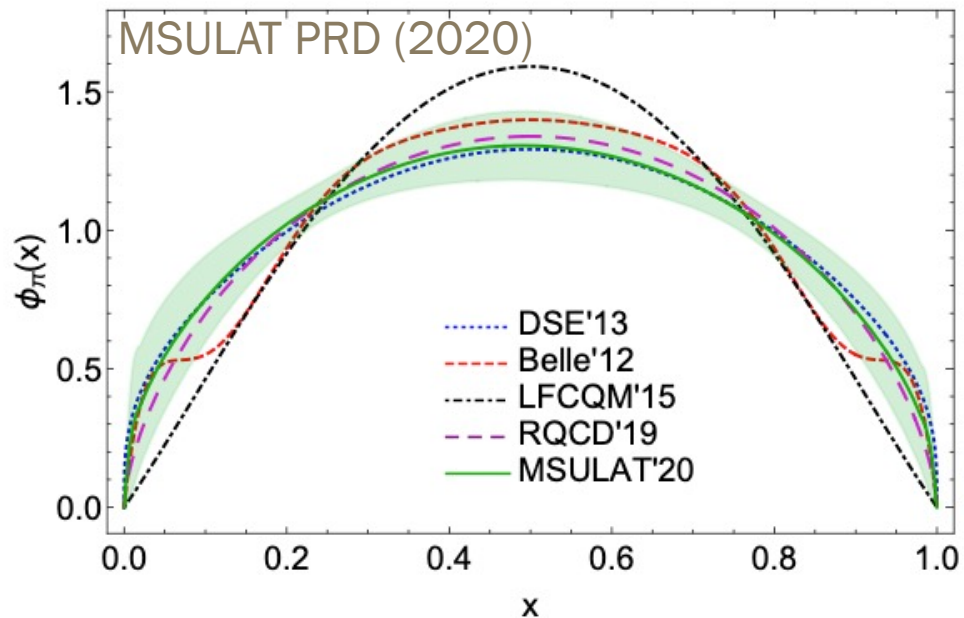
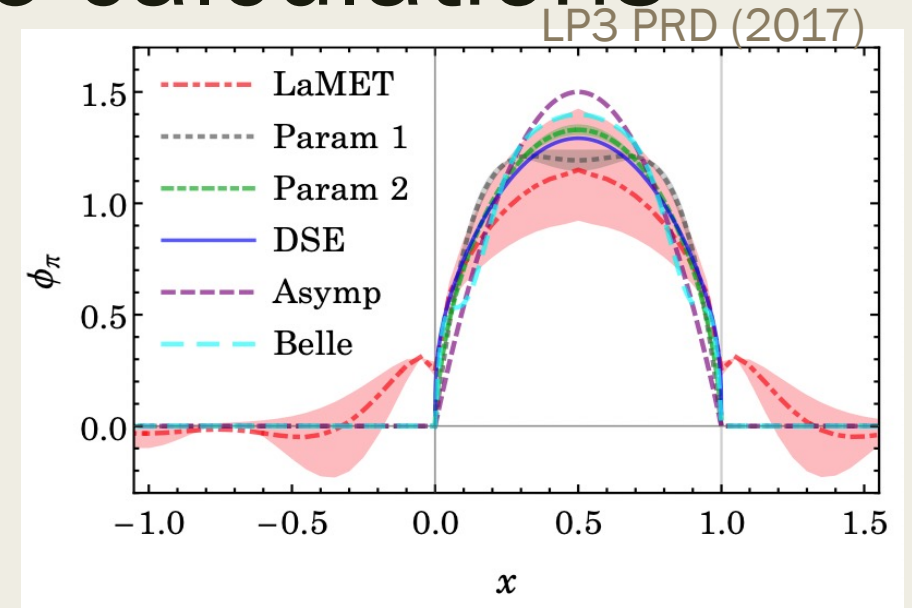
# Progress in $x$ -dependence calculations

- Lattice efforts
  - *Boost to large momentum*



# Progress in $x$ -dependence calculations

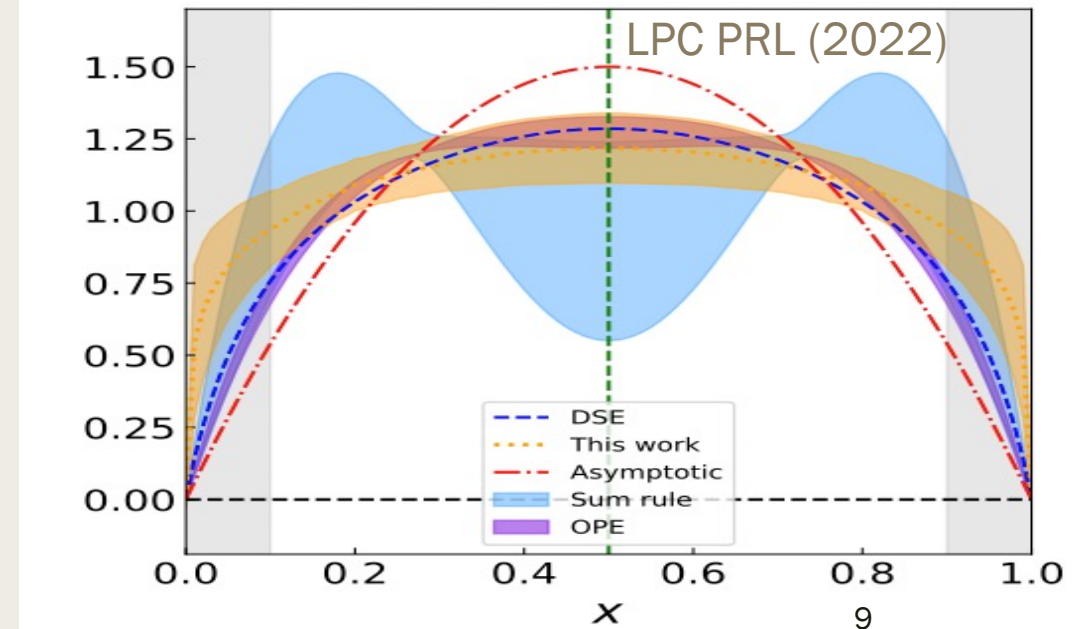
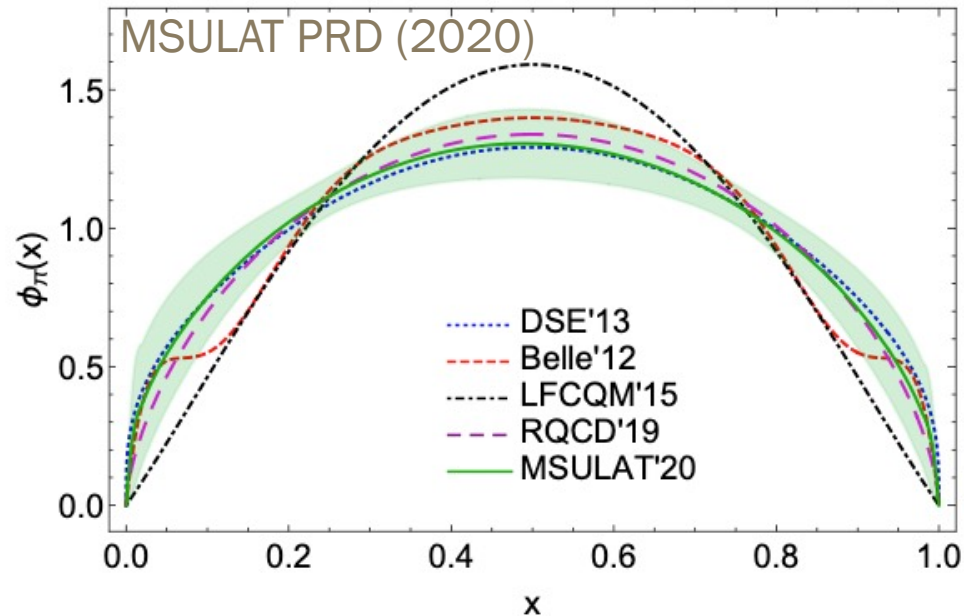
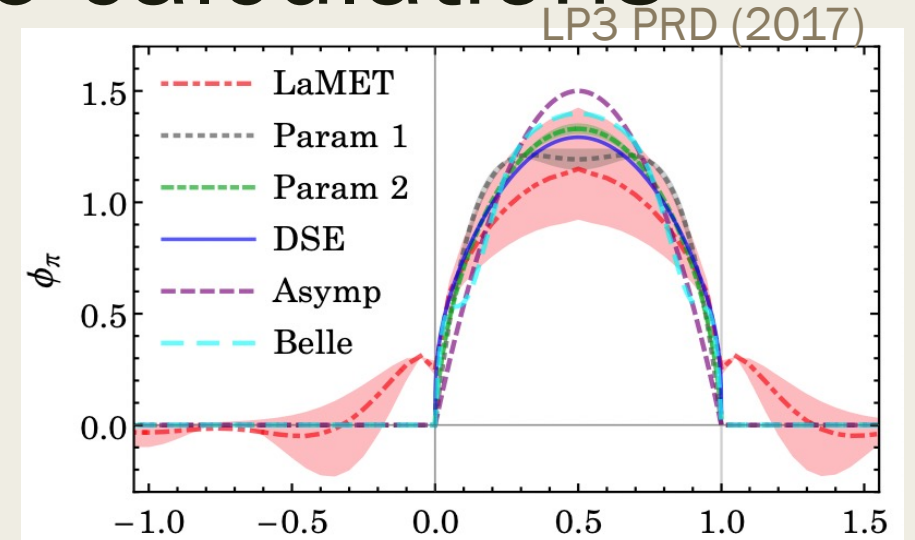
- Lattice efforts
  - *Boost to large momentum*
  - *Continuum extrapolation*





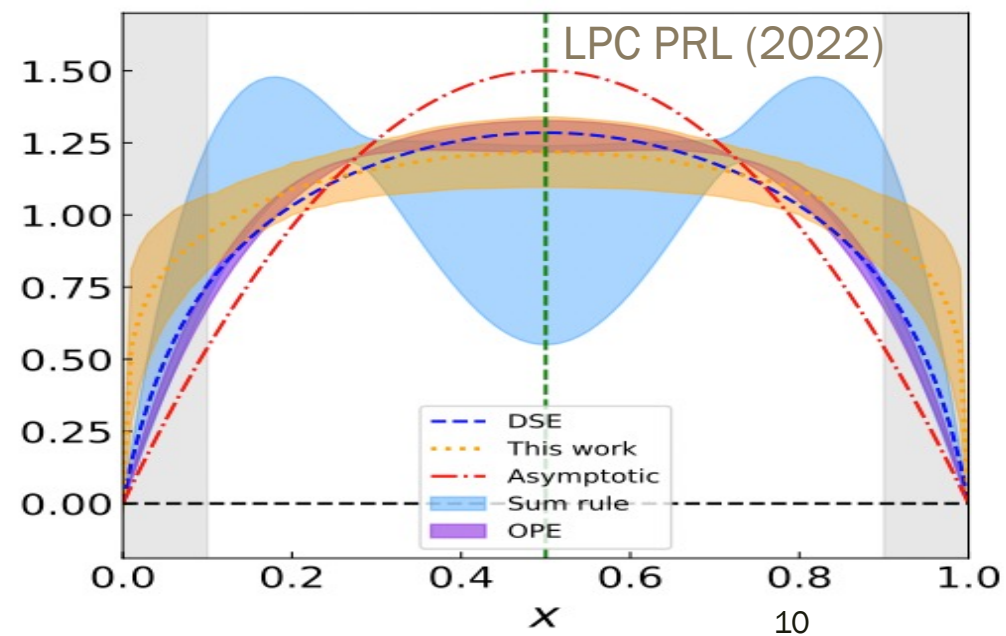
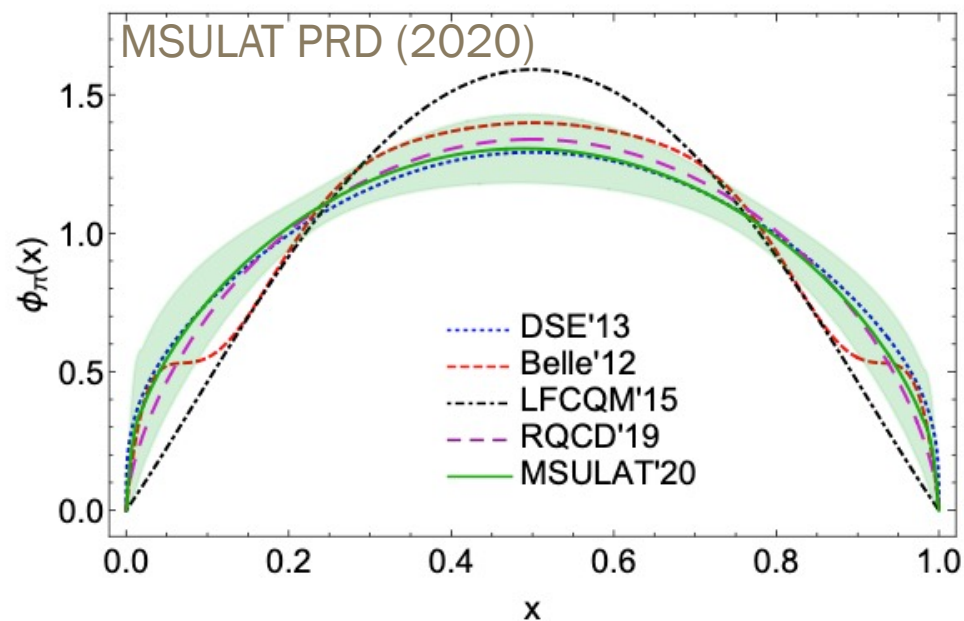
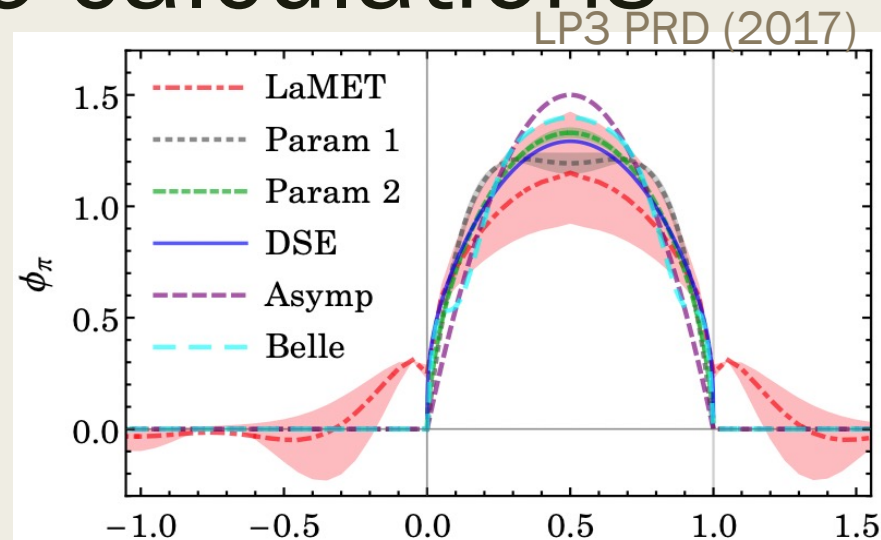
# Progress in $x$ -dependence calculations

- Lattice efforts
  - Boost to large momentum
  - Continuum extrapolation
  - Physical pion mass

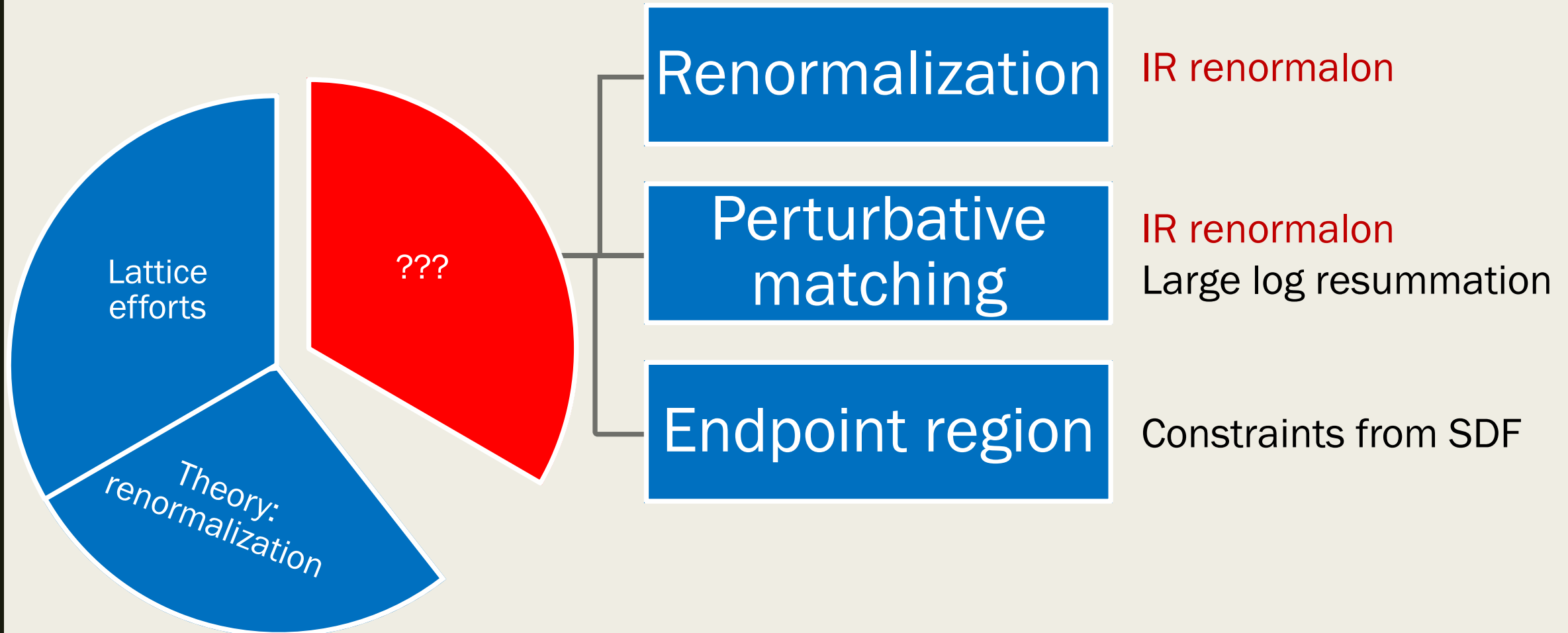


# Progress in $x$ -dependence calculations

- Lattice efforts
  - Boost to large momentum
  - Continuum extrapolation
  - Physical pion mass
- Theory: renormalization scheme



# What is still missing?



# Outline

Introduction to Pion DA

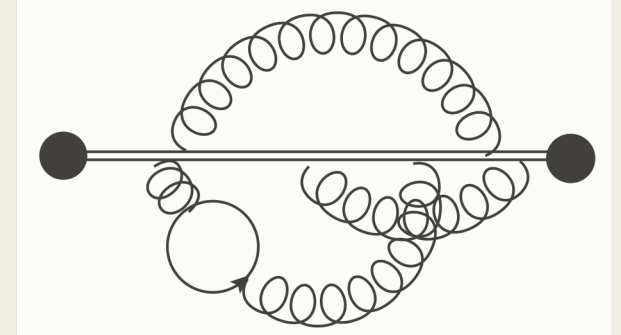
Renormalization of DA

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# Why is renormalization necessary

- Non-local operator:  $\hat{O}(z) = \bar{q}(z)W(z, 0)\Gamma q(0)$
- Wilson line self energy:  $\delta m \sim \frac{1}{a}$ , linear divergence!
- Multiplicative renormalization



$$\hat{O}^R(z, a) = \hat{O}^{bare}(z, a) / Z^R(z, a)$$

$$a \rightarrow 0 \text{ 😊}$$

$$a \rightarrow 0 \text{ 😰}$$

$$\sim e^{\delta m z}$$

# Renormalization with lattice data

Use lattice data with the same divergence

➤ Ratio scheme:

Radyushkin, PRD (2017)

$$Z^R(z, a) = \langle P_z = 0 | \hat{O}(z) | 0 \rangle$$

The lattice correlator vanishes for  $P_z = 0$ , not applicable

➤ RI/MOM scheme:

Martinelli, et al. NPB (1995)  
Zhang, et al. PRD (2020)

$$Z^R(z, a) = \frac{\langle q | \hat{O}(z) | q \rangle}{\langle q | \hat{O}(z) | q \rangle_{\text{tree}}}$$

Extra non-perturbative effects at large  $z$ .

The scheme conversion to  $\overline{\text{MS}}$  scheme at large  $z$  is not perturbative.

# Improvements from perturbation theory

## ➤ Self renormalization:

LPC, NPB (2021)

- Fit the  $a$  dependence from  $P_z = 0$  lattice data
- Match the lattice data to perturbative results in  $\overline{\text{MS}}$  scheme

## Hybrid scheme framework:

Short distance  $|z| < z_s : Z^R(z, a)$  from another scheme

Ji, et al., NPB (2021)

Long distance  $|z| > z_s : Z^R(z_s, a) e^{\delta m(|z| - z_s)}$

LPC, PRL (2022)

The scheme conversion to  $\overline{\text{MS}}$  scheme is perturbative in all regions.

## How to determine $\delta m$ ?

- Perturbation theory/Fit lattice data (from self-renormalization)?

# Outline

Introduction to Pion DA

Renormalization

Power Accuracy in LaMET

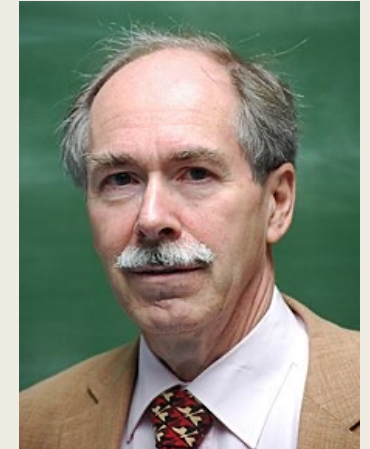
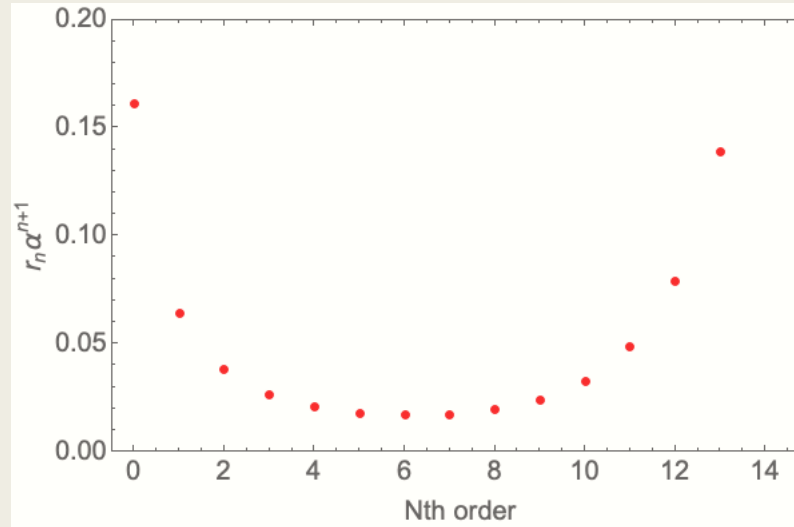
Application to DA Analysis



# Renormalon in perturbation series

- $\delta m = \frac{1}{a} \sum \alpha_s^{n+1} (a) r_n$
- At higher orders:
  - $r_n \sim n!$
  - Divergent for any  $\alpha_s$
  - No well-defined sum

Data from Bali, et al. PRD(2013)



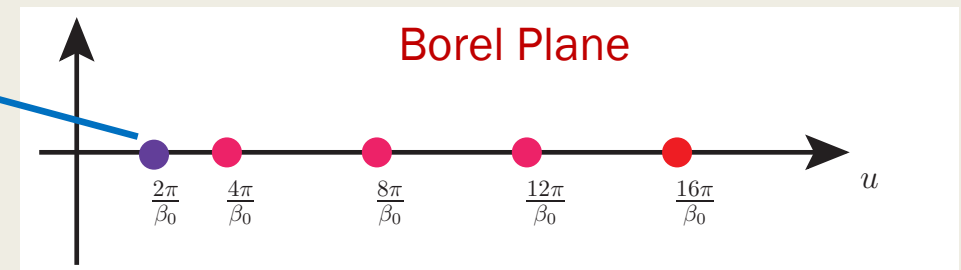
Gerard 't Hooft  
1999 Nobel Prize

$n!$  growth comes from IR renormalon

Need a regularization!

Ambiguity  $\Delta(\delta m) \sim O(\Lambda_{\text{QCD}})$

The matching  $C(x, y, \mu, P_Z)$  also contains renormalons



Braun, et al. PRD (2019)

# Ambiguity in extracting $\delta m$ from data

- Following the self-renormalization, we can parametrize the **linear divergence** and **logarithmic divergence** in  $a$  for  $C_0(z, a)$

LPC, NPB (2021)

$$C_0(z, a) = \exp[g(z)] \exp\left[\frac{kz}{a \ln a\Lambda}\right] \exp\left[\frac{3C_F}{\beta_0} \ln\left(\ln\frac{1}{a\Lambda}\right) + \ln\left(1 + \frac{d}{\ln a\Lambda}\right)\right]$$

Physical  $z$  dependence

- When fitting to lattice data,  $k$  and  $\Lambda$  are correlated and uncertain

$$\Delta(\delta m) = \Delta\left(\frac{k}{a \ln a\Lambda}\right) \sim O(\Lambda_{\text{QCD}})$$

This uncertainty comes from the renormalon ambiguity

$$\Delta(g(z)) \sim O(z\Lambda_{\text{QCD}})$$

The physical  $z$  dependence is dependent on our choice of  $k$  and  $\Lambda$

# Power Accuracy

- Leading renormalon ambiguity results in  $\mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right)$  correction in the LaMET matching

$$\tilde{\phi}(x, P_z) = \int_0^1 dy C(x, y, \mu, P_z) \phi(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{x^2 P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{(1-x)^2 P_z^2}\right)$$

- $P_z \rightarrow \infty$ :  $\mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right)$  not important
- $P_z \sim \text{GeV}$ :  $\mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right)$  dominate the power correction

Need an approach to eliminate  $\mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right)$  in the matching: Power Accuracy!

# How to Achieve Power Accuracy?

- LaMET Matching:

$$\tilde{\phi}(x, P_z) = \int_0^1 dy C(x, y, \mu, P_z) \phi(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{x^2 P_z^2}\right)$$

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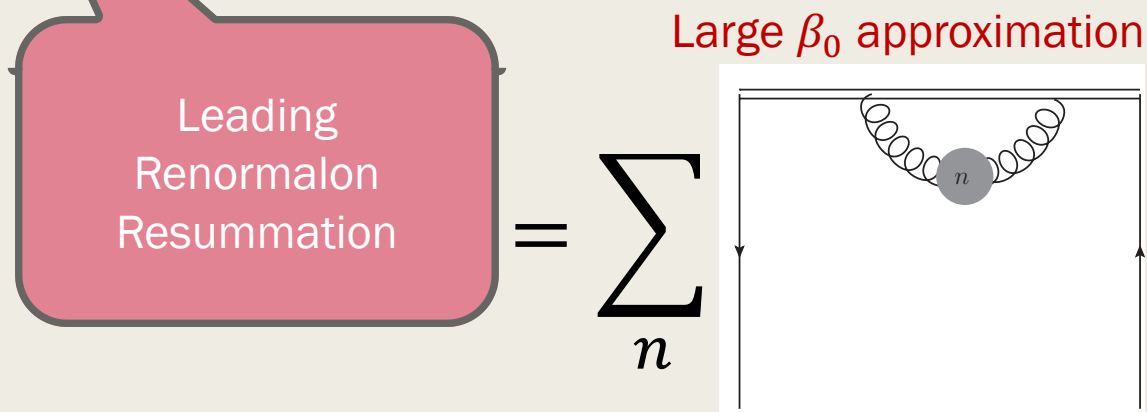


Leading  
Renormalon  
Resummation

# How to Achieve Power Accuracy?

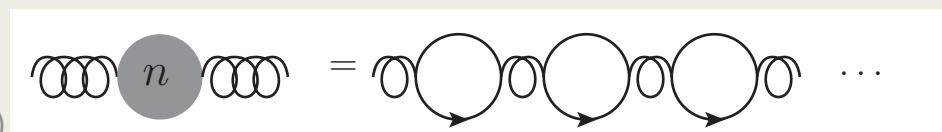
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Why large- $\beta_0$  approximation?

1. Pert theory very difficult at high order
2. The approximation is good in  $\delta m$  Bali, et al. PRD(2013)
3. Independent of Dirac structure and hadron momentum



Braun, et al. PRD (2019)

# How to Achieve Power Accuracy?

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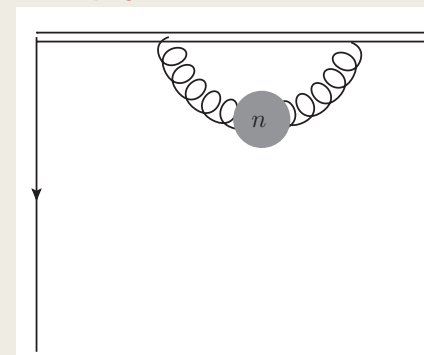
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Renormalization  
with  $e^{\delta m \cdot z + m_0 \cdot z}$

Leading  
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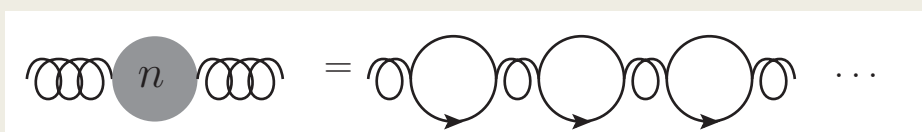
$$= \sum_n$$

Large  $\beta_0$  approximation



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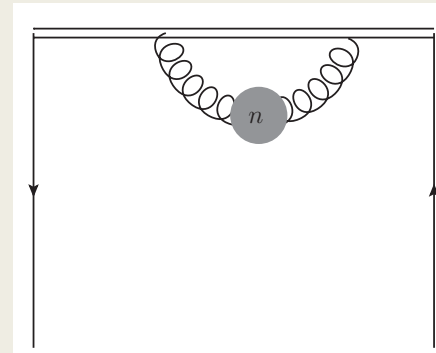
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Renormalization  
with  $e^{\delta m \cdot z + m_0 \cdot z}$

Leading  
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# How to calculate $m_0$

- Work with  $P_z = 0$  lattice data LPC, NPB (2021) Perturbative calculation

$$e^{m_0 \cdot z} e^{\delta m \cdot z} \langle P_z = 0 | \bar{q}(z) \gamma_t W(z, 0) q(0) | P_z = 0 \rangle = C_0^{\overline{MS}}(z, a)$$

Lattice data

# How to calculate $m_0$

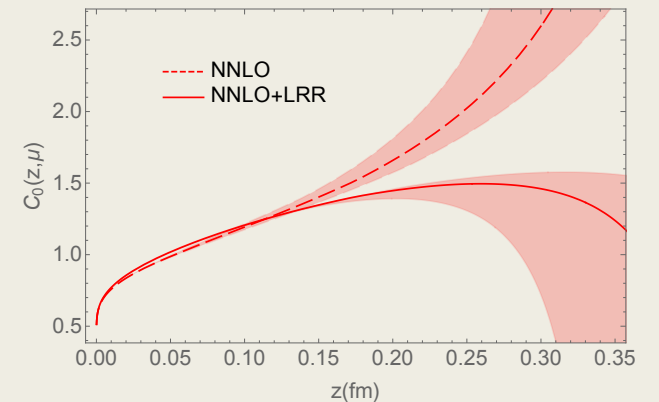
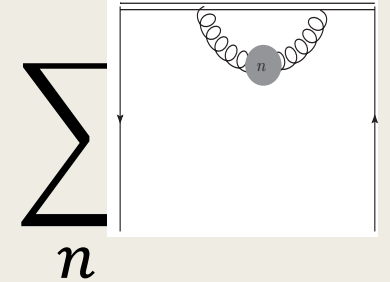
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Perturbative calculation

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Lattice data

Large  $\beta_0$  approximation

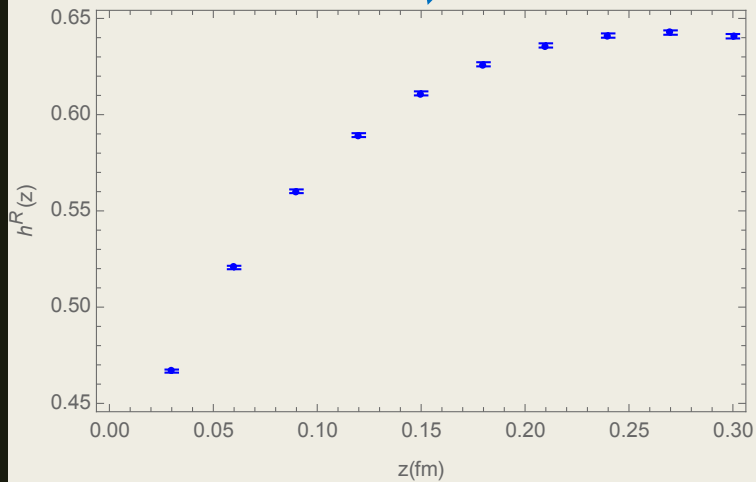


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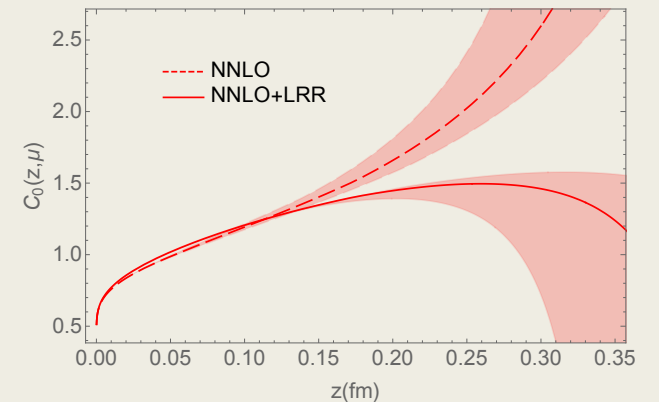
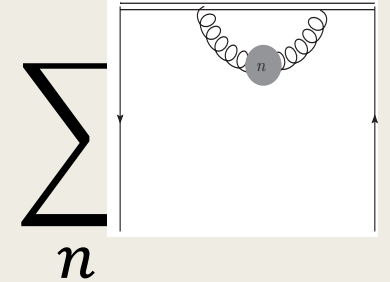
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Perturbative calculation

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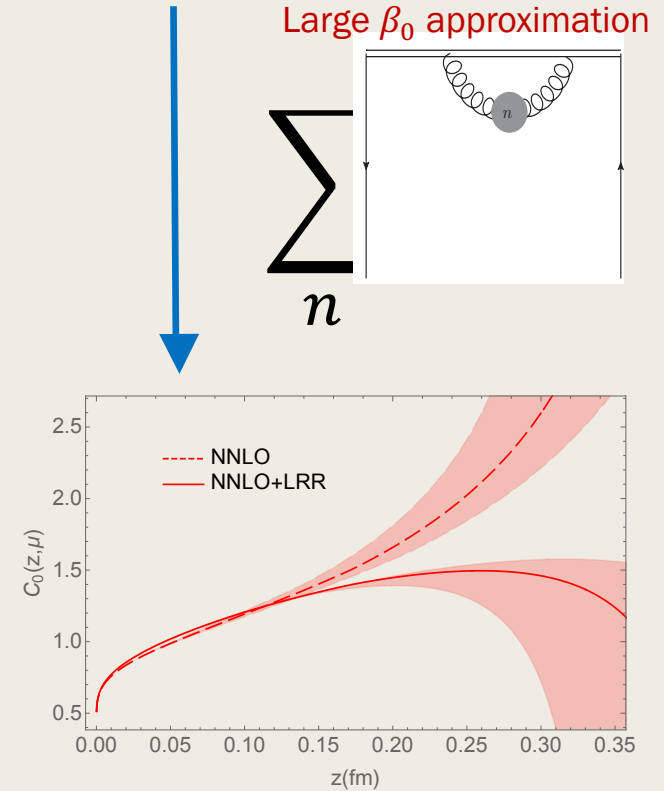
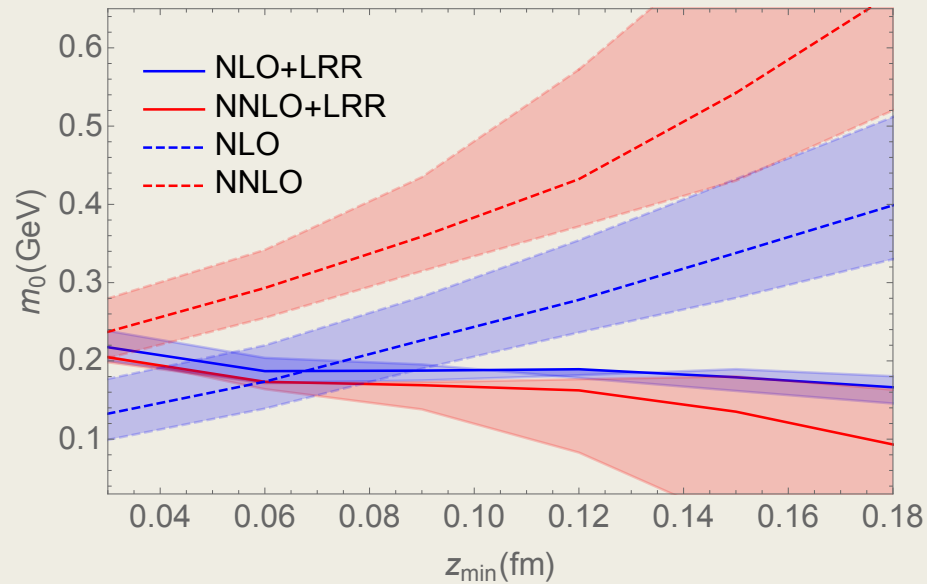
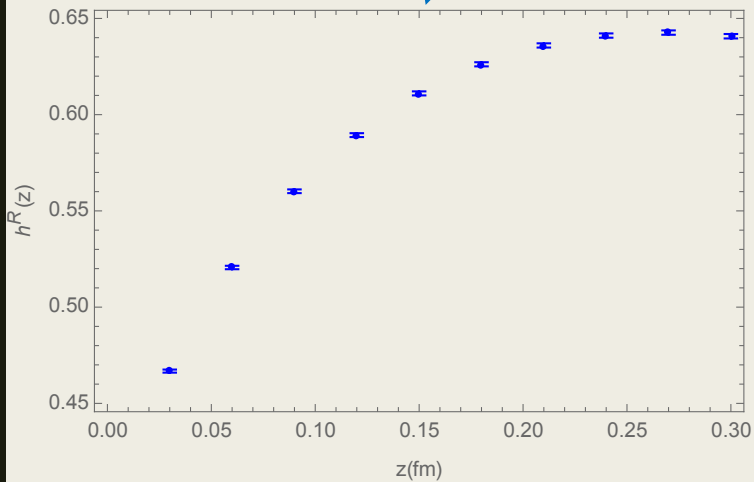
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Lattice data

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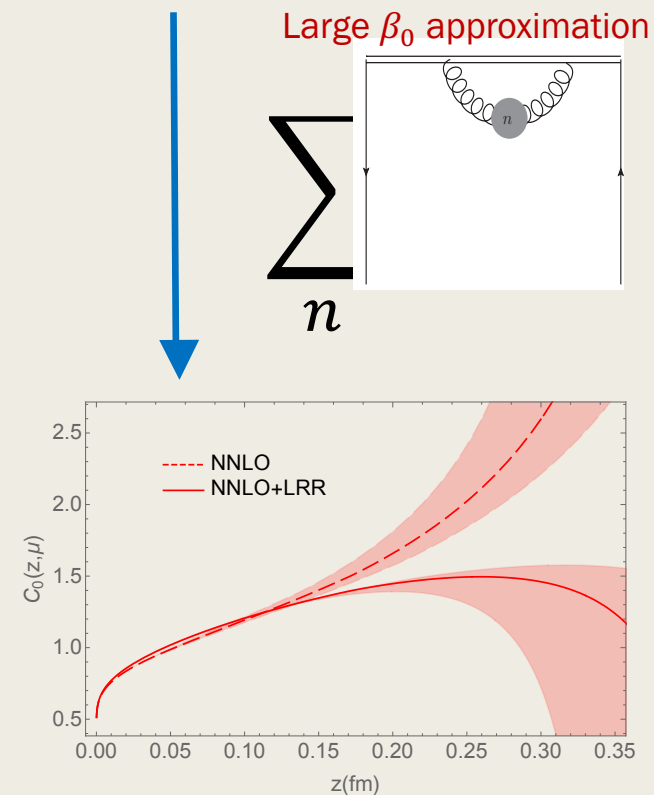
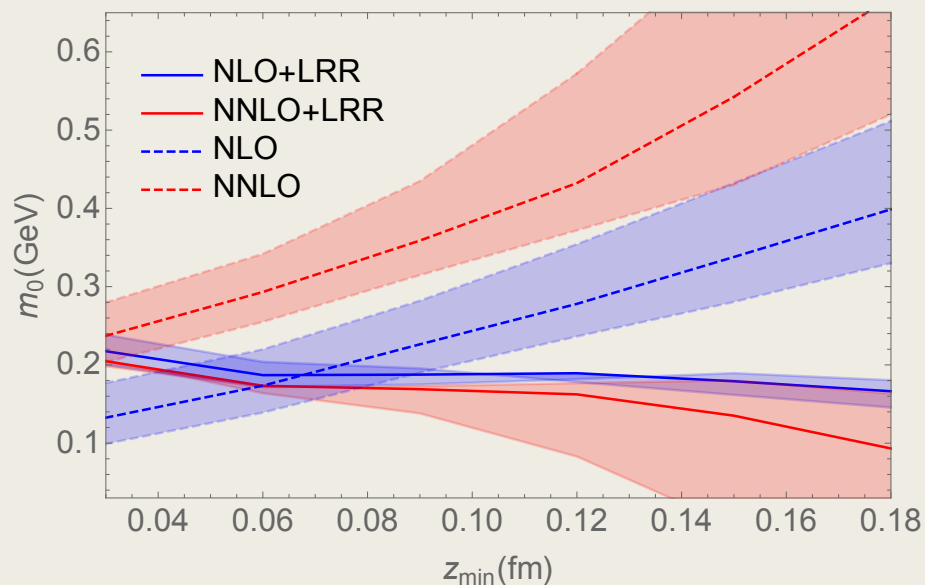
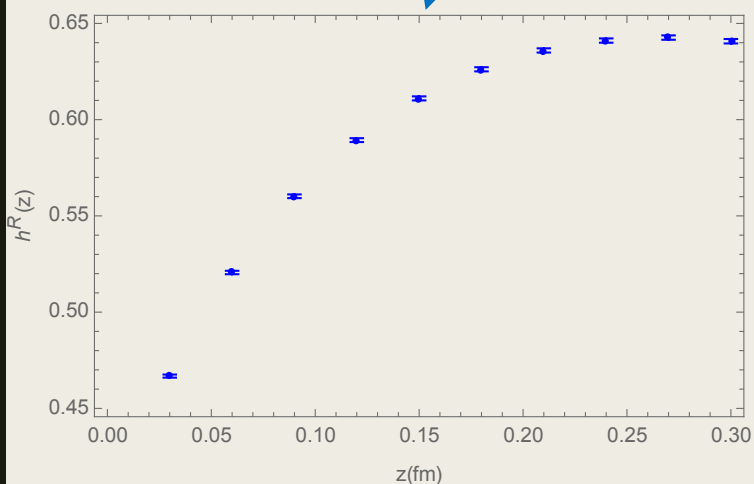
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Lattice data

Large  $\beta_0$  approximation



Such an extraction is independent of Dirac structure and external state's momentum

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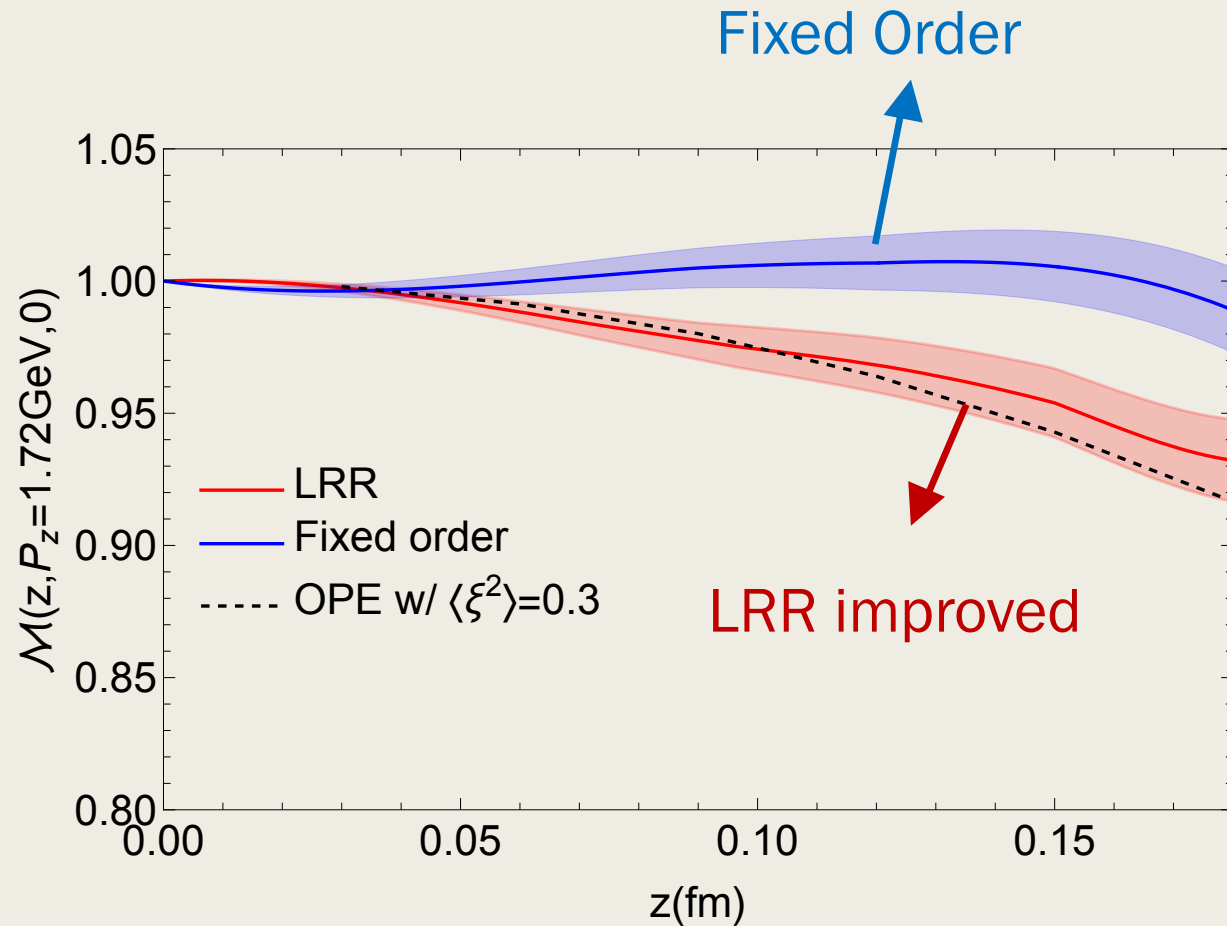
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# LRR-improved renormalization



- Fixed order renormalization suggests a hump at small  $z$
- LRR improved renormalization suggests a decaying distribution

Can be tested in short distance OPE



# Testing renormalization through moments extraction

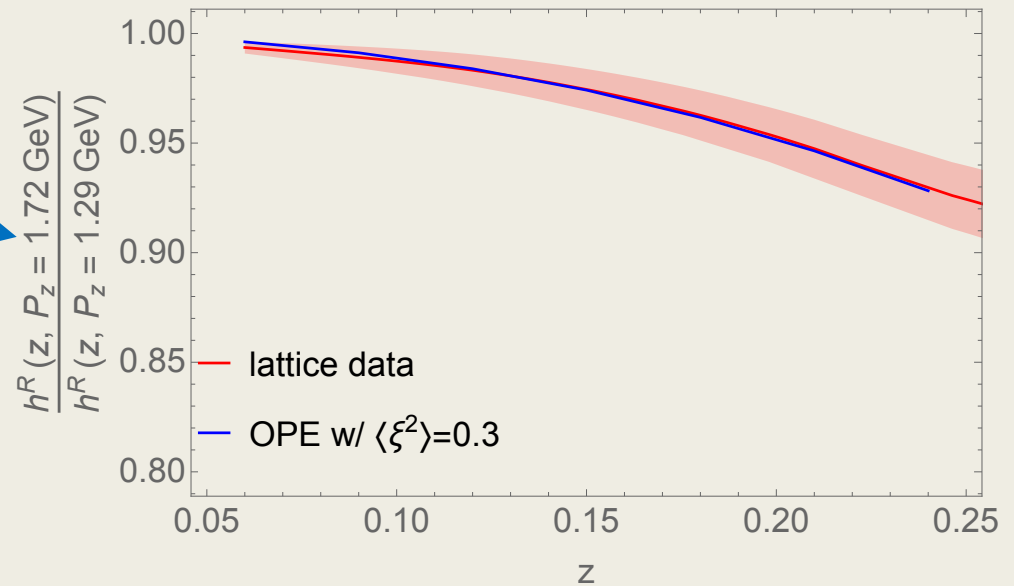
$$\text{OPE } (z \ll \Lambda_{\text{QCD}}^{-1}) : \quad \tilde{h}^R(z, P_z, \mu) = \sum_{n=0} \frac{\left(-\frac{izP_z}{2}\right)^n}{n!} \sum_{m=0}^n C_{nm}(z^2\mu^2) \langle \xi^m \rangle + O(z^2\Lambda_{\text{QCD}}^2)$$

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**Method 1:** Gao, et al., PRD (2022)

- Ratio:  $\frac{\tilde{h}^B(z, P_1)}{\tilde{h}^B(z, P_2)}$
- Renormalization independent
- $\langle \xi^2 \rangle = 0.298(39)$



# Testing renormalization through moments extraction

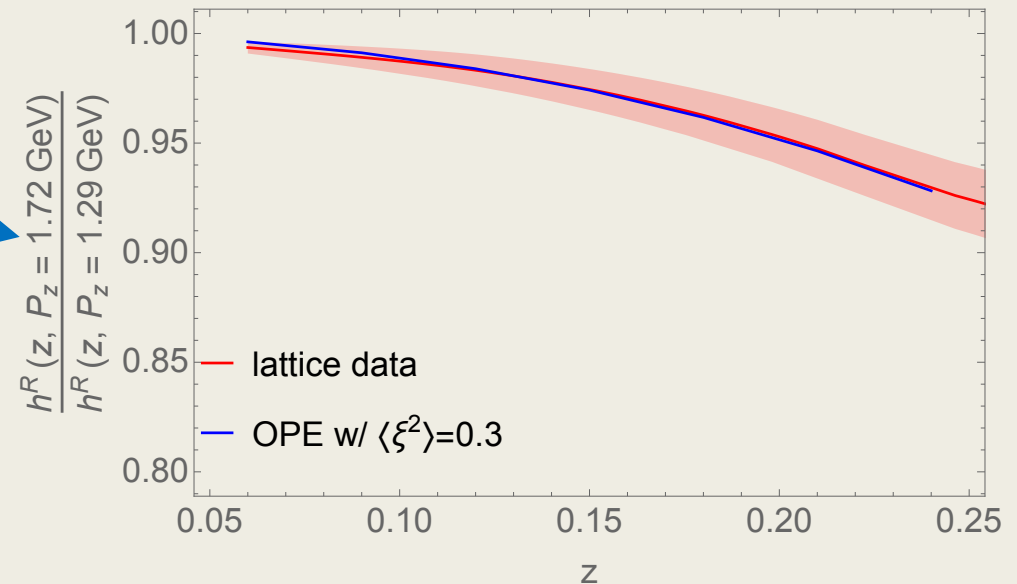
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**Method 1:** Gao, et al., PRD (2022)

- Ratio:  $\frac{\tilde{h}^B(z, P_1)}{\tilde{h}^B(z, P_2)}$
- Renormalization independent
- $\langle \xi^2 \rangle = 0.298(39)$

**Method 2:**

- Renormalized ME:  $\tilde{h}^R(z, P_z, \mu)$
- Fixed-order result suggests  $\langle \xi^2 \rangle < 0$
- LRR result suggests  $\langle \xi^2 \rangle \approx 0.3$



# Testing renormalization through moments extraction

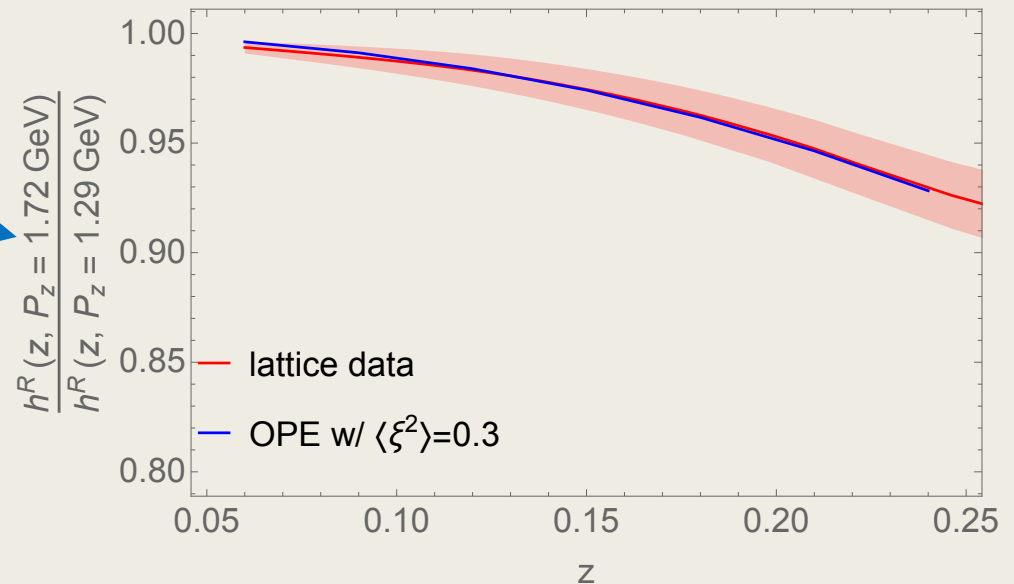
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Consistency suggests the LRR-renormalization is correct!

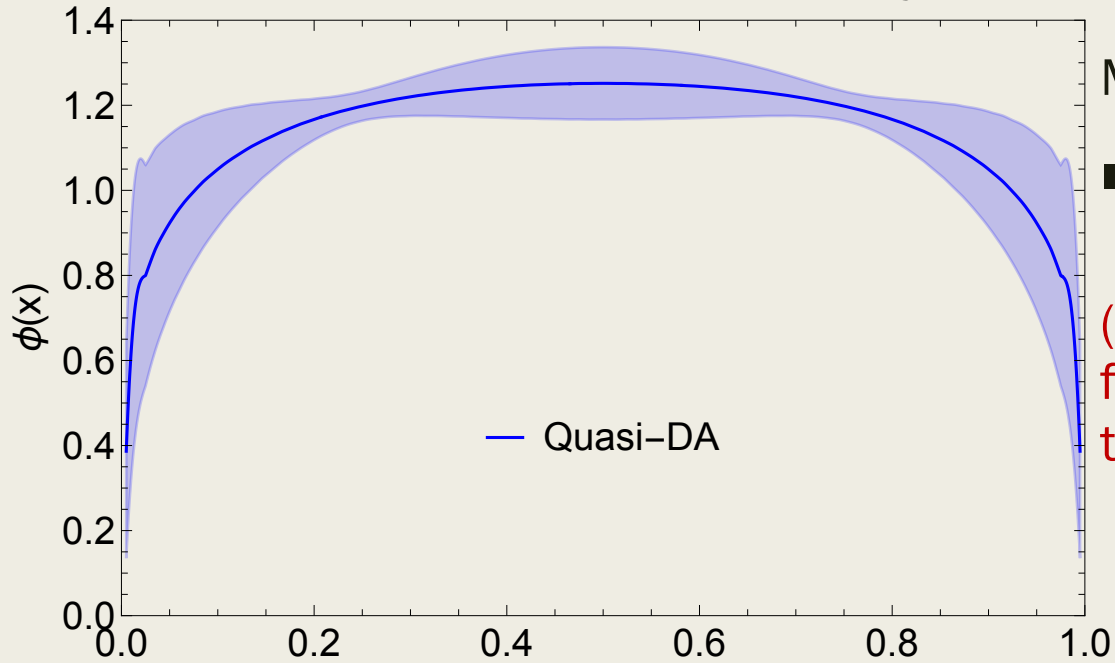
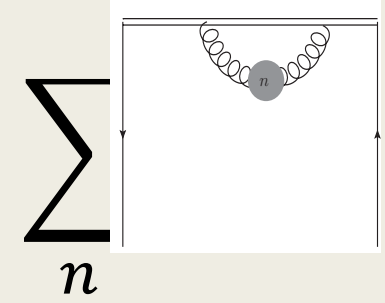
# How LRR improves matching?

The diagrams are the same for different momenta/Dirac structure:

$$\Delta Q(z, P_z, \mu) = \left( C_0^{\text{LRR}}(z, \mu) - C_0(z, \mu) \right) e^{-iyzP_z}$$

$$\Delta C(x, y, \mu, P_z) = \int_{|z| > z_s} \frac{dz}{2\pi} e^{-ixzP_z} \Delta Q(z, P_z)$$

Large  $\beta_0$  approximation



Matching with LRR (the same quasi-DA):

- LRR introduces correction to the matching, more important near endpoints

(This is just the effect of LRR matching, not the final result, which will be presented in Jack's talk.)

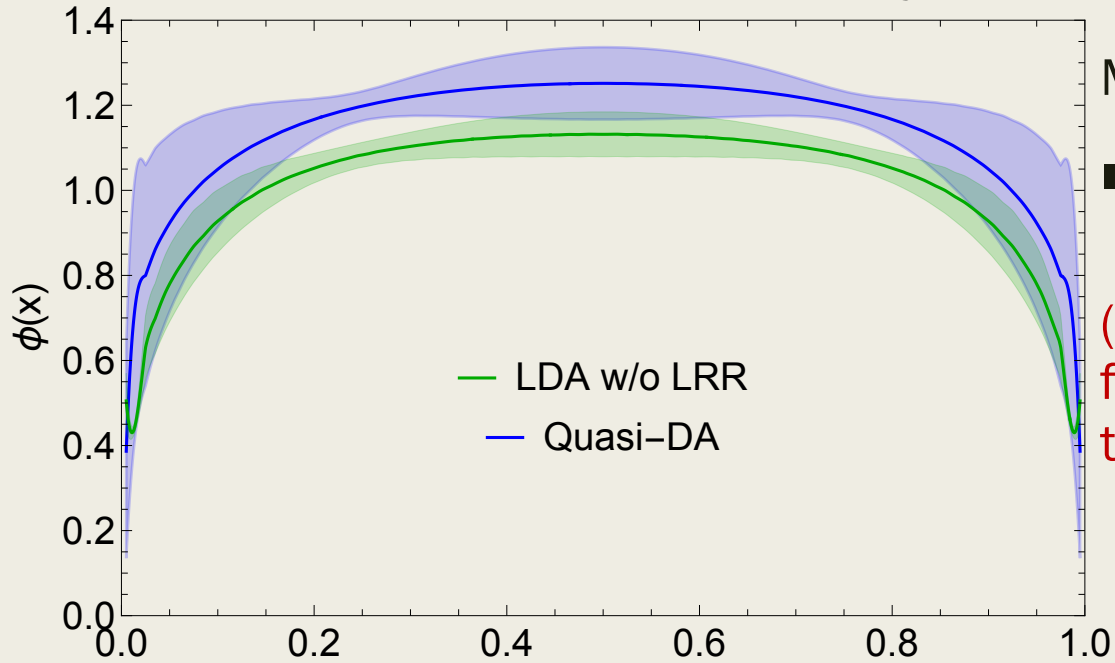
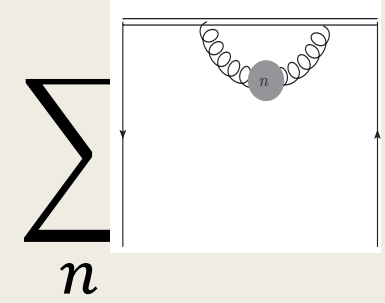
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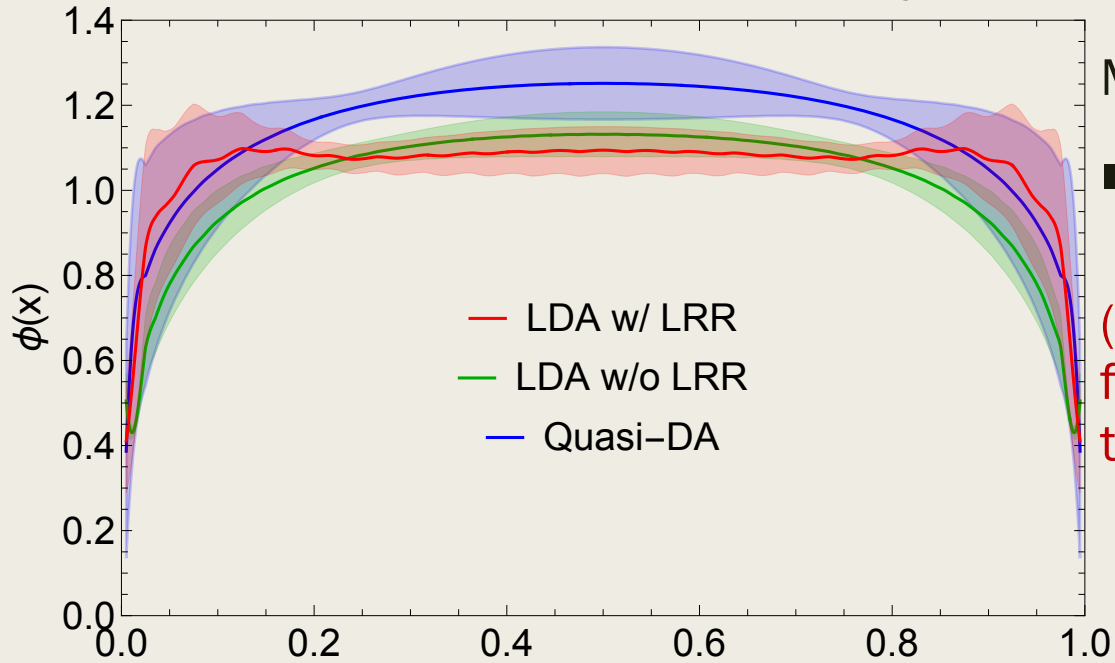
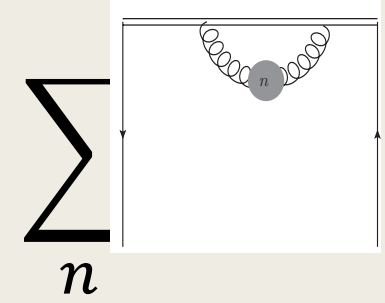
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# Conclusion

- Renormalons exist in the renormalization of the LaMET operator, and the perturbative matching kernel
- The leading renormalon results in an  $O\left(\frac{\Lambda_{\text{QCD}}}{xP_z}\right)$  correction to the LaMET factorization
- We perform a leading renormalon resummation in the large- $\beta_0$  limit and use a term  $m_0^{\text{eff}}$  in renormalization to eliminate the correction
- The LRR-improved renormalization improves the lattice calculation of short distance correlations, consistent with OPE
- A corresponding modification is made in the matching kernel





**THANK YOU!**