Unpolarized gluon PDF for the proton using the twisted mass formulation

Results

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Results

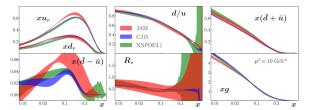
Outline

- 1 Introduction
 - Motivation
 - gPDFs on the Lattice
- 2 Calculation of gluon PDFs
 - The Pseudo-PDF Approach
 - Lattice Setup
- 3 Results
 - Matrix Elements
 - Double Ratio and Interpolation
 - ITD Development
 - Reconstructed Pseudo-PDF
 - Ongoing Work
 - Summary Future Work
- 4 Acknowledgements

Motivation

- Gluon contributions to physical quantities play a critical role in hadron structure
- Gluon contributions can be large, eg. gluon momentum fraction $\approx 40\%$
- Dedicated experimental efforts to understand gluonic structure of hadron

[Moffat et al, PRD 104, 016015 (2021)] [Ball et al, EPJC 77, 663 (2017)] [Accardi et al, PRD 93, 114017 (2016)]



Lattice studies of gPDFs can assist in constraining global analysis

gPDFs on the Lattice

- Several challenges in extracting reliable results
 - purely disconnected diagram
 - at least an order of magnitude more statistics than quark counterparts
 - unavoidable mixing with quark singlet PDFs
- x-dependence of gluon PDFs even more challenging
- Inverse problem in reconstruction of x-dependence due to limited lattice data

4/22

Pseudo-PDF Approach (in a nutshell)

- Matrix elements of non-local operators and momentum-boosted proton states
 - Several choices for the form of the gluon operator consisting of two field-strength tensors, separated by spatial distance z, and two straight Wilson lines, connecting points 0 → z and z → 0

$$M_{\mu i;\nu j}(P,z) = \langle N(P)|F_{\mu i}(z)W(z,0)F_{\nu j}(0)W(0,z)|N(P)\rangle$$

- lacksquare Choice of indices for $F_{\mu
 u}$ not unique
- This operator avoids finite mixing under renormalization
 - must subtract vacuum expectation value

$$\mathcal{O} = \frac{1}{2} \sum_{i \neq 3} F_{i3}(x + z\hat{z}) W(x + z\hat{z}, x) F_{i3}(x) - \sum_{i \neq j \neq 3} F_{ij}(x + z\hat{z}) W(x + z\hat{z}, x) F_{ij}(x)$$

- Matrix elements extracted from ratio of 2pt- and 3pt- functions
- Ground state from plateau fit

$$\frac{C^{3pt}(t,\tau,0,\vec{P})}{C^{2pt}(t,0,\vec{P})} \stackrel{0<<\tau<< t}{=} \mathcal{M}(\nu,z^2)$$

Pseudo-PDF Approach (in a nutshell)

- \blacksquare \mathfrak{M} Euclidean, renormalized in ratio scheme, scale 1/z
 - lacksquare want: Minkowski, renormalized in $\overline{\mathrm{MS}}$, scale μ
 - we need matching to the light-cone, scheme conversion and evolution
- Form the double ratio (reduced ITD) with zero-momentum and local matrix elements for renormalization and to reduce higher twist contributions [Orginos et al., Phys.Rev.D 96 (2017) 9, 094503]

$$\mathfrak{M}(\nu, z^2) \equiv \left(\frac{M(\nu, z^2)}{M(\nu, 0)|_{z=0}}\right) / \left(\frac{M(0, z^2)|_{p=0}}{M(0, 0)|_{p=0, z=0}}\right)$$

- Scale evolution and apply matching kernel on ITD
 - neglect mixing with quark singlet
 - \blacksquare normalize with $\langle x \rangle_g$

$$\mathcal{Q}(\nu, z^2, \mu^2) = \mathfrak{M} + \frac{\alpha_s N_c}{2\pi} \int_0^1 du \, \mathfrak{M}(u\nu, z^2) \left\{ \ln\left(\frac{z^2 \mu^2 e^{2\gamma_E}}{4}\right) B(u) + 4\left[\frac{u + \ln(\overline{u})}{\overline{u}}\right]_+ + \frac{2}{3} \left[1 - u^3\right]_+ \right\}$$

■ Reconstruct x-dependence (Backus-Gilbert, fitting ansatz, Fourier tansform, etc.)

$$Q(\nu, z^2, \mu^2) = \int_0^1 dx \cos(x\nu) x g(x, \mu^2)$$

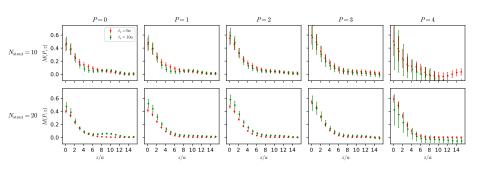
Lattice Parameters and Statistics

- $N_f = 2+1+1$ ensemble of twisted-mass clover fermions and Iwasaki improved gluons
 - $m_{\pi} = 260 \text{ MeV}$
 - a = 0.093 fm
 - $L^3 \times T = 32^3 \times 64$
 - \blacksquare $Lm_{\pi}=4$
- Stout smearing ($\omega = 0.129$)
 - field-strength tensor: 10, 20 steps
 - Wilson line: 0, 10 steps
- Momentum smearing (optimized value $\xi = 0.6$) used for P = 2, 3, 4 [Bali et al, PRD 93, 094515 (2016)]
- Excited states:
 - Numerical results relatively good up to $t_s = 10a$
- Statistics
 - Average over all 6 spatial directions of Wilson line / momentum $(\pm x, \pm y, \pm z)$
 - Statistics much higher than quark PDFs

$ \mathbf{P_3} \left[\frac{2\pi}{L}\right]$	P ₃ [GeV]	$N_{\rm confs}$	N _{src}	N _{dir}	Total statistics
0 to 4	0 - 1.67	1,134	200	6	1,360,800

Matrix Elements: Excited States Contamination and Effect of N_{stout}

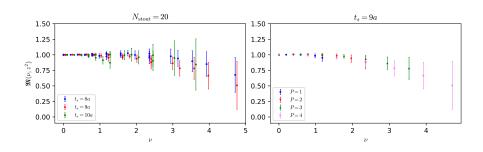
Results



- Various values of t_s and two stout steps for the gluon operator
- Statistical errors increase with momentum boost and ts
- MEs have expected behavior (higher boosts decay faster to 0)
- Final results use $N_{\rm stout} = 20$

Double Ratio (Reduced ITD)

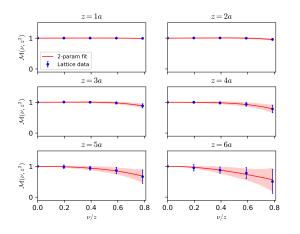
Double-ratio analysis to narrow down ts



- Values at different t_s are compatible within uncertainties
- Final results use $t_s = 9a$
 - $z_{max} = 6a = 0.568$ fm motivated by signal
- Lattice data form a smooth function
 - Must interpolate for evolution and matching

Interpolation of Double Ratio

- We interpolate the double-ratio at each z to get a continuous function for the integration
 - interpolation done with linear and 2nd-order polynomial fits

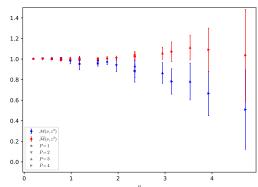


- 2nd-order polynomial fits prove to be the most suitable for evolution and matching
- Choice of fit mostly irrelevant below z = 4a

ITD Development

 \blacksquare Apply the evolution kernel to the reduced matrix elements to the scale $\mu=2$ GeV ahead of final conversion to $\overline{\rm MS}$ scheme and matching to the light-cone

$$\tilde{\mathfrak{M}}(\nu,z^2,\mu^2)=\mathfrak{M}+\frac{\alpha_{s}\textit{N}_{\textit{c}}}{2\pi}\int_{0}^{1}\textit{d}u\ln(\frac{z^2\mu^2\textrm{e}^{2\gamma_{\textit{E}}}}{4})\textit{B}(u)\mathfrak{M}(u\nu,z^2)$$

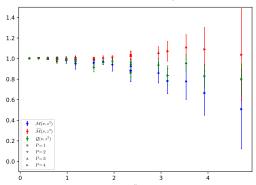


- Data from different (P, z) pairs fall on a universal curve
- We find good agreement up to z = 6a

ITD Development

 \blacksquare Apply the matching kernel to convert to $\overline{\mathrm{MS}}$ scheme and match to the light-cone

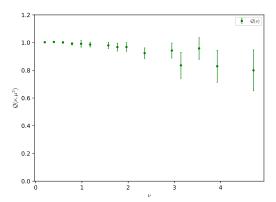
$$\mathcal{Q}(\nu, z^2, \mu^2) = \tilde{\mathfrak{M}}(\nu, z^2, \mu^2) + \frac{\alpha_s N_c}{2\pi} \int_0^1 du \ L(u) \mathfrak{M}(u\nu, z^2)$$



- We continue to find good agreement between common values of loffe time from different combinations of momenta and Wilson line lengths
- Matching effects in opposite direction of evolution

ITD Development

lacksquare Average over common u for final pseudo-ITD



 \blacksquare No information remains regarding initial (P, z) pairs

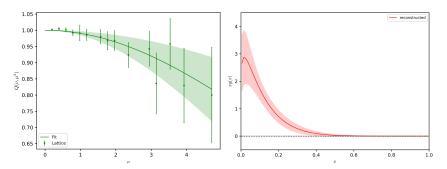
PDF Reconstruction

Introduction

Reconstructed Pseudo-PDF

The fit is chosen by the minimization of

$$\chi^2 = \sum_{\nu=0}^{\nu_{max}} \frac{\left(Q(\nu,\mu^2) - \int_0^1 dx \cos(\nu x) N \, x^a (1-x)^b\right)^2}{\sigma_Q^2(\nu,\mu^2)}$$



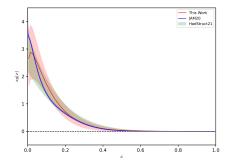
- PDF is normalized using gluon momentum fraction $\langle x \rangle_a^{\overline{\rm MS}} (\mu = 2 \, {\rm GeV}) = 0.427(92)$ [Alexandrou et al, PRD 101, 094513 (2020)]
- Other reconstruction methods (naive Fourier-transform, Backus-Gilbert method) have proven less suitable [Bhat et al, PRD 103, 034510 (2021)] gPDF from Lattice QCD

Comparison with Other Works

Introduction

Reconstructed Pseudo-PDF

- Comparison with lattice results from HadStruc collaboration [Khan et al, PRD 104, 094516 (2021)]
 - $m_{\pi} = 358$ MeV, a = 0.094, $L^3 \times T = 32^3 \times 64$
- JAM20 global analysis [Moffat et al, PRD 104, 016015 (2021)], $\langle x \rangle_q = 0.40(1)$



- We find agreement between all results
- This work: $\nu_{max} = 4.71$
- HadStruc: $\nu_{max} = 7.07$

Quark Singlet Contribution

- All calculations neglect the mixing between quark-singlet and gluon PDFs (singlet-guark PDFs challenging to calculate)
- A proper analysis requires inclusion of the singlet-quark contributions
- The mixing for non-local operators appears in the matching equations:

$$\begin{split} \mathfrak{M}(\nu, \boldsymbol{z}^2, \boldsymbol{\mu}^2) &= \frac{\mathcal{I}_g(\nu, \boldsymbol{\mu}^2)}{\mathcal{I}_g(0, \boldsymbol{\mu}^2)} - \frac{\alpha_s N_c}{2\pi} \int_0^1 d\boldsymbol{u} \; \frac{\mathcal{I}_g(\boldsymbol{\mu}\nu, \boldsymbol{\mu}^2)}{\mathcal{I}_g(0, \boldsymbol{\mu}^2)} \bigg\{ \ln \bigg(\frac{^z 2\boldsymbol{\mu}^2 e^{2\gamma} E}{4} \bigg) B_{gg}(\boldsymbol{u}) + 4 \bigg[\frac{\boldsymbol{u} + \ln(\overline{\boldsymbol{u}})}{\overline{\boldsymbol{u}}} \bigg]_+ \\ &+ \frac{2}{3} \big[1 - \boldsymbol{u}^3 \big]_+ \bigg\} - \frac{\alpha_s C_F}{2\pi} \ln \bigg(\frac{^z 2\boldsymbol{\mu}^2 e^{2\gamma} E}{4} \bigg) \int_0^1 d\boldsymbol{w} \; \frac{\mathcal{I}_S(\boldsymbol{w}\nu, \boldsymbol{\mu}^2)}{\mathcal{I}_g(0, \boldsymbol{\mu}^2)} \mathfrak{B}_{gq}(\boldsymbol{w}) \end{split}$$

We attempt to extend our calculation by incorporating the single-guark PDFs calculated on the same ensemble

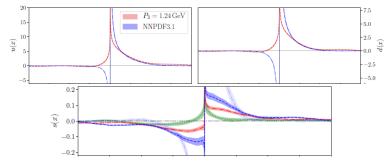
PHYSICAL REVIEW D 104, 054503 (2021)

Flavor decomposition of the nucleon unpolarized, helicity, and transversity parton distribution functions from lattice QCD simulations

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Karl Jansen, ⁴ and Floriano Manigrasso^{1,5,6}

Disconnected Matrix Elements

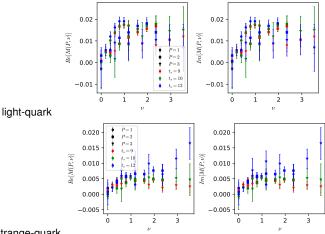
Published work of [C. Alexandrou et al., PRD 104 (2021) 5, 054503]
 uses quasi-PDFs method



Here, we combine the light-quark and strange-quark looks with the two point functions produced for the gluon PDF to increase statistics of disconnected diagram

Disconnected Matrix Elements

■ Light and strange disconnect matrix elements for $t_s = 9, 10, 12$

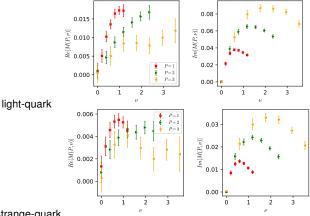


strange-quark

Large errors: compatibility between data on different t_s values

Disconnected Matrix Elements

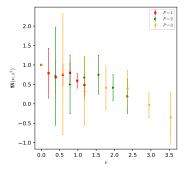
Focus on matrix elements for $t_s = 9$ for the pseudo-PDFs analysis



- strange-quark
- Dependence on momentum expected in matrix elements (ME)
- tiny strange-quark ME but non-zero signal achieved with advanced methods like hierarchical probing [A. Stathopoulos et al., arXiv:1302.4018]

Strange Quark Double Ratio

- Light-quark PDFs require inclusion of the connected diagram (in progress)
- Strange quark PDF is obtained exclusively from the disconnected diagram
- Reduced-ITD for strange quark contributions:



- Signal quality decays fast with increase of P and z.
- Study of correlations between matrix elements on-going

Summary - Future Work

Summary

- We presented a calculation of gluon PDFs using stout smearing to suppress the statistical noise
- We employ the pseudo-ITD method to overcome difficulties with direct evaluation of renormalization functions
- We study excited-state effects using various $T_{\rm sink}$ values
- Comparison of the unpolarized gluon PDF with other lattice results and global analysis shows compativle results
- First results on the quark singlet PDFs using pseudo-ITD (on-going)

Future work

- Completion of the analysis and elimination of mixing with singlet quark PDFs
- \blacksquare Address sources of systematic effects (excited states, ν dependence of fits)
- \blacksquare Study discretization effects: new ensemble at finer lattice spacing (a=0.08 fm) in production

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