



Beta decay recoil-order terms for studies of physics beyond the Standard Model

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Argonne National Laboratory
21 July 2024



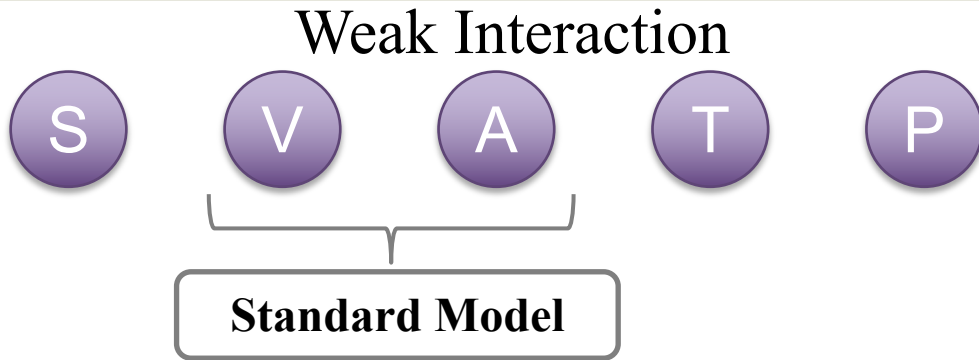
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ENERGY

Office of
Science

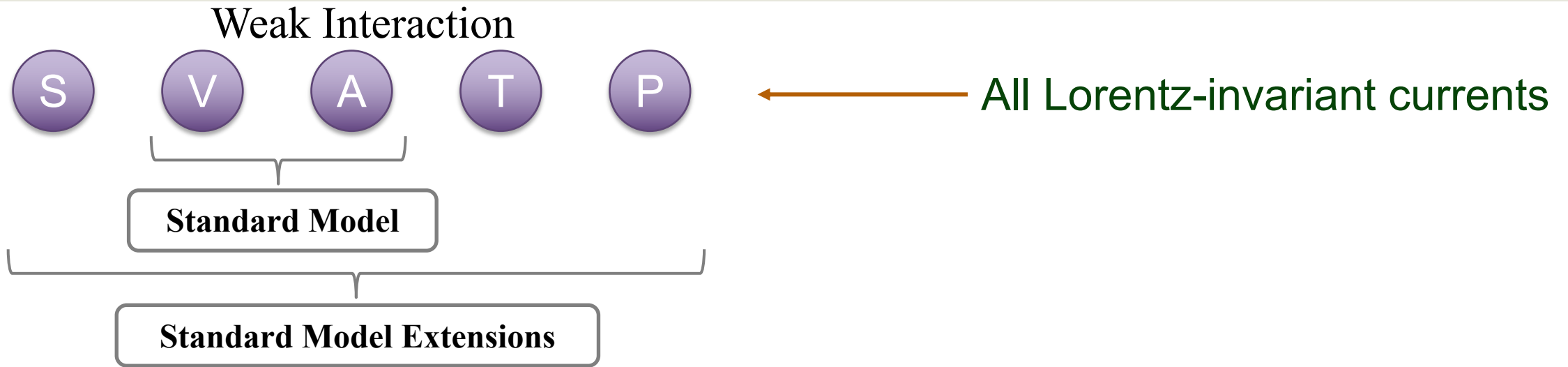
The structure of the Weak interaction



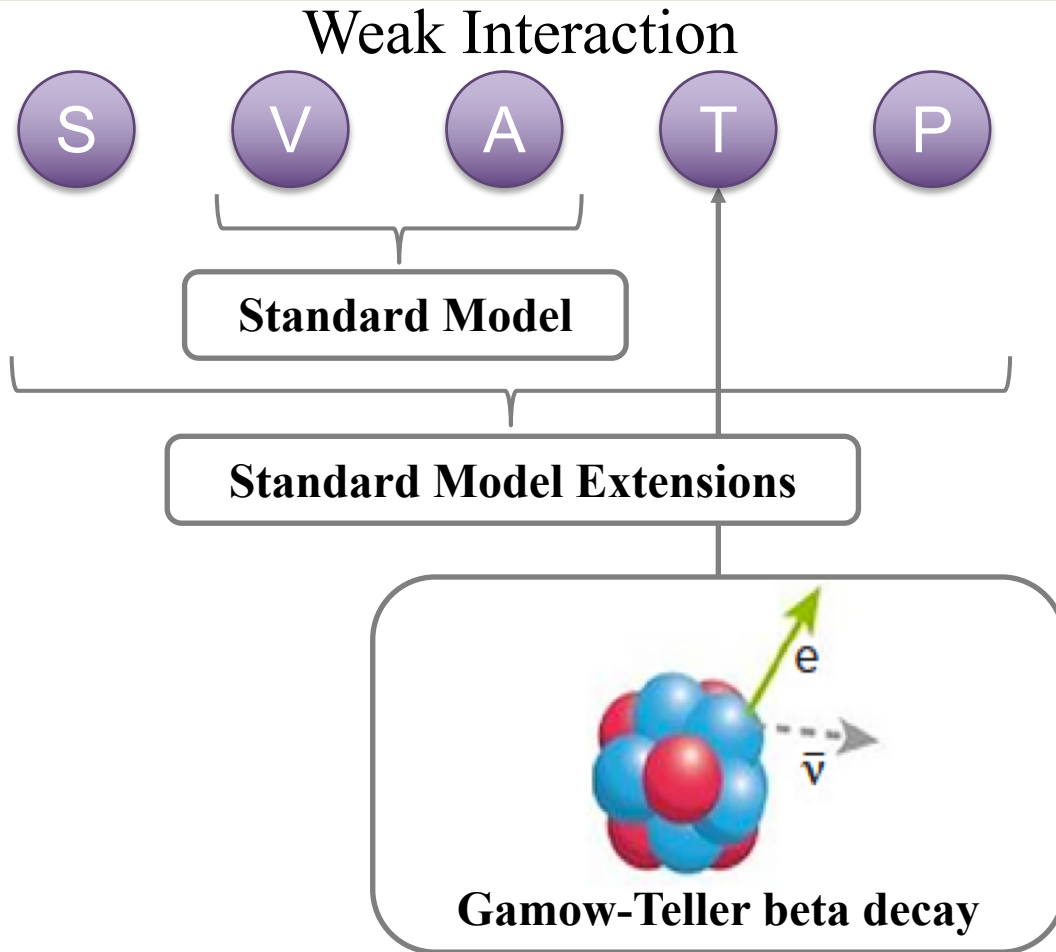
A series of β -decay experiments lead to the formulation of the V – A structure of the weak interaction:

- ❖ C. S. Wu, *et al.*, Phys. Rev. 105, 1413 (1957).
- ❖ W. B. Herrmannsfeldt, *et al.*, Phys. Rev. 107, 641 (1957).
- ❖ C. Johnson, *et al.*, Phys. Rev. 132, 1149 (1963).

The structure of the Weak interaction



The structure of the Weak interaction



← All Lorentz-invariant currents

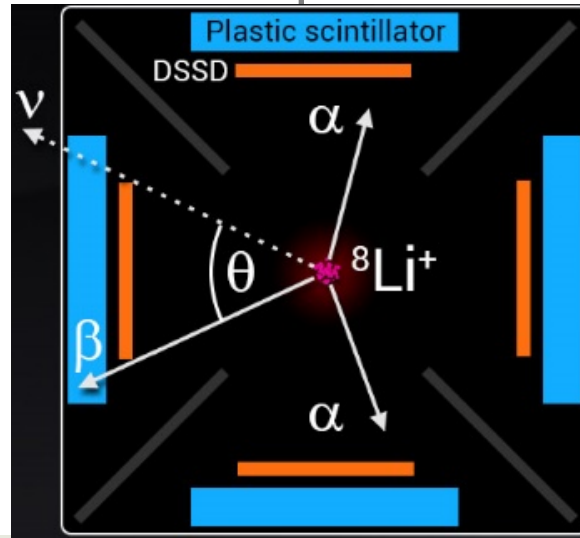
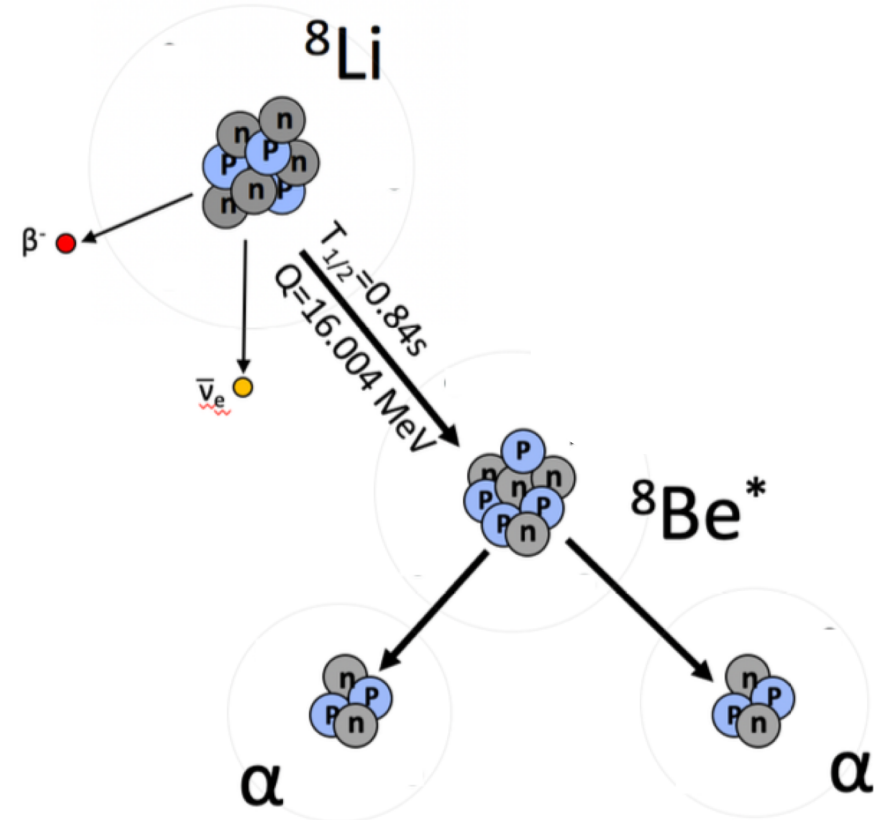
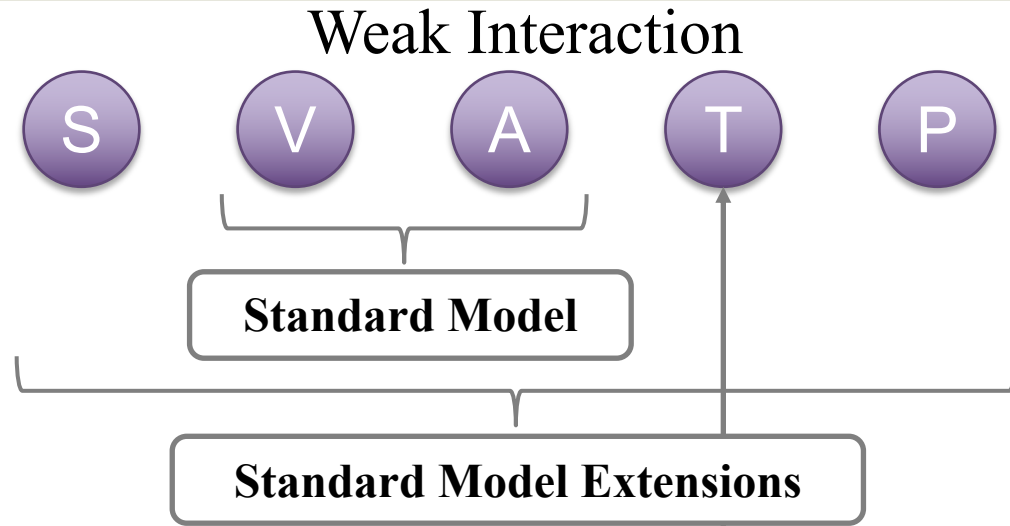
$$a_{\beta\nu} = -\frac{1 |C_A|^2 - |C_T|^2}{3 |C_A|^2 + |C_T|^2}$$

In Standard Model $a_{\beta\nu} = -\frac{1}{3}$

Before 2022 most precise constraint from corrected 1963 measurement of ${}^6\text{He}$ decay: $a_{\beta\nu} = -0.3308 \pm 0.0030$

C. Johnson, *et al.*, Phys. Rev. 132, 1149 (1963).
F. Gluck, Nucl. Phys. A 628, 493 (1998)

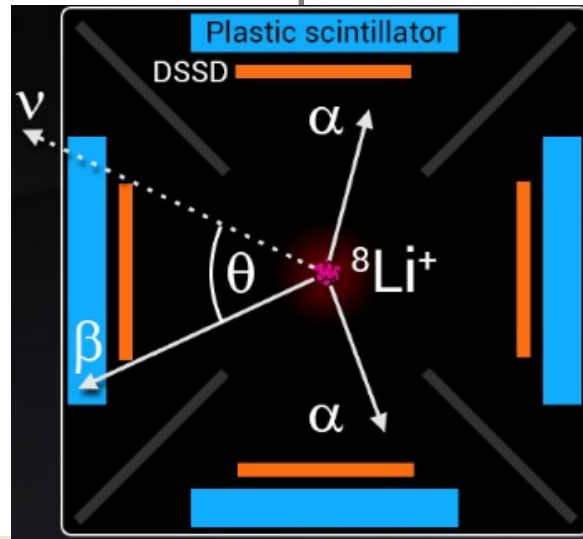
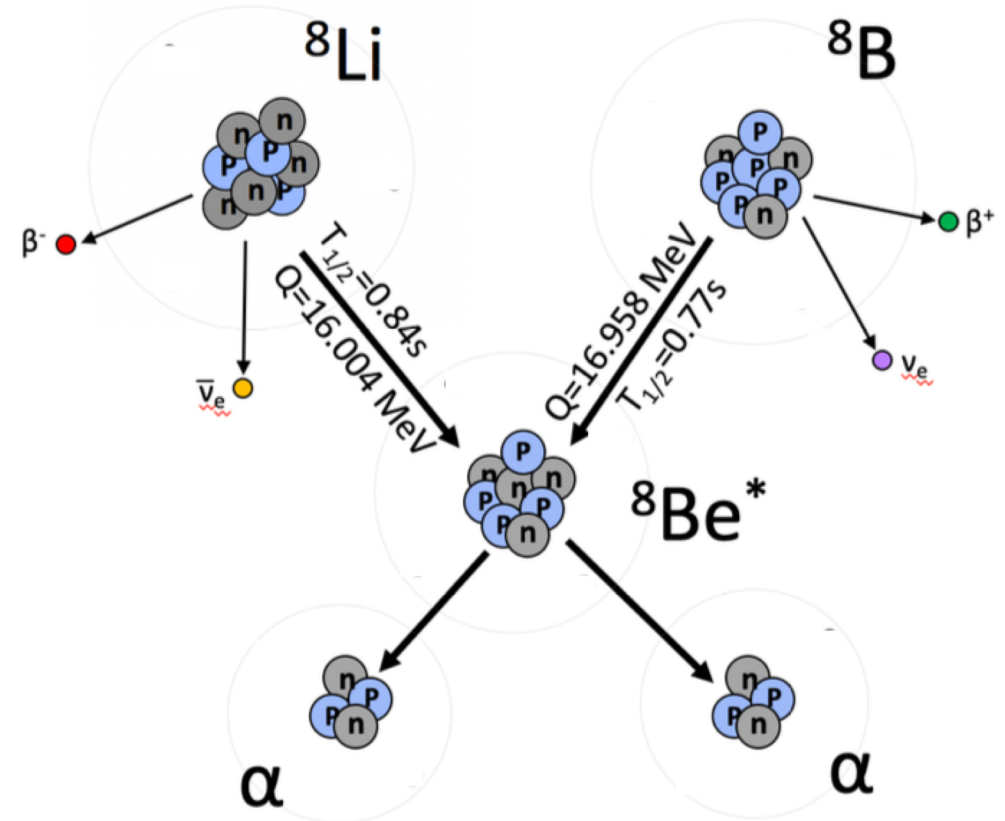
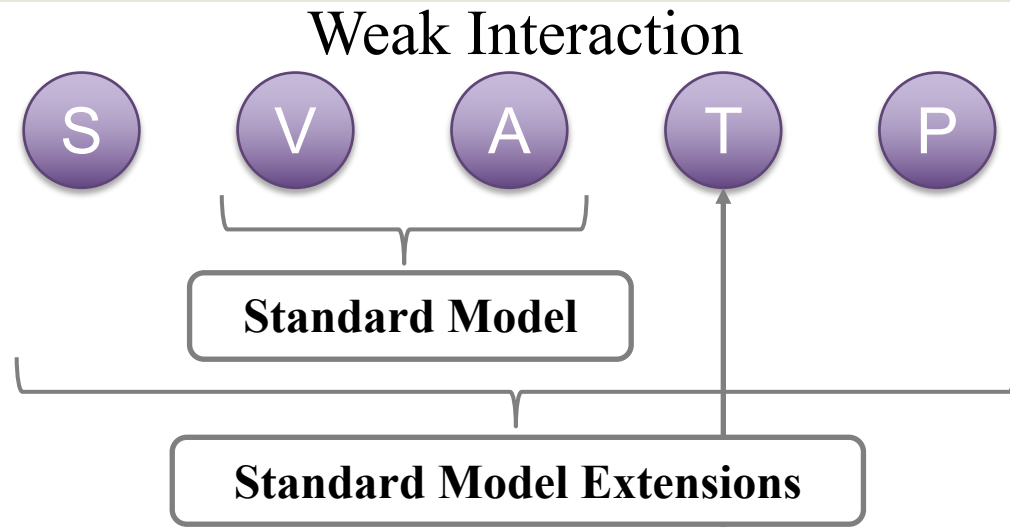
Precision measurements of ^8Li and ^8B beta decay to probe BSM physics



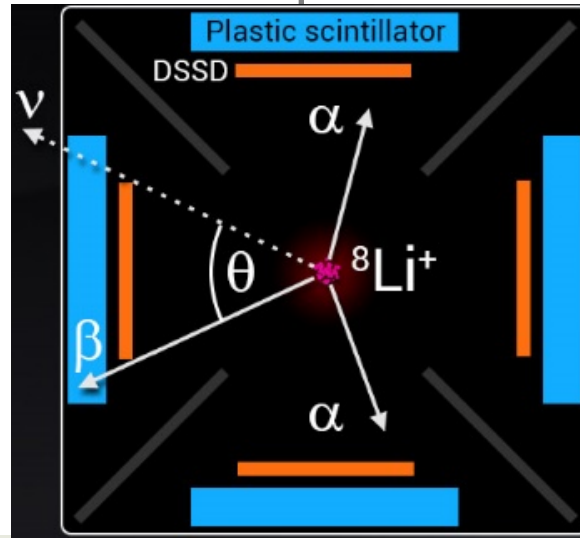
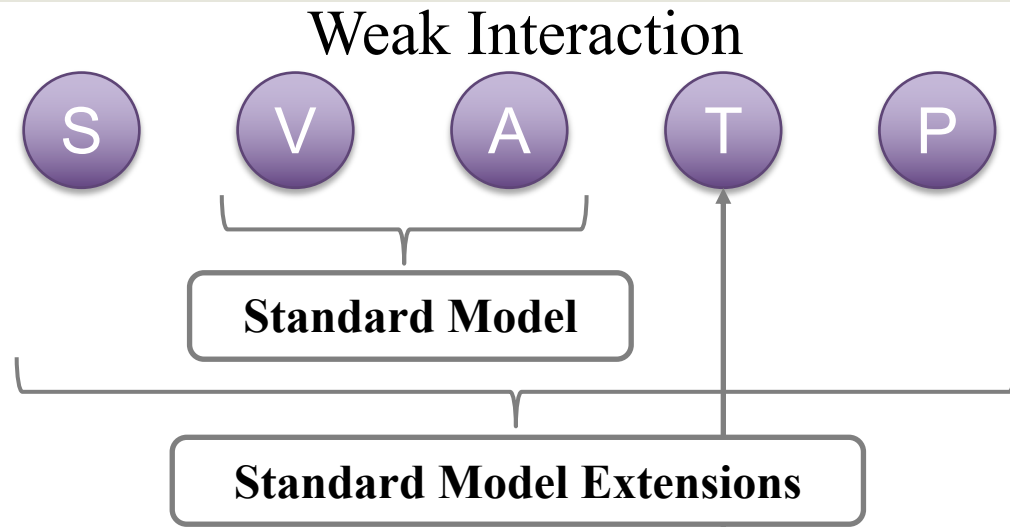
$|C_T/C_A|^2 < 0.011$ (95.5 C.L) with a statistical error of 0.0038 and systematic error of 0.0043,
 Which corresponds to $a_{\beta\nu} = -0.3342 \pm 0.0039$

M. G. Sternberg, R. Segel, N. D. Scielzo, *et al.*, PRL **115**, 182501 (2015).

Precision measurements of ${}^8\text{Li}$ and ${}^8\text{B}$ beta decay to probe BSM physics



Precision measurements of ^8Li and ^8B beta decay to probe BSM physics



Systematic Uncertainty	$\Delta C_T/C_A ^2$
Calibration	1.4×10^{-4}
α energy corrections	1.17×10^{-3}
Cuts to the data	1.25×10^{-3}
Radiative and recoil order terms	3.36×10^{-3}
α Si detector lineshape	6.3×10^{-4}
β Scattering	5.0×10^{-4}
Total	3.62×10^{-3}

From Mary Burkey's PhD Thesis (U. Chicago/ANL/LLNL, 2019)

$$\langle A \rangle = \underbrace{a \times 1}_{\text{Leading order (Gamow-Teller)}} + \underbrace{b \times \frac{q}{M} + c \times \frac{q^2}{M^2} + \dots}_{\text{Recoil-order}}$$

For ^8Li and ^8B beta decays $q/M \sim 0.002$

Experiment needs reliable β -decay recoil-order terms

Largest recoil term matrix elements in impulse approximation

${}^8\text{Be}^*$

$\xleftarrow{\beta\text{-decay}}$

${}^8\text{Li}$

$$j_K \sim \langle \Psi_f || \sum_i^A \tau_i^\pm [\sigma_i \times \hat{Q}_2(\hat{r}_i)]^K || \Psi_0 \rangle$$

Nuclear recoil form-factor

$$c_0 \sim \langle \Psi_f || \sum_i^A \tau_i^\pm \sigma_i || \Psi_0 \rangle$$

Gamow-Teller matrix element

B. R. Holstein, Rev. Mod. Phys. 46, 789 (1974)

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From Mary Burkey's PhD Thesis (U. Chicago/ANL/LLNL, 2019)

Previous values from Sumikama *et al.* measurements [PRC 83, 065501 (2011)]:

$$j_2/A^2 c_0 = -490 \pm 70$$

$$j_3/A^2 c_0 = -980 \pm 280$$

Heroic measurements, but low statistics ➔ “the obtained [recoil] terms were considered as the value averaged over the analyzed energy region”



Experiment needs reliable β -decay recoil-order terms

Largest recoil term matrix elements in impulse approximation

${}^8\text{Be}^*$

β^- -decay

${}^8\text{Li}$

$$j_K \sim \langle \Psi_f || \sum_i^A \tau_i^\pm [\sigma_i \times \hat{Q}_2(\hat{r}_i)]^K || \Psi_0 \rangle$$

Nuclear recoil form-factor

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B. R. Holstein, Rev. Mod. Phys. 46, 789 (1974)

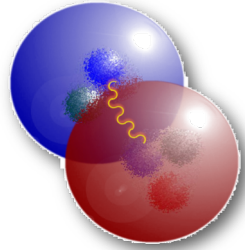
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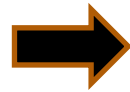
- Need to decrease uncertainty on $j_{2,3}/A^2 c_0$ (and other recoil-order terms)
- *Ab initio* methods to rescue

Ab initio methods in nuclear physics

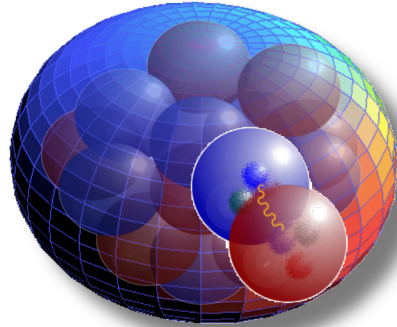
First Principles



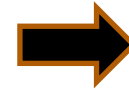
Realistic
Interactions



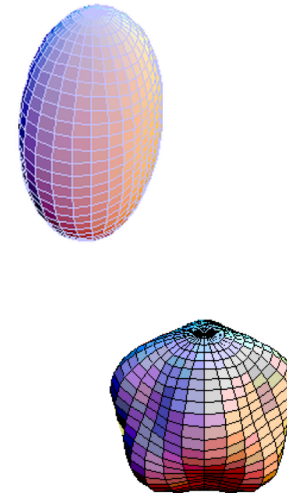
Many-body Dynamics



Symmetry-Adapted
No-Core Shell Model
(SA-NCSM)



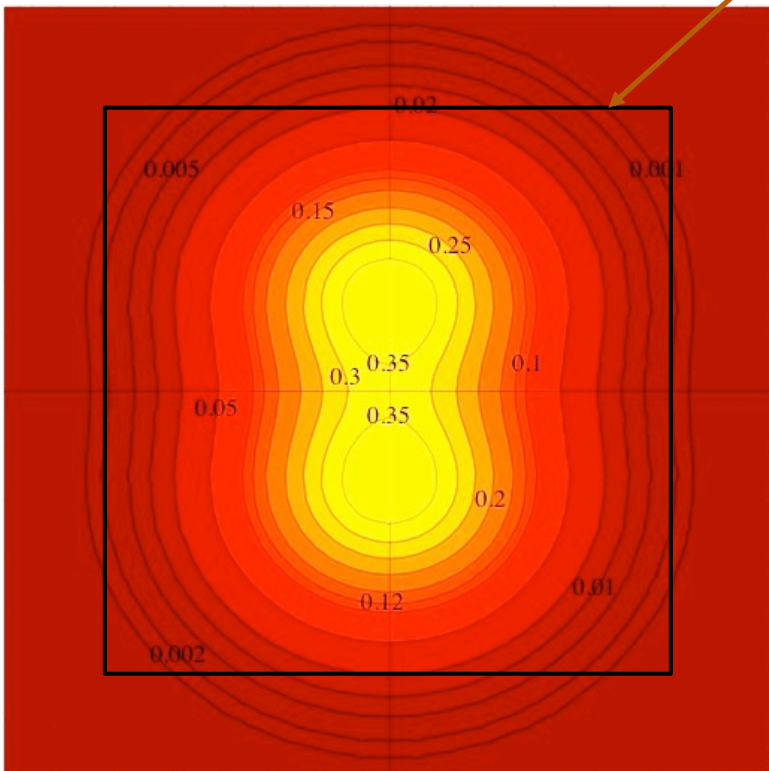
Properties of Nuclei



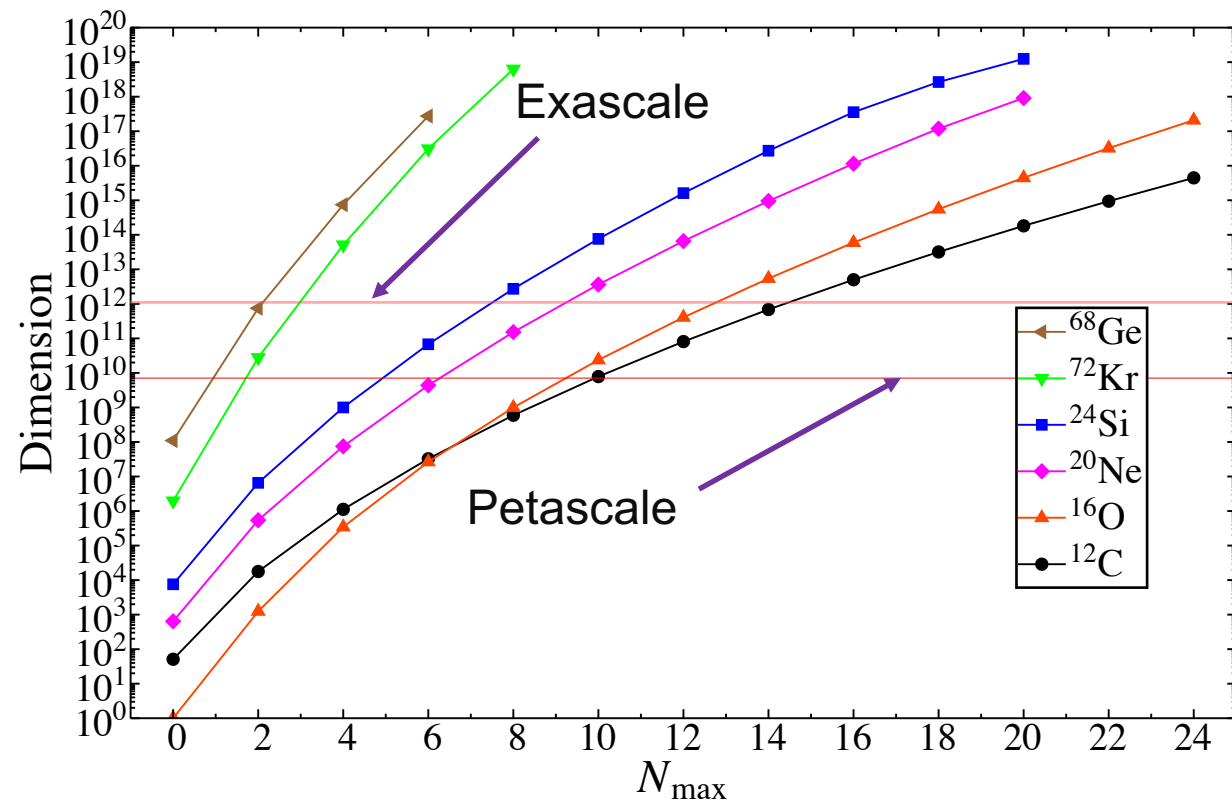
Explosive growth of the model space

Nucleus in
model space

Conventional Shell Model



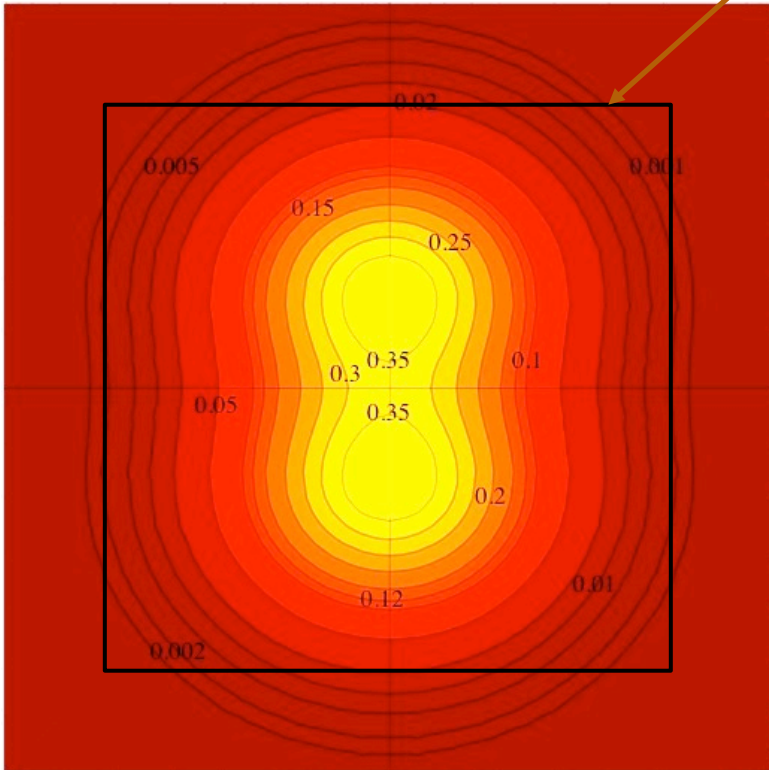
N_{\max}



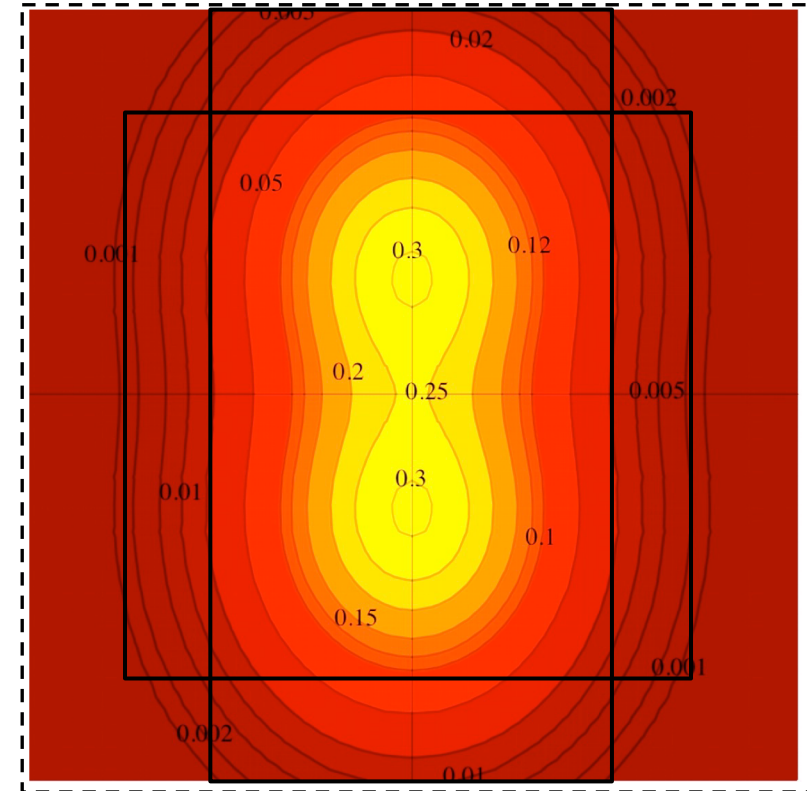
Symmetry-adapted basis helps dramatically reduce the models space size

Nucleus in model space

Conventional Shell Model

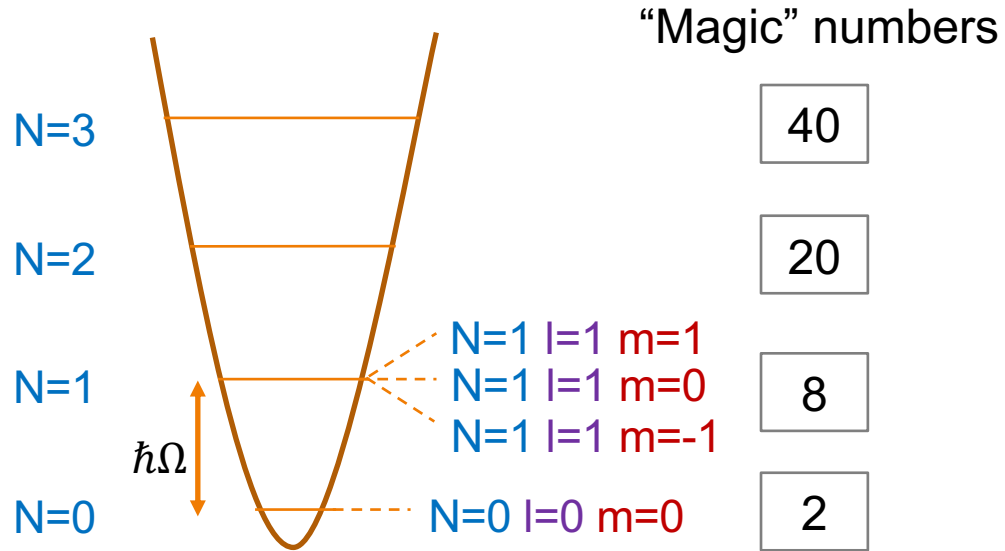


Ab initio Symmetry-adapted No-core Shell Model (SA-NCSM)



SU(3) and symplectic symmetry

Quantum Harmonic Oscillator Basis



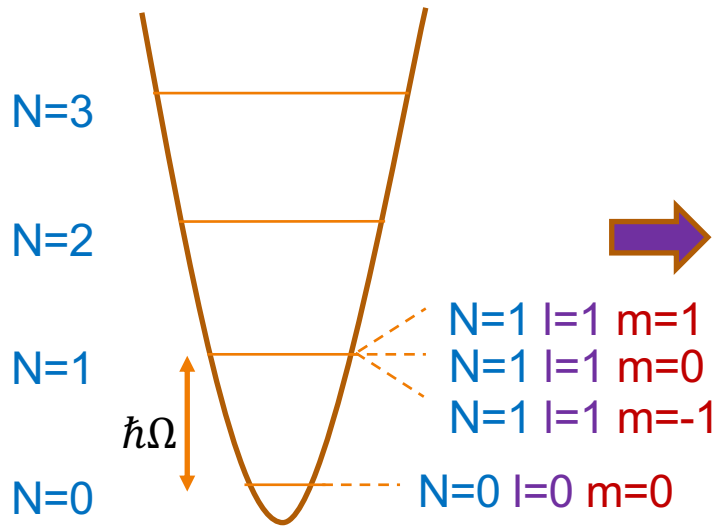
Filling of a harmonic oscillator (HO) shell corresponds to a “magic” number



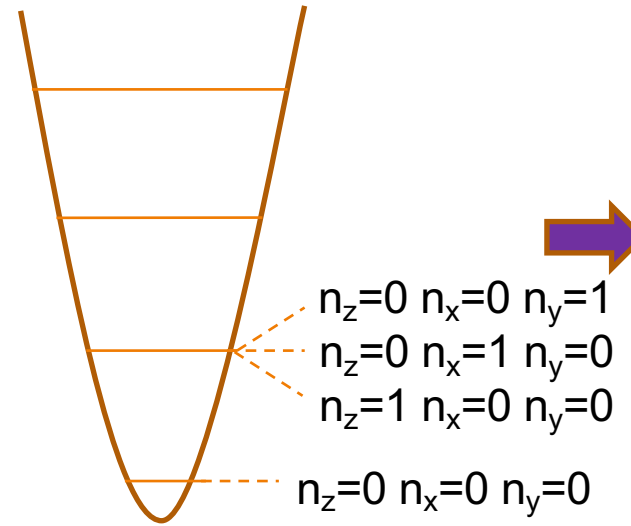
Maria Goeppert Mayer and Hans Jansen
 Nobel Prize in Physics (1963)
 “...for their discoveries concerning nuclear shell structure”



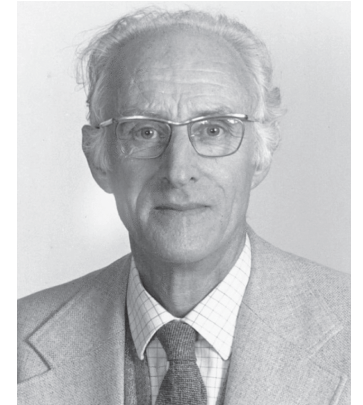
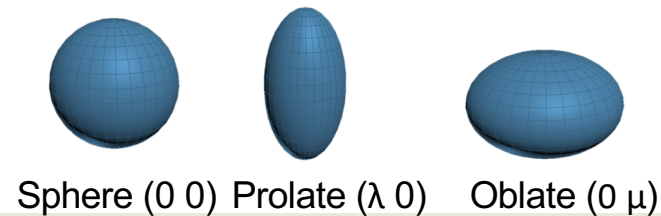
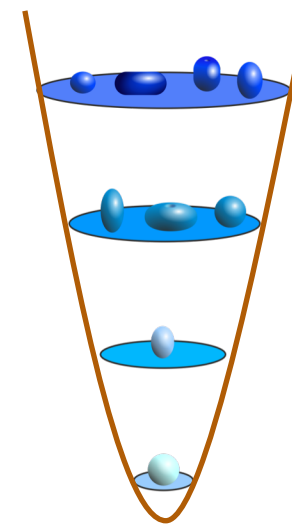
Symmetry-adapted Basis: SU(3)-coupled



Spherical harmonic oscillator (HO): basis states given by $\{N \mid m\}$

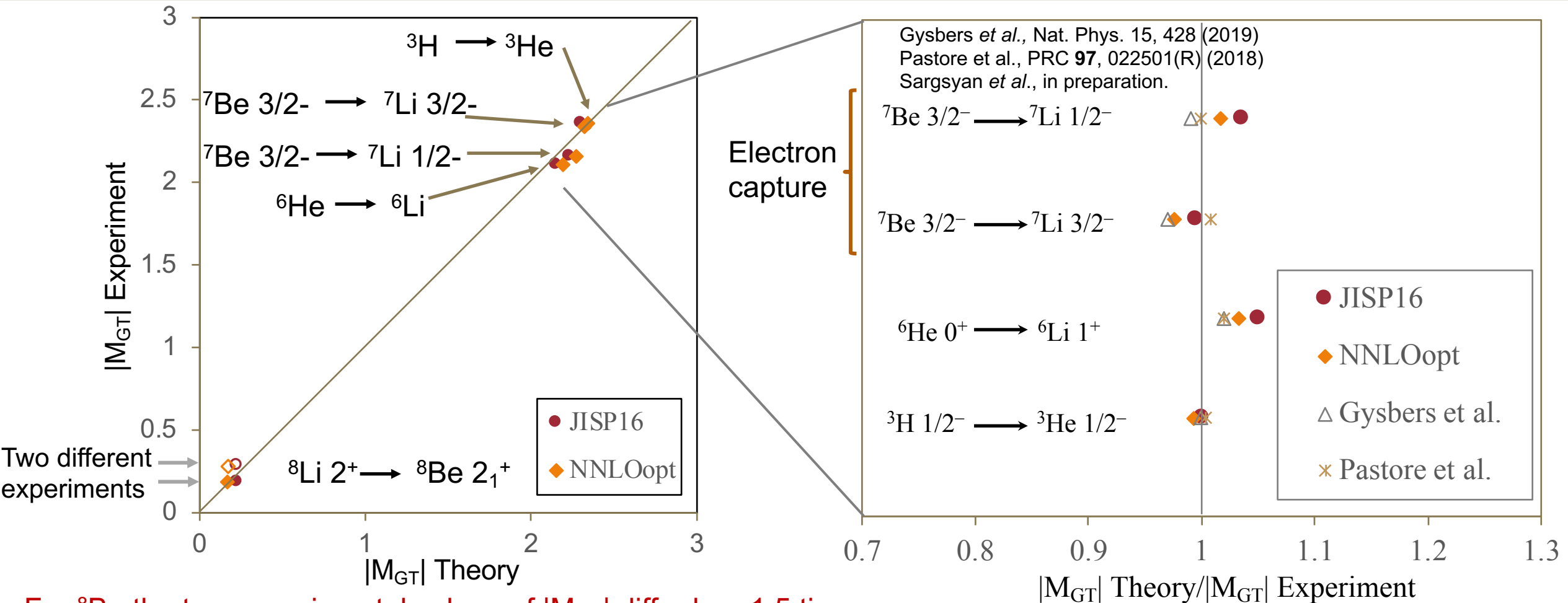


$N = n_z + n_x + n_y$
 Basis states given by $\{n_x \ n_y \ n_z\}$



James Philip Elliott

Beta decays with SA-NCSM

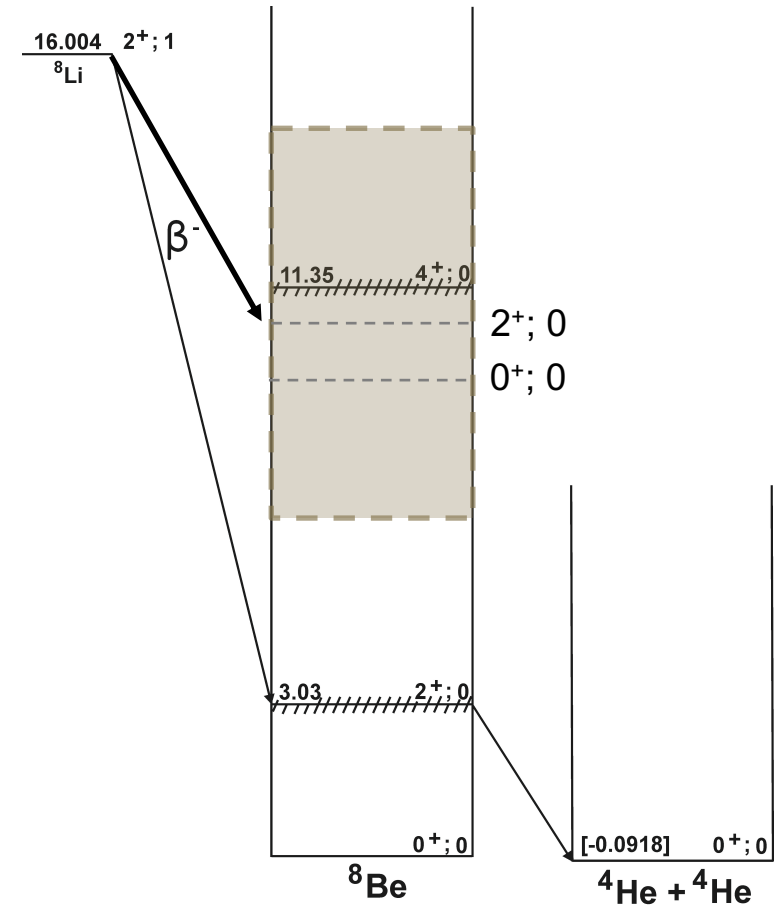
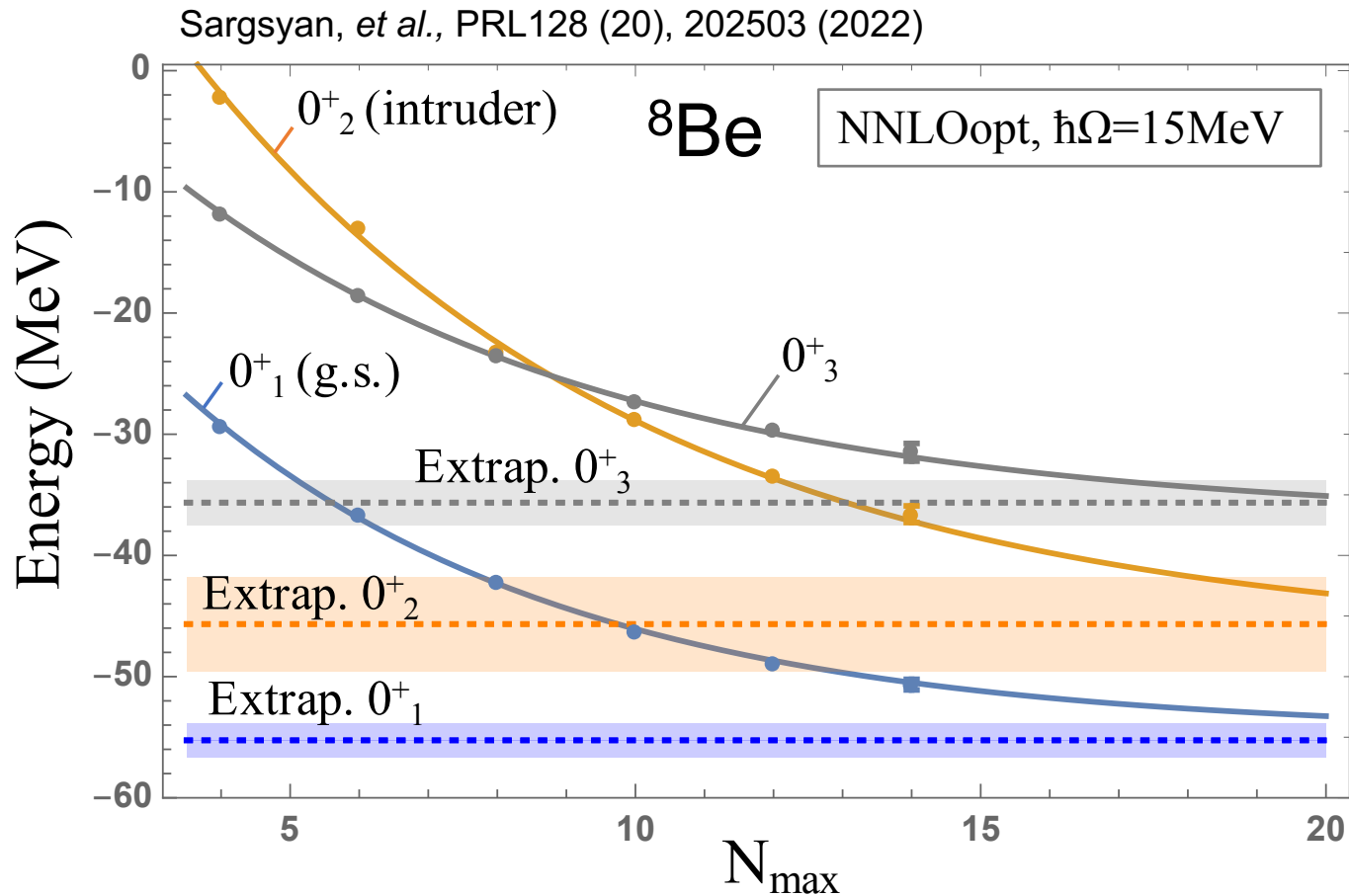


For ${}^8\text{Be}$ the two experimental values of $|M_{GT}|$ differ by ~ 1.5 times

Non-renormalized interactions, unquenched g_A



Possible intruder states in ${}^8\text{Be}$ can explain the discrepancy in ${}^8\text{Li}$ beta decay



Adapted from <https://nucldata.tunl.duke.edu>



0⁺ and 2⁺ intruder states in ⁸Be

PHYSICAL REVIEW C, VOLUME 64, 051301(R)

Intruder states in ⁸Be

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¹Institut de Recherches Subatomiques, IN2P3-CNRS-Université Louis Pasteur, Bâtiment 27/11, F-67037 Strasbourg Cedex 2, France

²Lawrence Livermore National Laboratory, L-414, P. O. Box 808, Livermore, California 94551

³Department of Physics and Astronomy, Iowa State University, Ames, Iowa 50011

(Received 11 July 2001; published 4 October 2001)

Low-lying intruder $T=0$ states in ⁸Be have been posited and challenged. To address this issue, we performed *ab initio* shell model calculations in model spaces consisting of up to $10\hbar\Omega$ excitations above the unperturbed ground state with the basis state dimensions reaching 1.87×10^8 . To gain predictive power we derive and use effective interactions from realistic nucleon-nucleon (NN) potentials in a way that guarantees convergence to the exact solution with increasing model space. Our $0\hbar\Omega$ dominated states show good stability when the model space size increases. At the same time, we observe a rapid drop in excitation energy of the $2\hbar\Omega$ dominated $T=0$ states. In the $10\hbar\Omega$ space the intruder 0^+0 state falls below 18 MeV of excitation and, also, below the lowest 0^+1 state. Our extrapolations suggest that this state may stabilize around 12 MeV. We hypothesize that these states might be the broad resonance intruder states needed in R -matrix analysis of α - α elastic scattering. In addition, we present our predictions for the $A=8$ binding energies with the CD-Bonn NN potential.

Measurement of the full excitation spectrum of the ⁷Li(p, γ) $\alpha\alpha$ reaction at 441 keV

Michael Munch*, Oliver Sølund Kirsebom, Jacobus Andreas Swartz, Karsten Riisager, Hans Otto Uldall Fynbo

Department of Physics and Astronomy, Aarhus University, Denmark



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Radiative decay width
Light nuclei

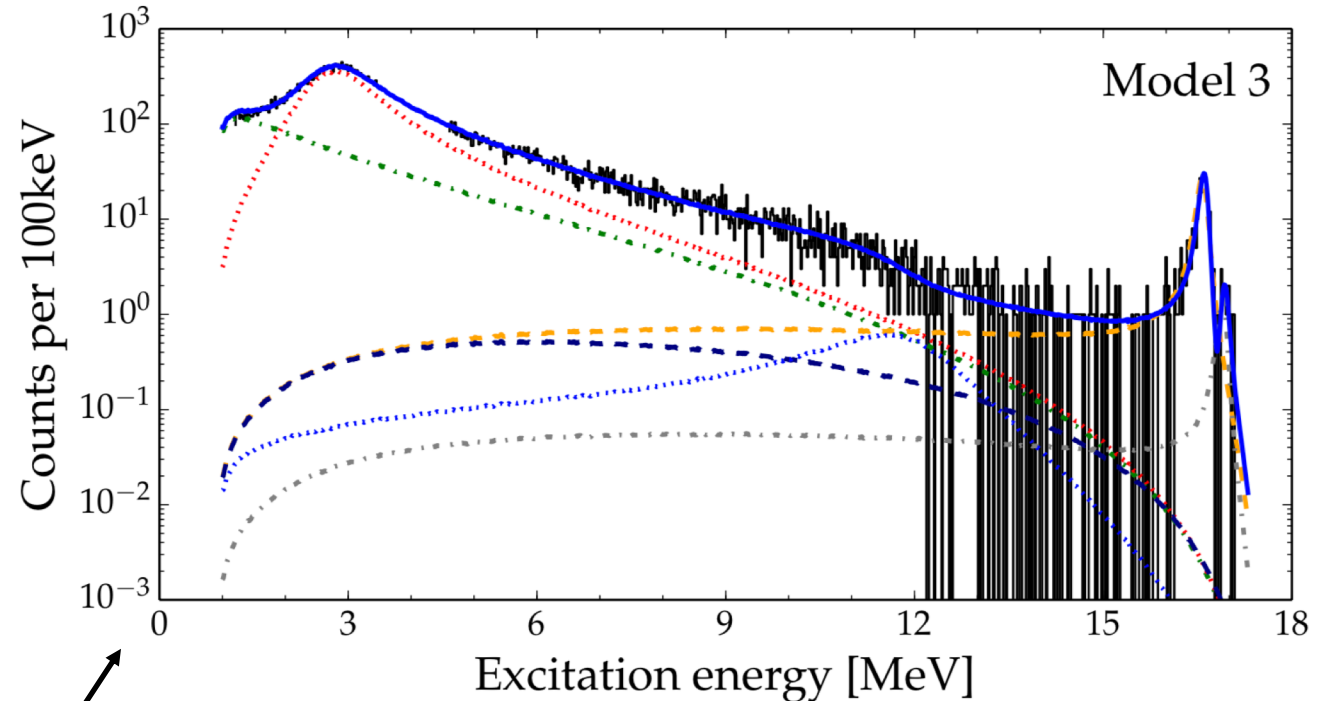
ABSTRACT

A current challenge for *ab initio* calculations is systems that contain large continuum contributions such as ⁸Be. We report on new measurements of radiative decay widths in this nucleus that test recent Green's function Monte Carlo calculations.

Traditionally, γ ray detectors have been utilized to measure the high energy photons from the ⁷Li(p, γ) $\alpha\alpha$ reaction. However, due to the complicated response function of these detectors it has not yet been possible to extract the full γ ray spectrum from this reaction. Here we present an alternative measurement using large area Silicon detectors to detect the two α particles, which provides a practically background free spectrum and retains good energy resolution.

The resulting spectrum is analyzed using a many-level multi channel R-matrix parametrization. Improved values for the radiative widths are extracted from the R-matrix fit. We find evidence for significant non-resonant continuum contributions and **tentative evidence for a broad 0^+ resonance at 12 MeV.**

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- - - 0₁⁺ - - - 2₁⁺ - - - 2₃⁺ — Total
- - - 0₂⁺ - - - 2₂⁺ - - - 2₄⁺

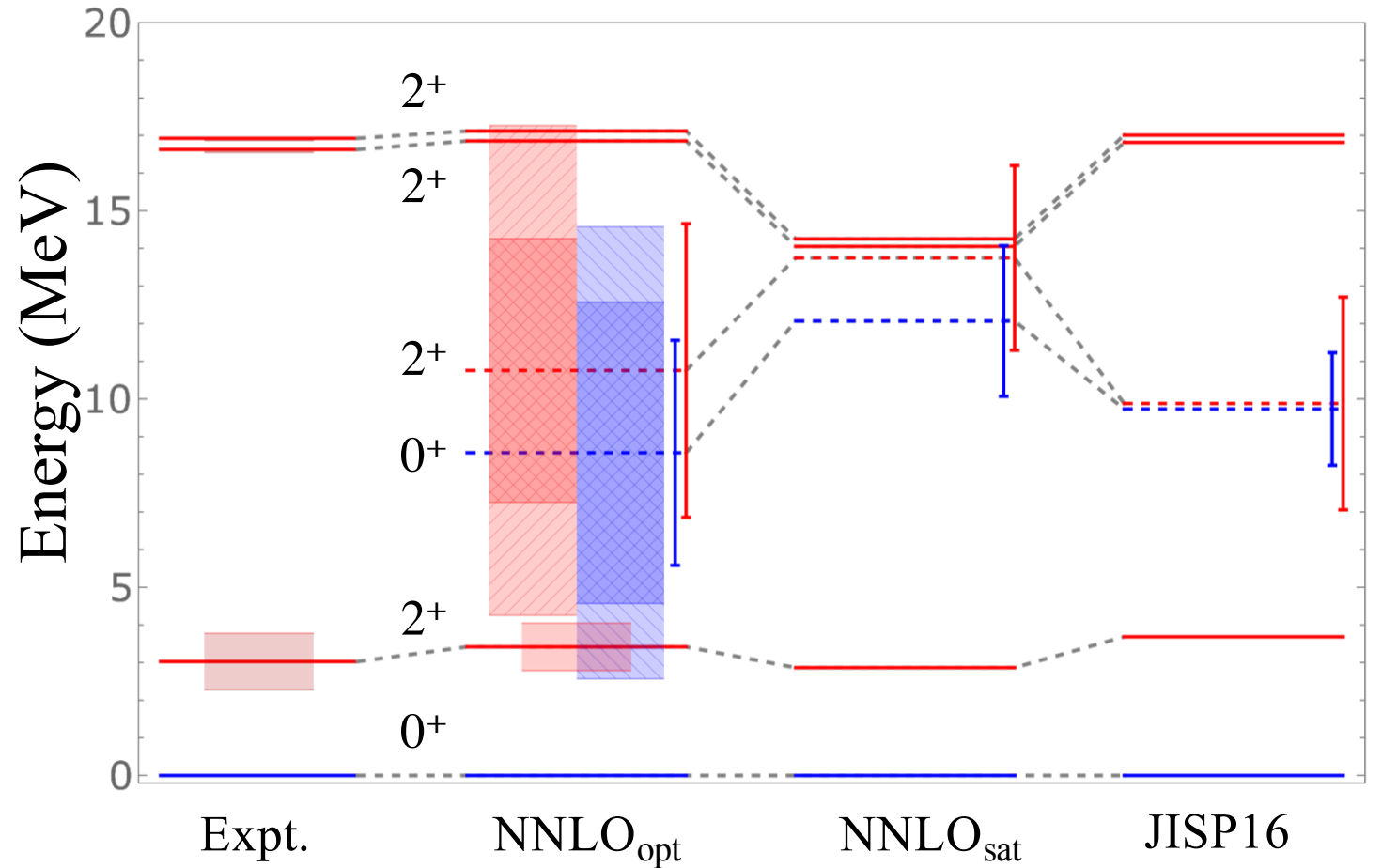
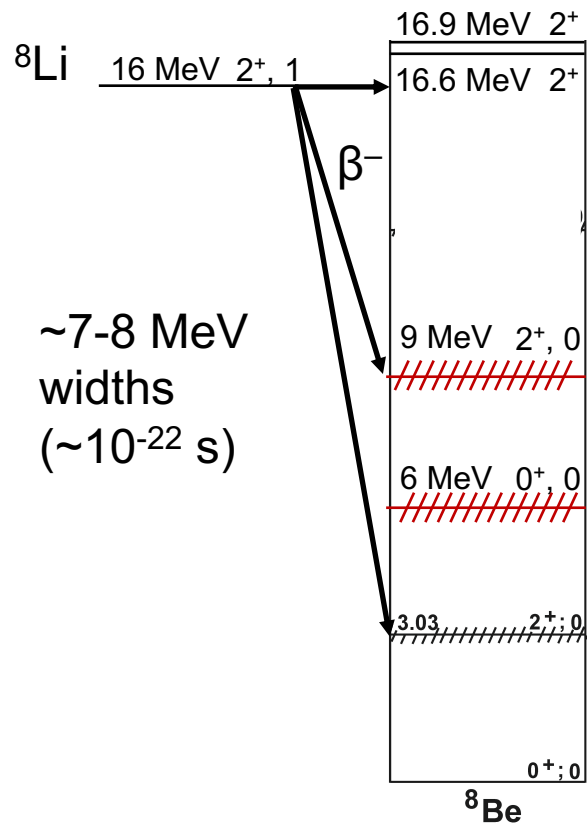
Munch et al., Phys. Lett. B 782 (2018) 779–784



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0⁺ and 2⁺ intruder states in ⁸Be

➤ First proposed by Barker in 1968-69



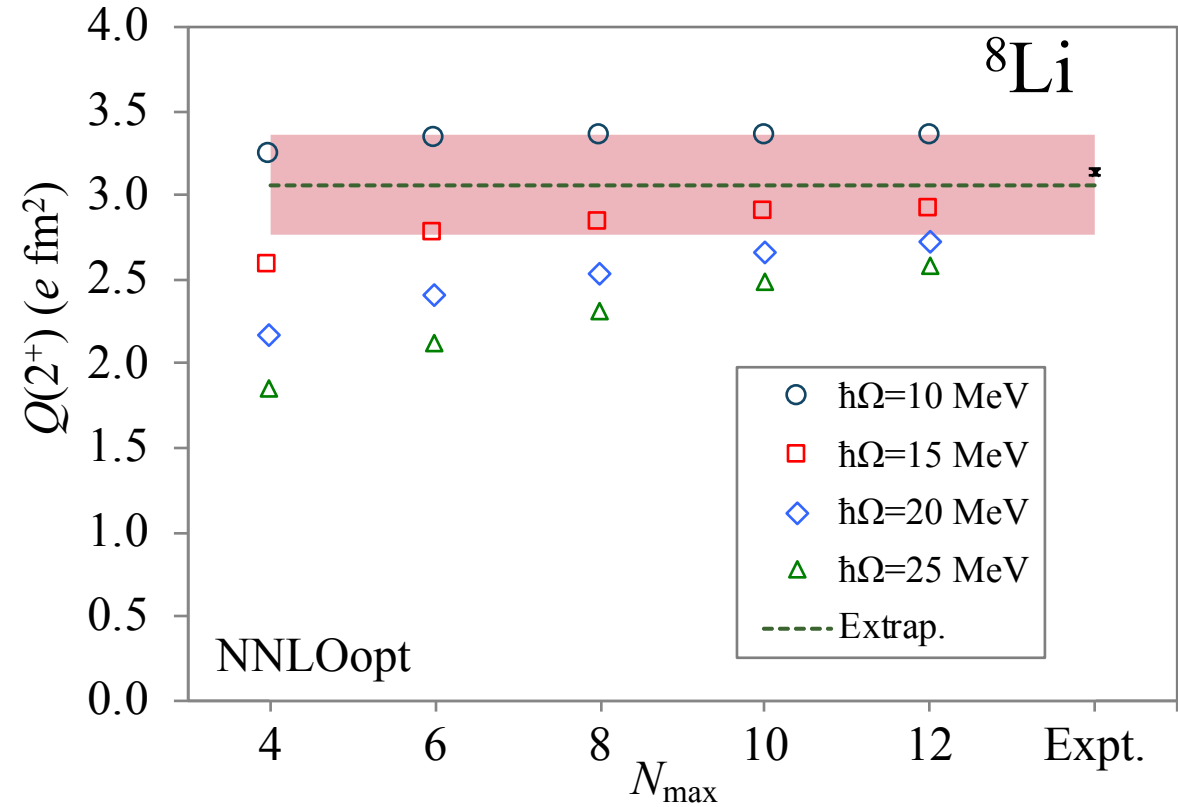
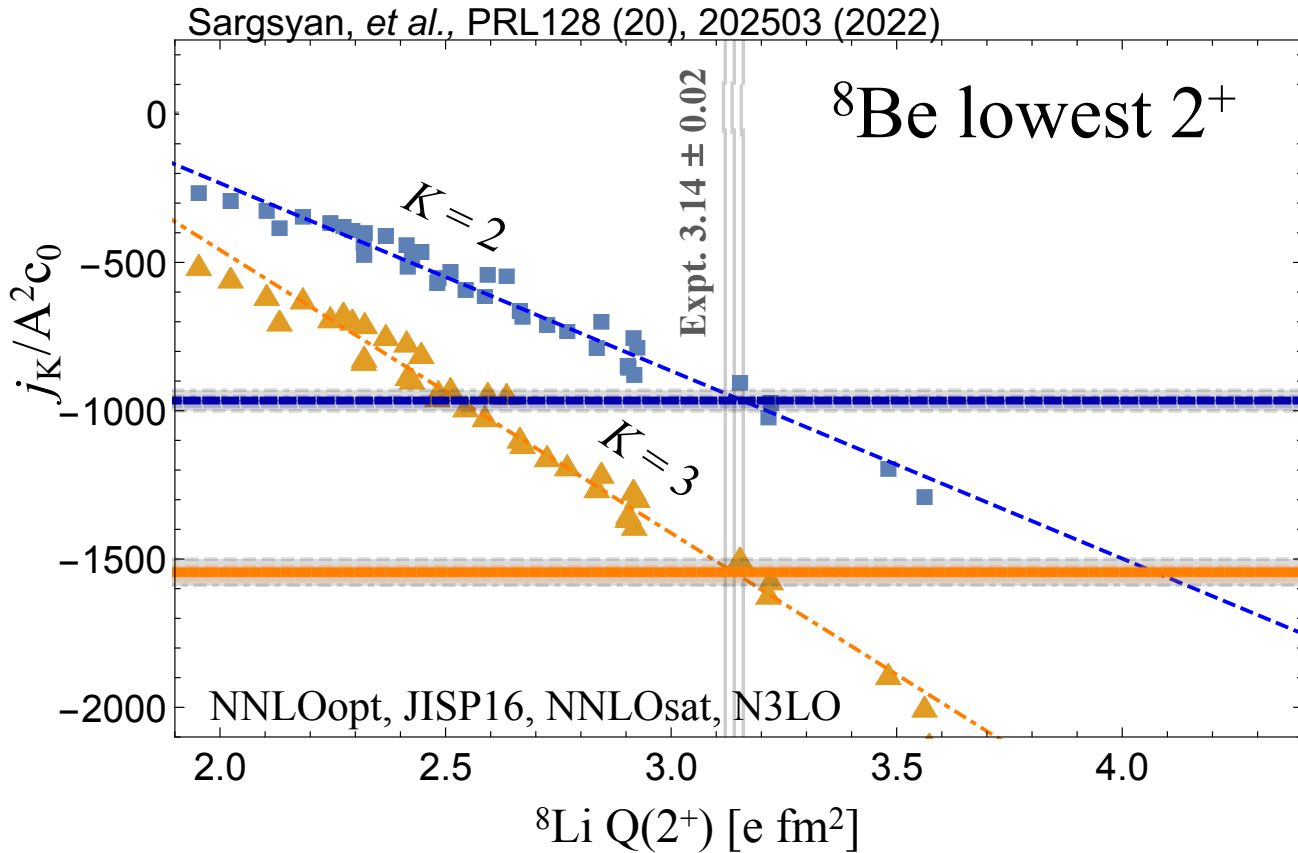
F. C. Barker. Australian Journal of Physics, vol. 21, 239–257, 1968.

F. C. Barker. Australian Journal of Physics, vol. 22, 293–316, 1969.

Sargsyan, *et al.*, PRL128 (20), 202503 (2022)



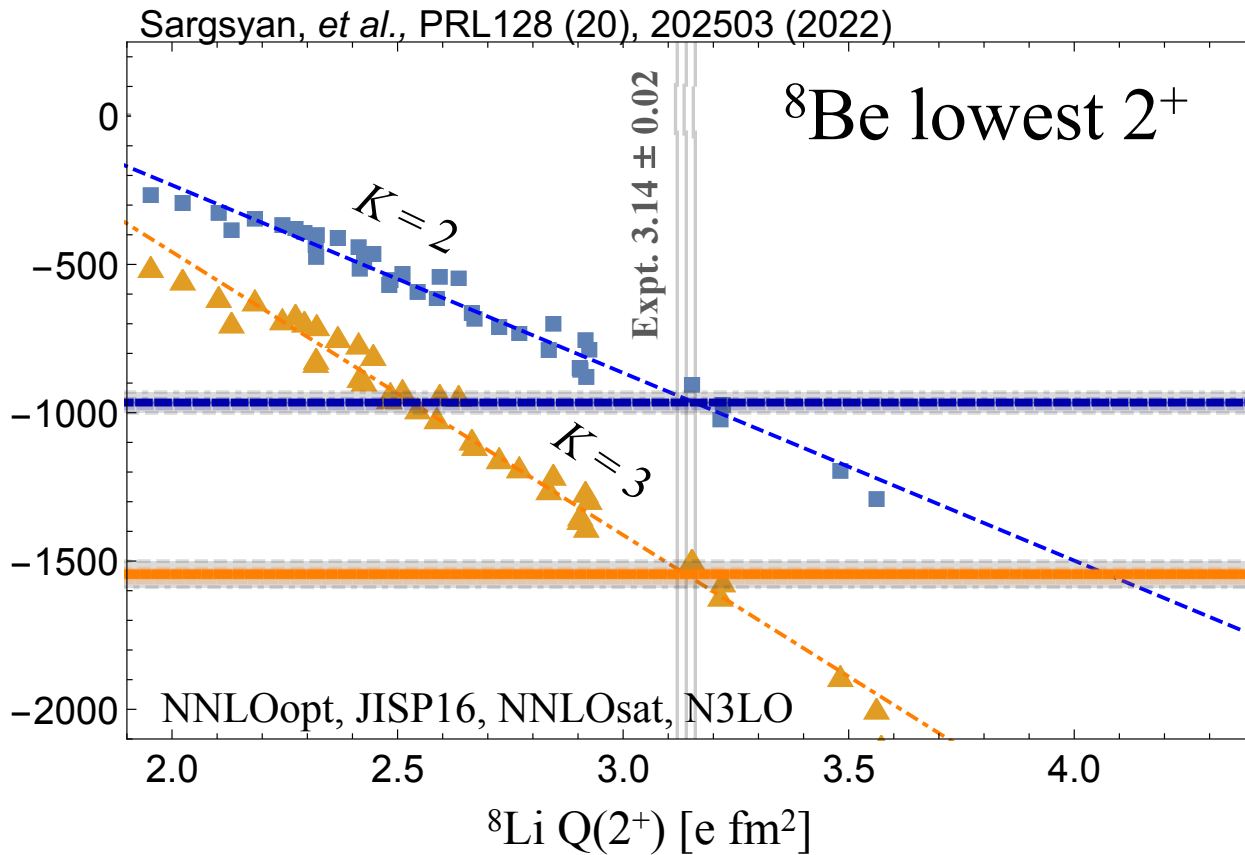
Correlation between j_K and Q helps constrain recoil order terms



$$j_K \propto \langle \Psi_f || \sum_i^A \tau_i^\pm [\sigma_i \times \hat{Q}_2(\hat{r}_i)]^K || \Psi_0 \rangle \quad \longleftarrow \text{Nuclear recoil term}$$

Most precise beta-decay measurement of its type in 50 years!

Mary Burkey's PhD Thesis (U. Chicago/ANL/LLNL, 2019)



Improved by nearly 50%!

Systematic Uncertainty	$\Delta C_T/C_A ^2$
Calibration	1.4×10^{-4}
α energy corrections	1.17×10^{-3}
Cuts to the data	1.25×10^{-3}
Radiative and recoil order terms	3.36×10^{-3}
α Si detector lineshape	6.3×10^{-4}
β Scattering	5.0×10^{-4}
Total	3.62×10^{-3}

TABLE I. Summary of dominant systematic uncertainties, listed at 1σ .

	Systematic Uncertainty	$\Delta C_T/C_A ^2$
Theory	Intruder State (added linearly)	0.0005
	Recoil-Order Terms & Radiative Corrections	0.0015
Experiment	α -Energy Calibration	0.0007
	Detector Lineshape	0.0009
	Data Cuts	0.0009
	β Scattering	0.0010
Total		0.0028

MT Burkey, *et al.* (incl. Sargsyan), PRL 128 (20), 202502 (2022).

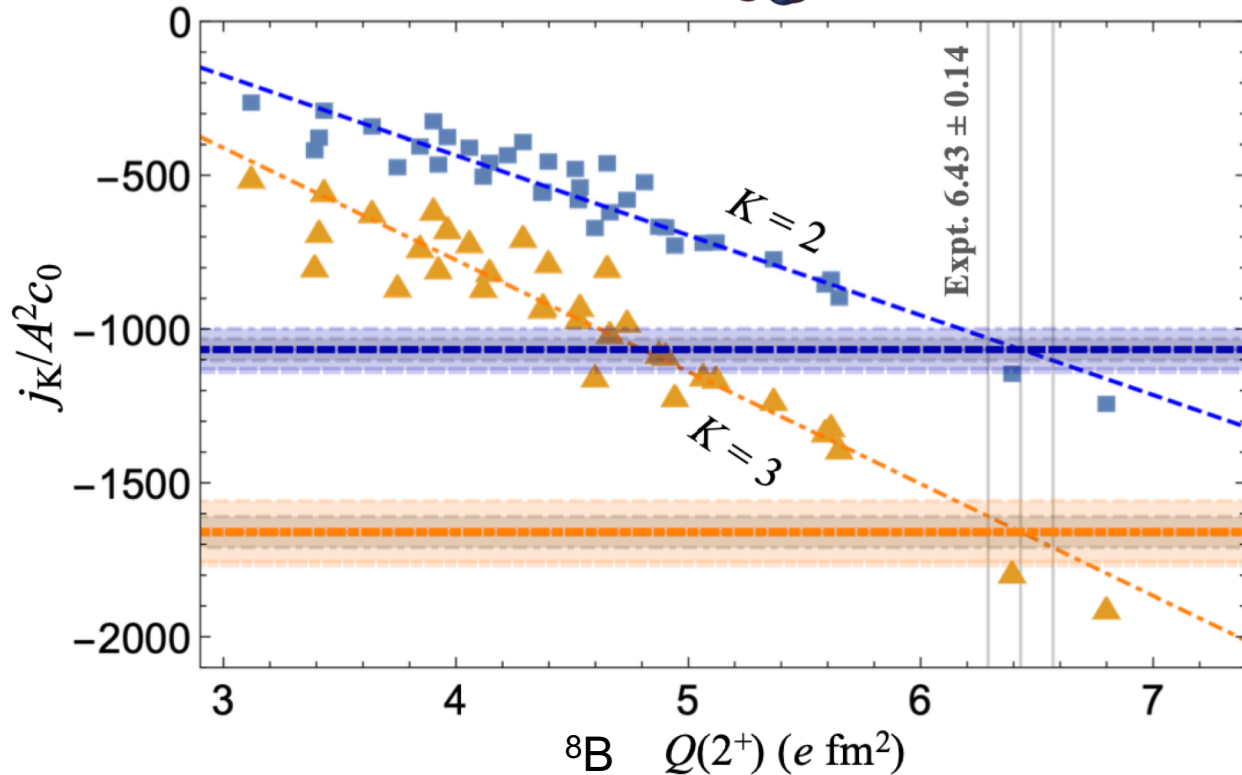
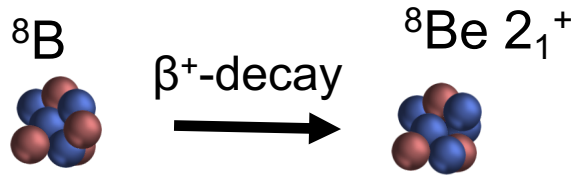
The most precise constraint now

$$|C_T/C_A|^2 = -0.0012 \pm 0.0019_{\text{stat}} \pm 0.0028_{\text{syst}}$$

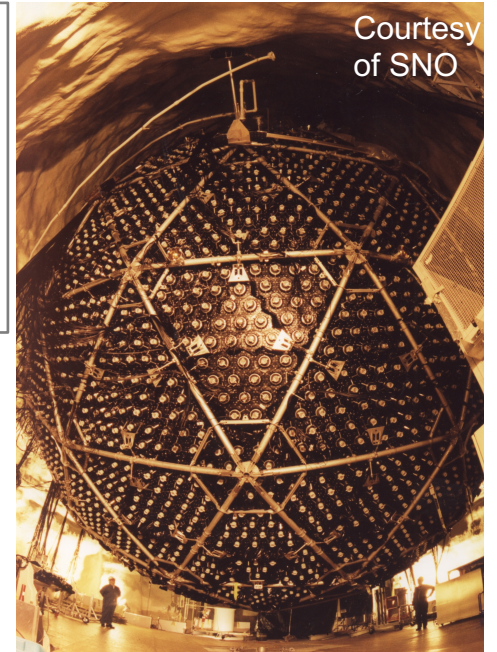
$$a_{\beta\nu} = -0.3325 \pm 0.0023$$



Recoil terms for ${}^8\text{B}$ to inform precision beta decay experiments



Also important for precision measurements of solar neutrinos!



$$|C_T/C_A|^2 = -0.0017 \pm 0.0029_{\text{stat}} \pm 0.0031_{\text{syst}}$$

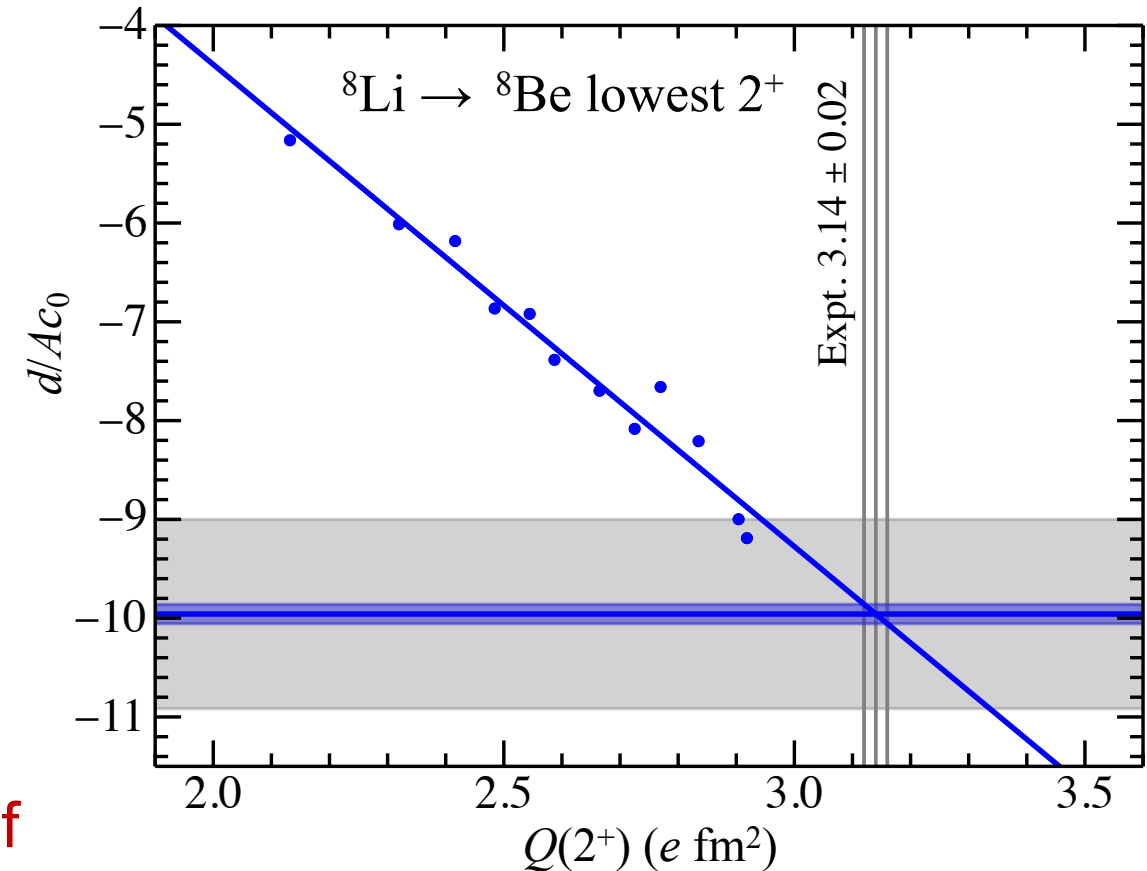
$$a_{\beta\nu} = -0.3345 \pm 0.0028$$

Longfellow, Gallant, Sargsyan, *et al.* PRL **132**, 142502 (2024).

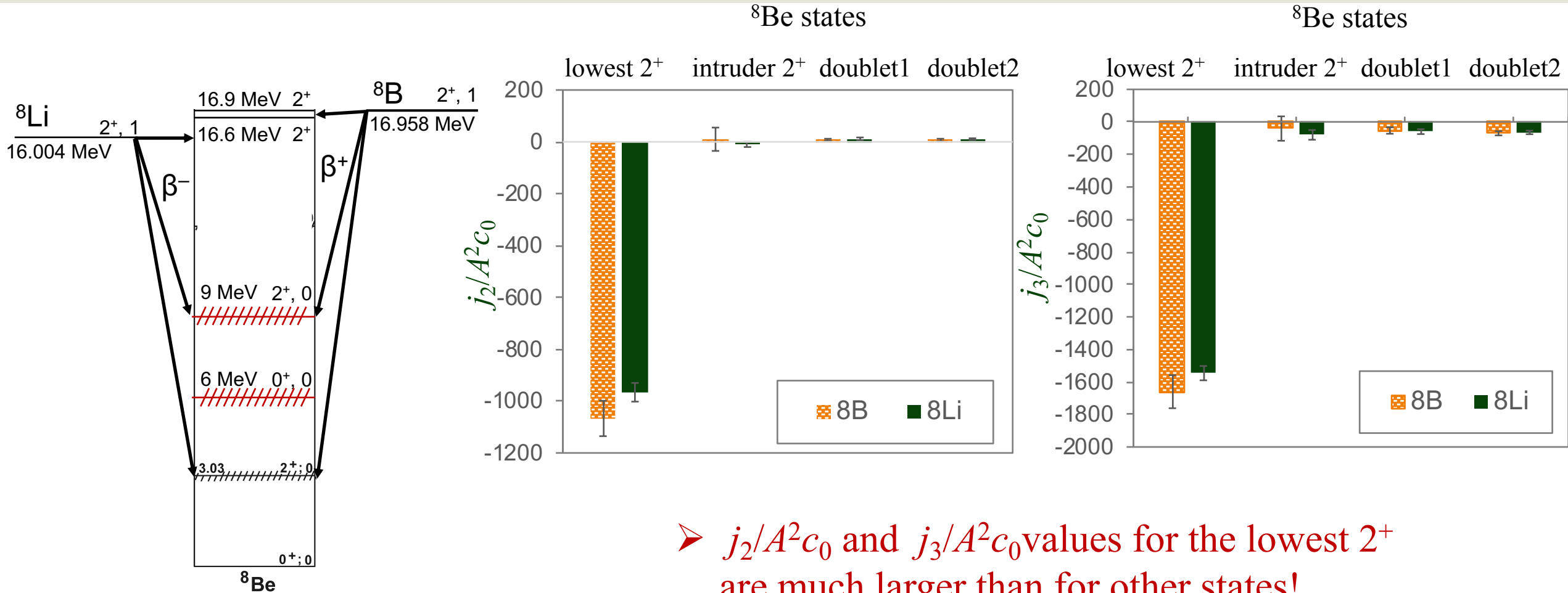


Weak magnetism and induced tensor recoil-order terms

- Weak magnetism (b) and induced tensor (d) recoil terms: next significant after j_2 and j_3
- Important for the tests of conserved vector current (CVC) hypothesis and existence of second class currents
- With SA-NCSM we can calculate these beta decay recoil-order terms for up to intermediate mass nuclei
- We are currently working on calculations of b and d for $^{22}\text{Na} \rightleftharpoons ^{22}\text{Ne}$ beta decay



Recoil terms for all ${}^8\text{Li}$ and ${}^8\text{B}$ β -decay accessible states



➤ j_2/A^2c_0 and j_3/A^2c_0 values for the lowest 2^+ are much larger than for other states!

04-2014

Strong correlation between j_2 and j_3

PHYSICAL REVIEW C **83**, 065501 (2011)

Test of the conserved vector current hypothesis by a β -ray angular distribution measurement in the mass-8 system

T. Sumikama,^{1,2} K. Matsuta,¹ T. Nagatomo,³ M. Ogura,¹ T. Iwakoshi,¹ Y. Nakashima,¹ H. Fujiwara,¹ M. Fu M. Mihara,¹ K. Minamisono,⁴ T. Yamaguchi,⁵ and T. Minamisono⁶

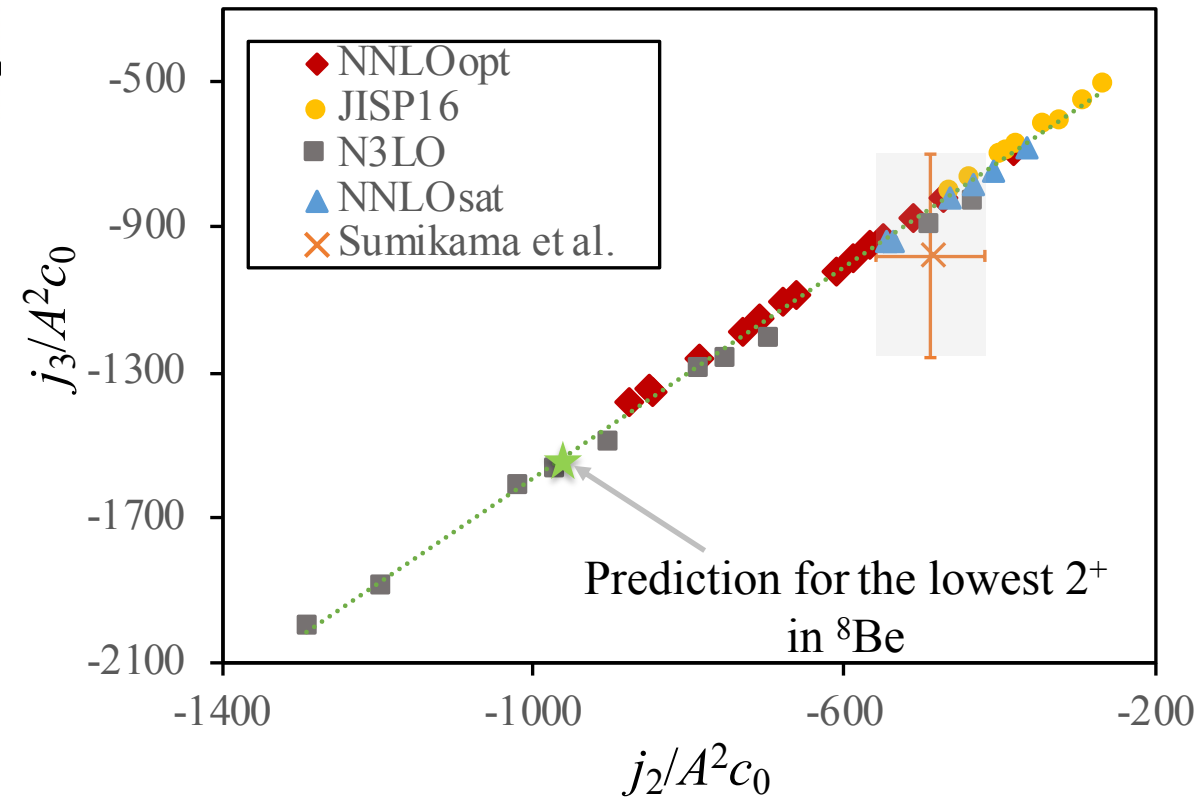
$$j_2/A^2c_0 = -490 \pm 70$$

$$j_3/A^2c_0 = -980 \pm 280$$

Extremely difficult measurements, hence “the obtained [recoil] terms were considered as the value averaged over the analyzed energy region”

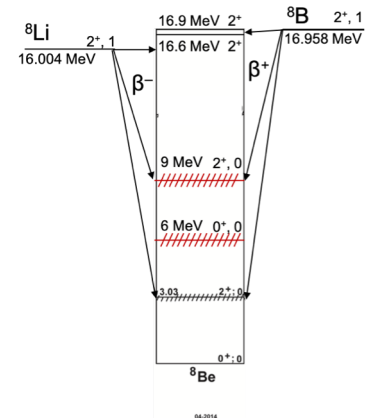
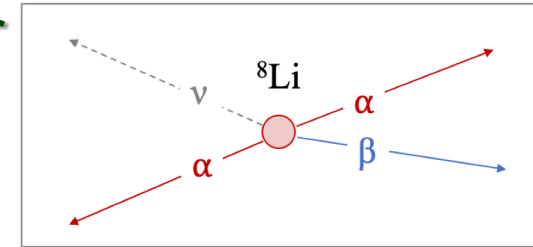
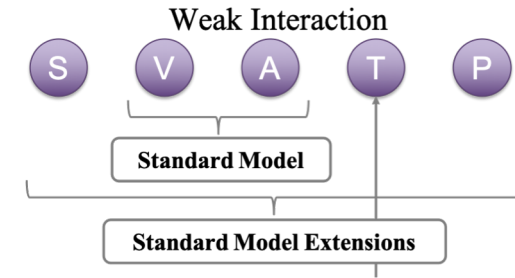
But energy dependence is important!!

Strong correlation between j_2 and j_3 recoil-order terms



Summary

- Precision measurements of ${}^8\text{Li}$ and ${}^8\text{B}$ beta decays have recently set the most stringent limit on the tensor coupling in weak interactions in the low-energy regime
- *Ab initio* SA-NCSM calculations of ${}^8\text{Li}$ and ${}^8\text{B}$ beta decay recoil-order terms helped experiment to decrease systematic uncertainties
- The calculated b/Ac_0 and d/Ac_0 values are important for tests of conserved vector current hypothesis
- Low-lying intruder states in ${}^8\text{Be}$ can have important implications for $A=8$ beta decays and related precision measurements



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Guy Savard



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Daniel Langr

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Czech Academy of Sciences



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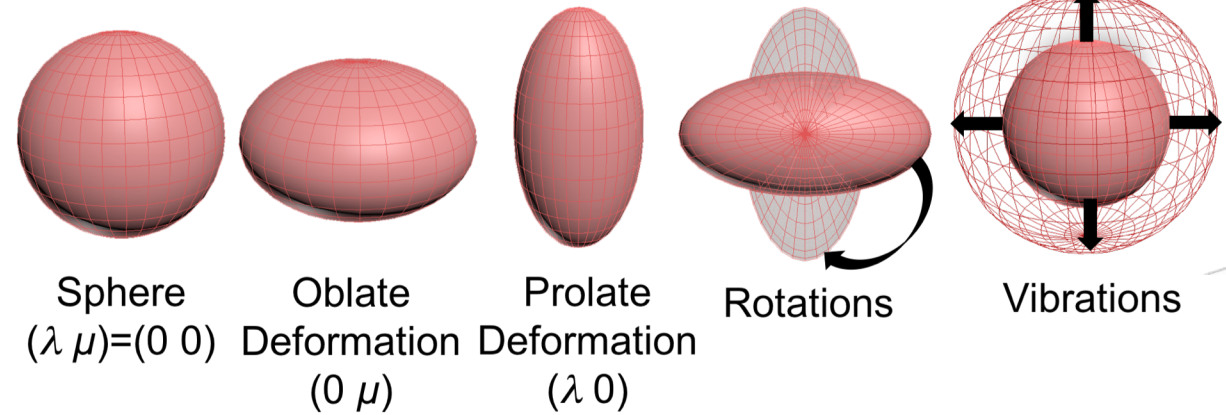
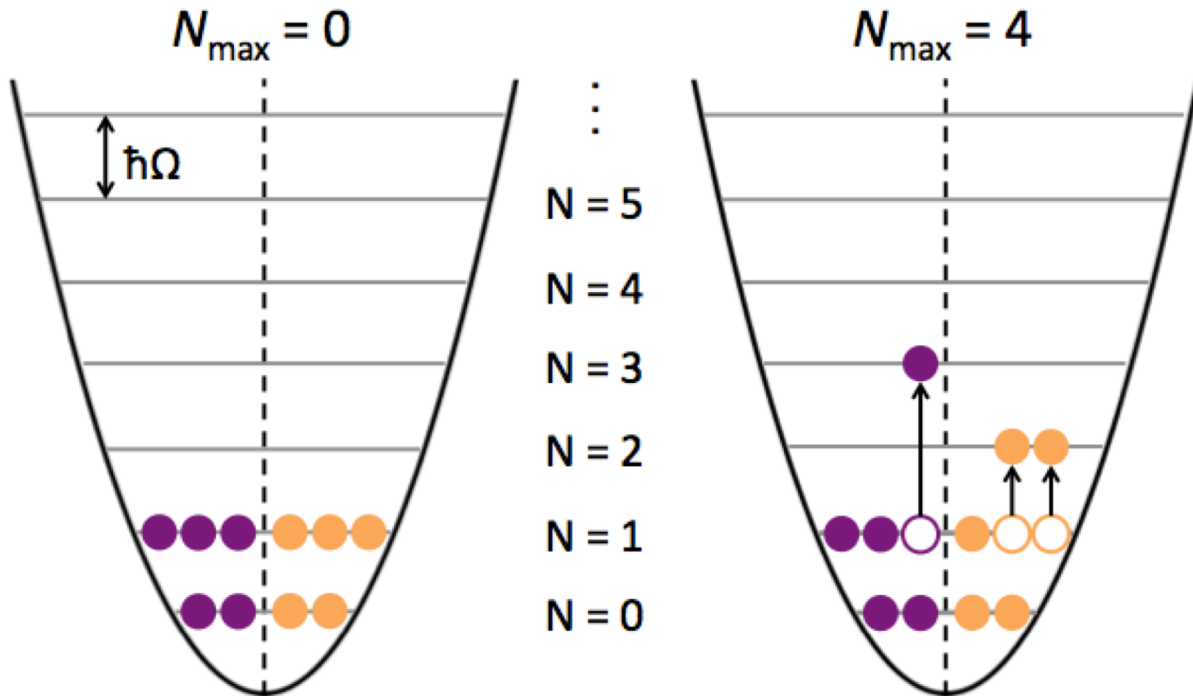
Experimental crew at the ATLAS facility at Argonne National Lab

Thank you!

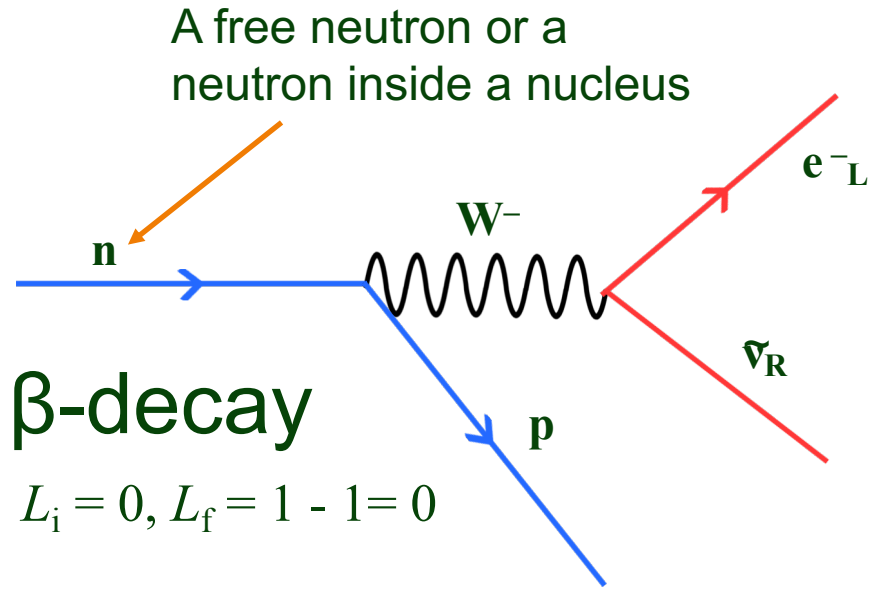


Facility for Rare Isotope Beams
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Backup slide zone



Beta decay as a probe for BSM studies

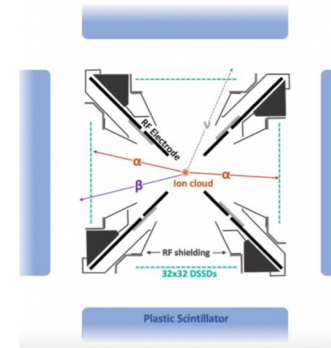


- Unitarity of the CKM quark mixing matrix

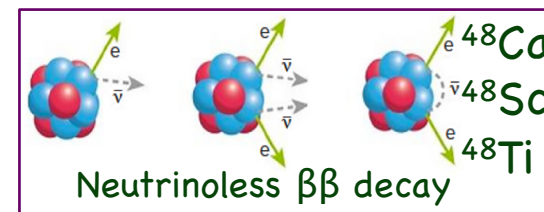
$$\begin{pmatrix} d_w \\ s_w \\ b_w \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Weak states **CKM mixing matrix** Mass eigenstates

- Beyond Standard Model (BSM) terms in weak interaction

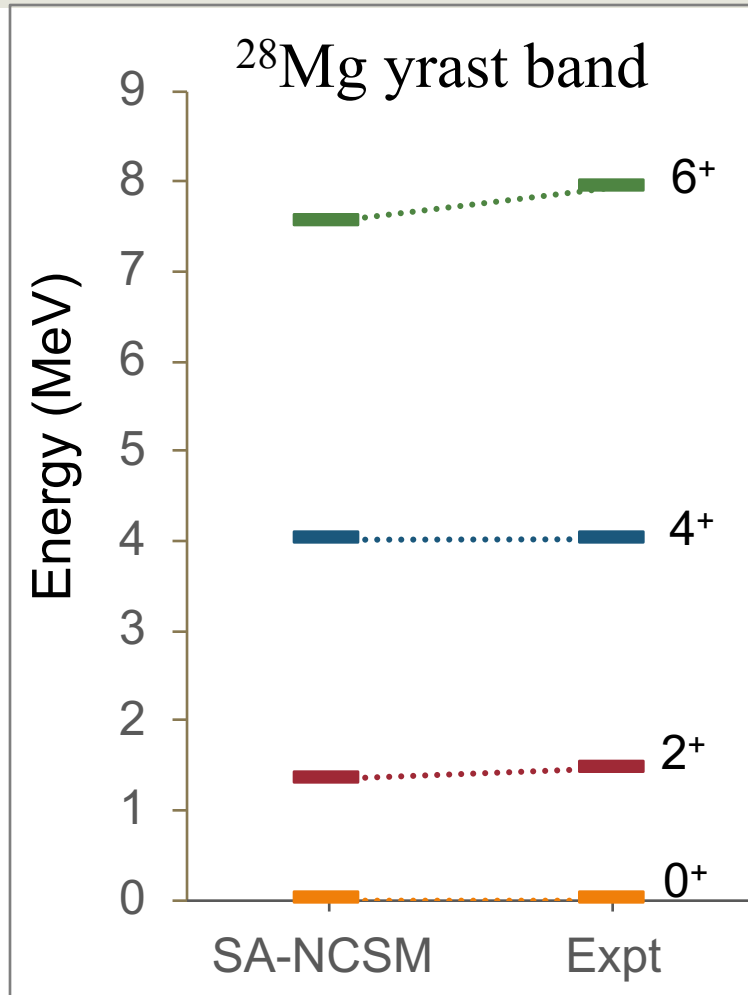


- Neutrinoless double beta decay

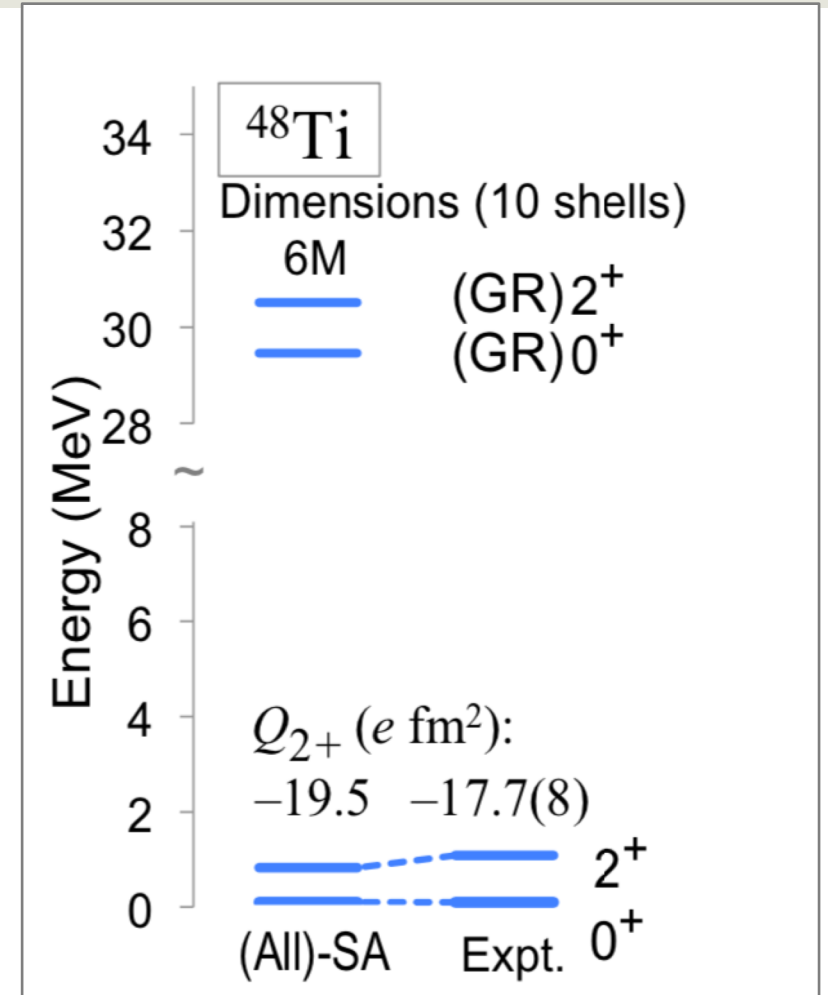


Precision measurements need input from nuclear theory

SA-NCSM can reach intermediate mass nuclei



NNLO_{opt} $\hbar\Omega = 15$ MeV



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Recoil-order terms in β -decay

Beta decay rate: $d\Gamma \propto |T|^2$

T matrix in SM (V-A): $T \propto l^\mu \langle \beta | V_\mu - A_\mu | \alpha \rangle$

Axial current matrix element

Leading order (Gamow-Teller) Recoil-order q/M

$$\begin{aligned}
 l^\mu \langle \beta, J' M' | A_\mu | \alpha, J M \rangle &= C_{J' M' 1 k}^{J M} \epsilon_{ijk} \epsilon_{ij\lambda\eta} \frac{1}{4M} \left[c(q^2) l^\lambda P^\eta - d(q^2) l^\lambda q^\eta \right] \\
 &+ \frac{1}{(2M)^2} h(q^2) q^\lambda P^\eta \mathbf{q} \cdot \mathbf{l} \\
 &+ C_{J' M' 2 k}^{J M} C_{1 n 2 n'}^{2 k} l_n (4\pi/5)^{1/2} Y_{2 n'}(\hat{q}) \frac{q^2}{(2M)^2} j_2(q^2) \\
 &+ C_{J' M' 3 k}^{J M} C_{1 n 2 n'}^{3 k} l_n (4\pi/5)^{1/2} Y_{2 n'}(\hat{q}) \frac{q^2}{(2M)^2} j_3(q^2) + \dots
 \end{aligned}$$

Lepton current matrix element

Recoil-order $(q/M)^2$

For ${}^8\text{Li}$ and ${}^8\text{B}$ beta decay $q/M \sim 0.002$

Systematic Uncertainty	$\Delta C_T/C_A ^2$
Calibration	1.4×10^{-4}
α energy corrections	1.17×10^{-3}
Cuts to the data	1.25×10^{-3}
Radiative and recoil order terms	3.36×10^{-3}
α Si detector lineshape	6.3×10^{-4}
β Scattering	5.0×10^{-4}
Total	3.62×10^{-3}

From Mary Burkey's PhD Thesis (U. Chicago/ANL/LLNL, 2019)