

# NEW RESULTS FROM THE HALL C J/ $\Psi$ -007 EXPERIMENT AND FUTURE PROSPECTS

## MEASUREMENTS TOWARDS THE GLUONIC GRAVITATIONAL FORM FACTORS

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With special thanks to:  
**Burcu Duran, Mark Jones, Shivangi Prasad,**  
Zein-Eddine Meziani, Chao Peng  
and the Hall C J/ $\Psi$ -007 collaboration

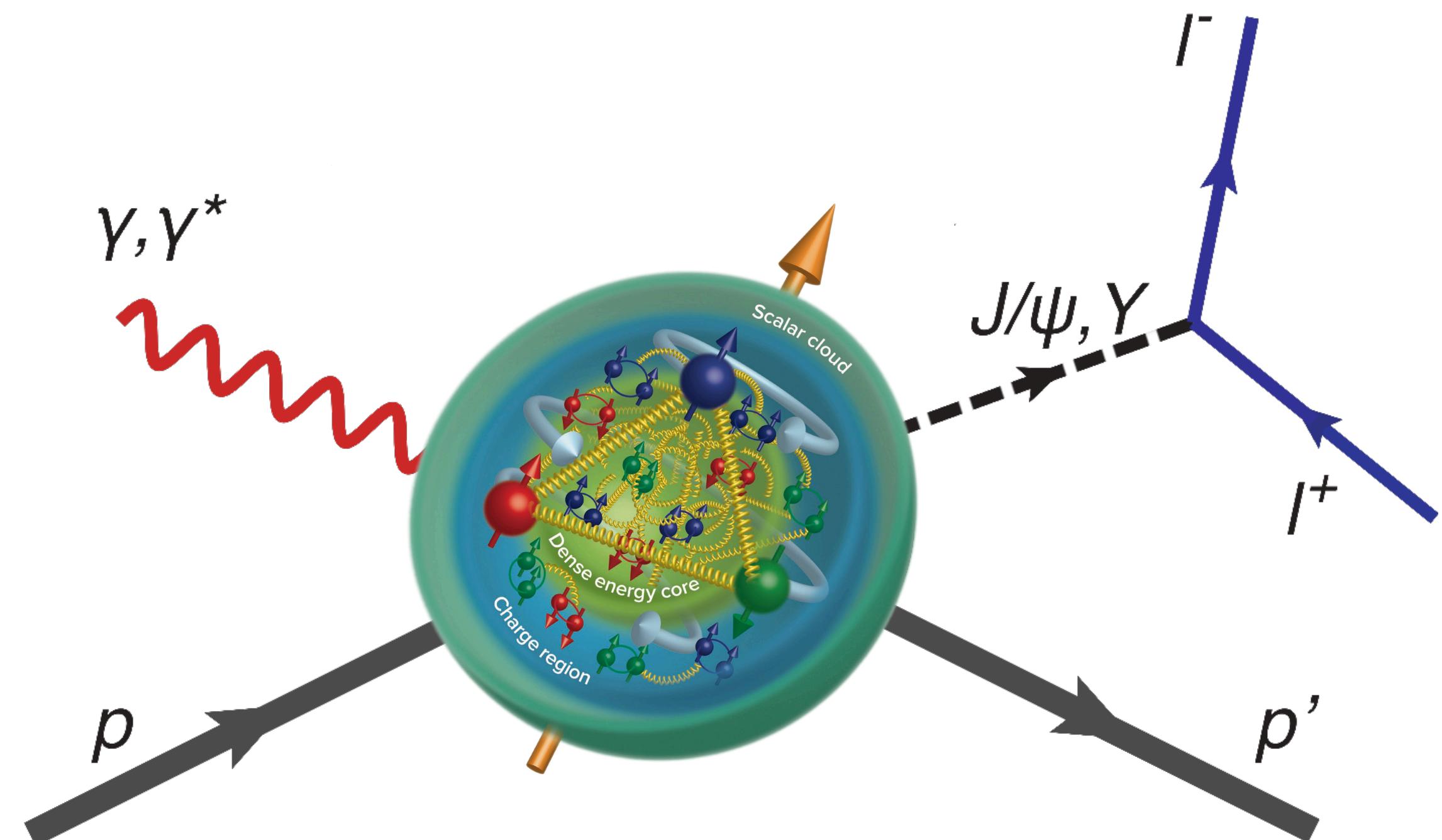
007 $J/\Psi$



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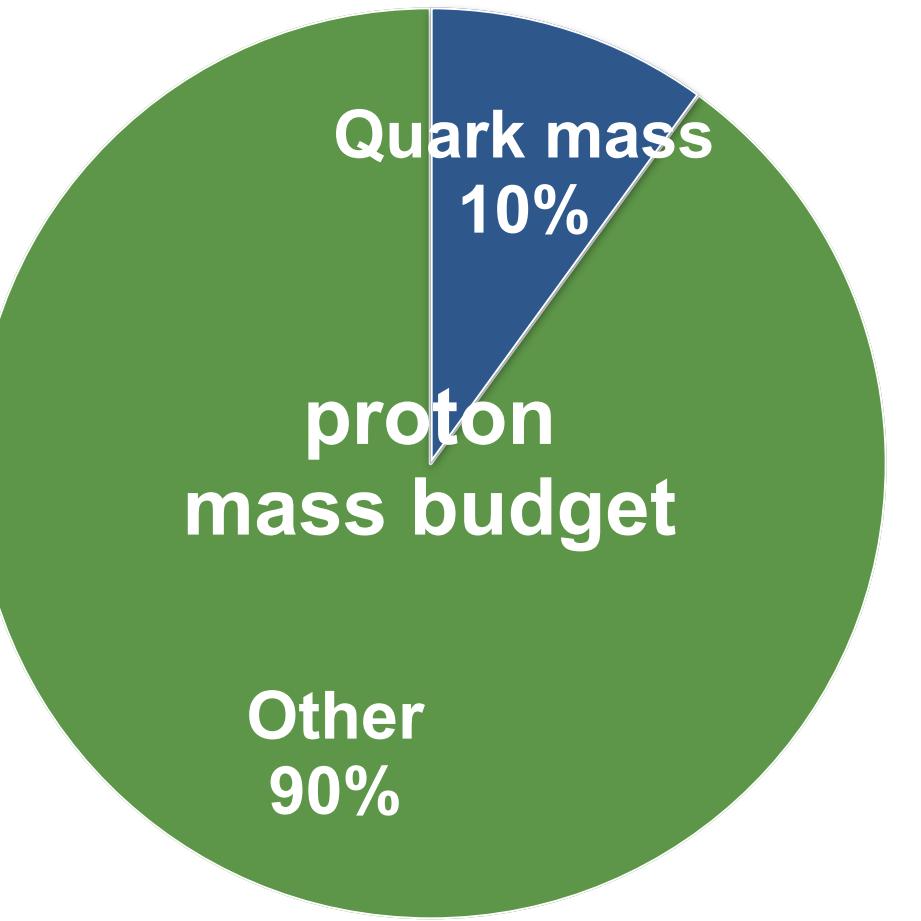
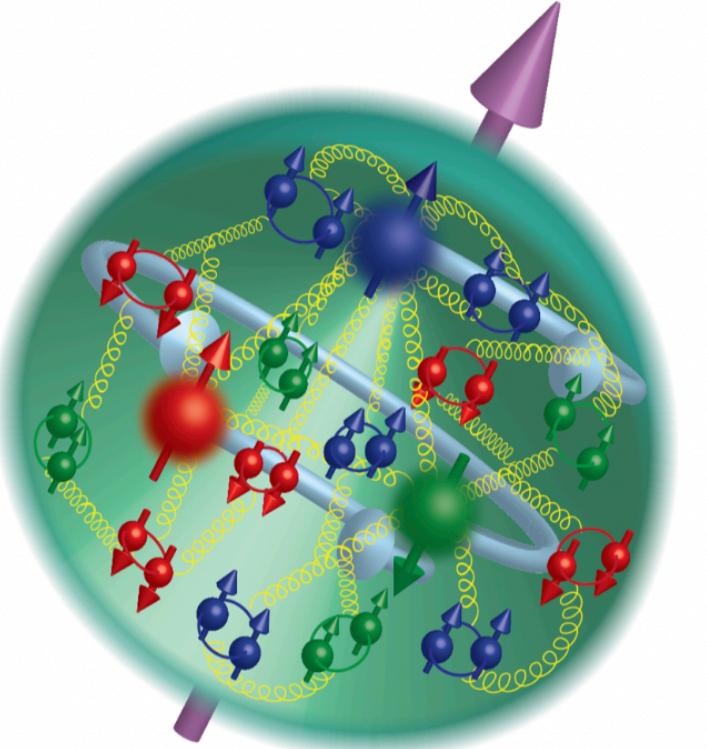
This work is supported by the U.S. Department of  
Energy, Office of Science, Office of Nuclear Physics,  
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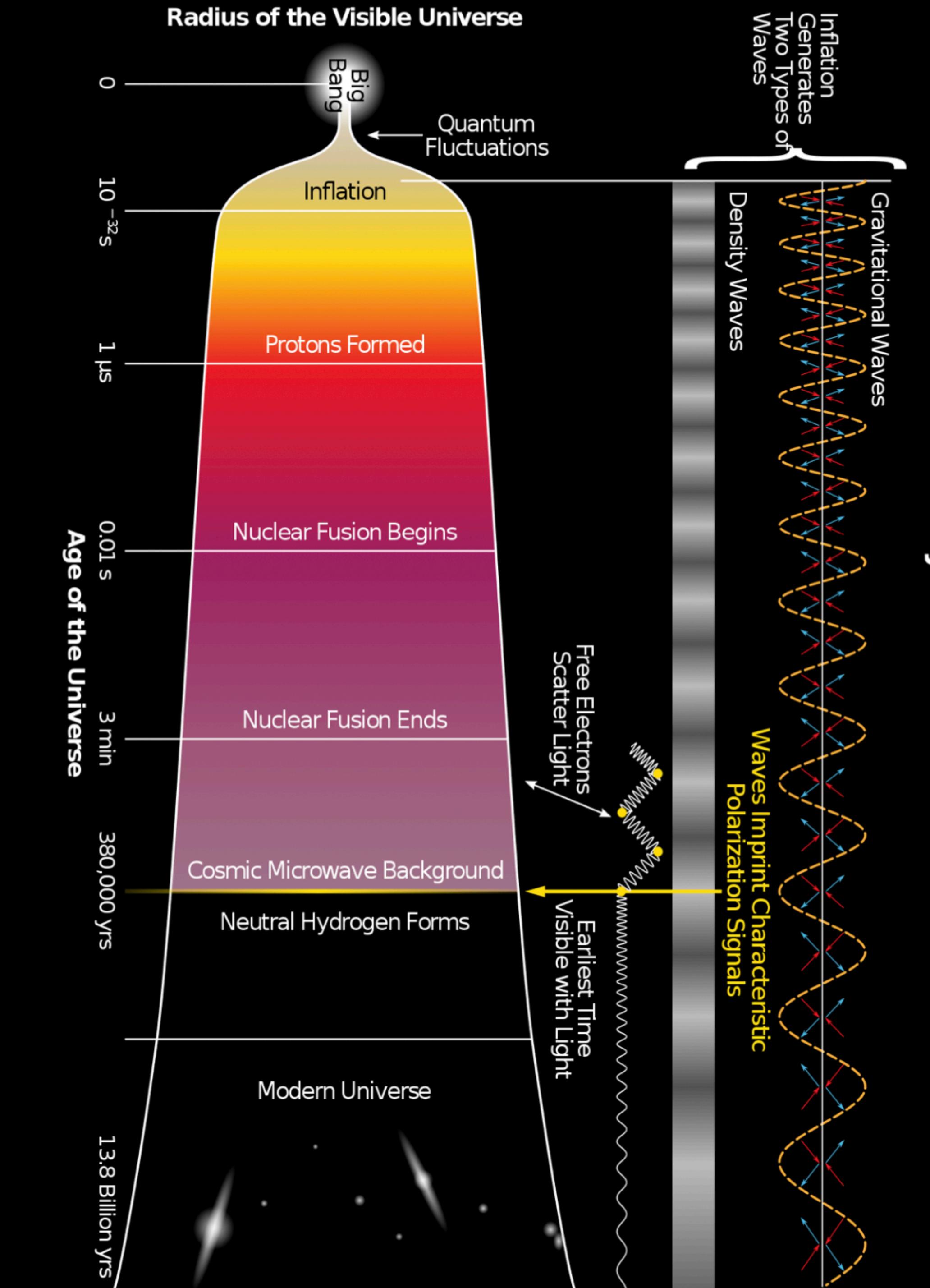


SoLID Science Workshop,  
Argonne, 2024

# The emergence of nucleon mass QCD IN THE STANDARD MODEL

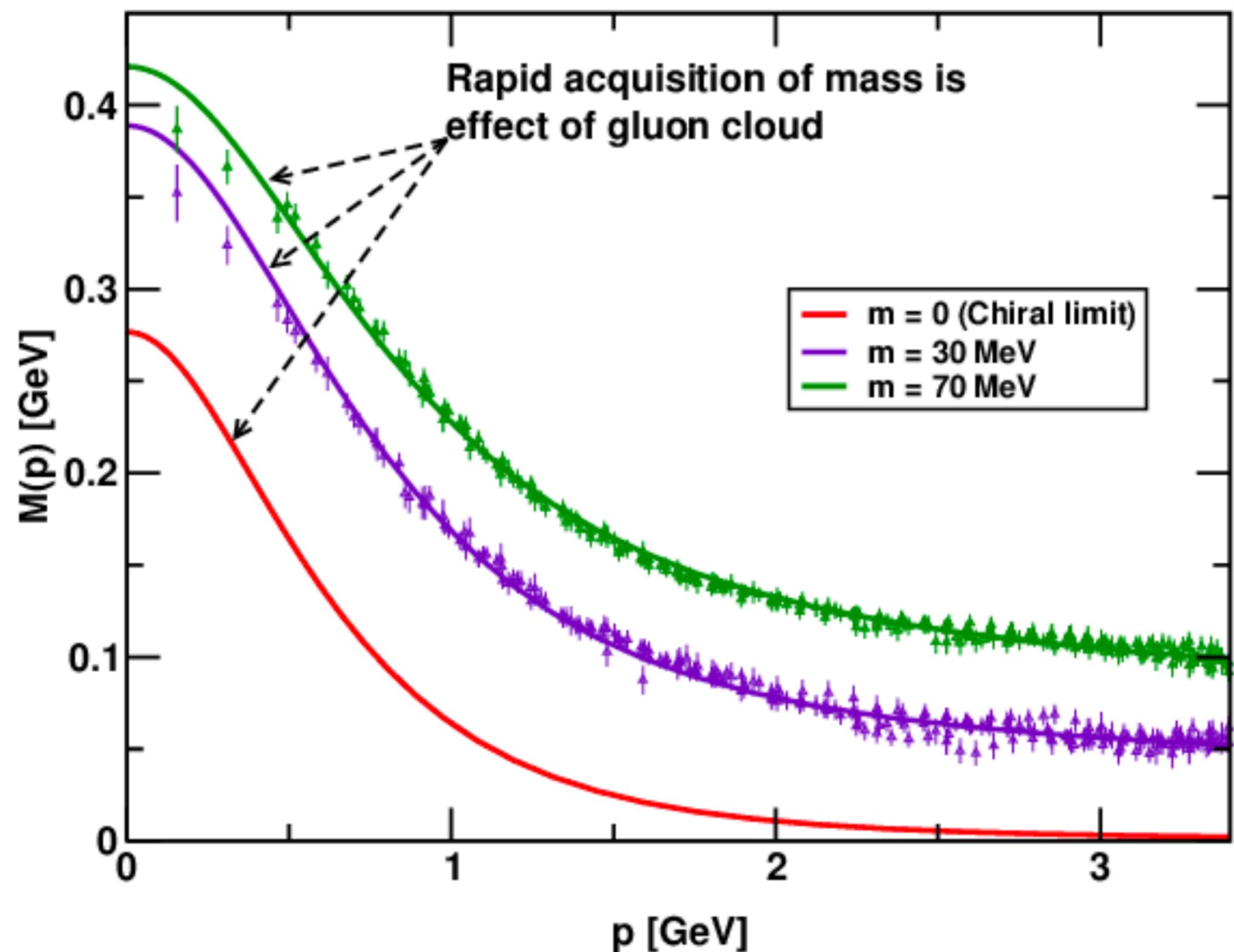


- Since the formation of protons and neutrons, most of the mass of the visible universe encapsulated in protons, neutrons, and nuclei.
- Surprising: nucleon mass much larger than sum of quark masses.
- How does QCD give rise to the 1GeV proton?*
- How is the proton mass distributed in its confinement size?*

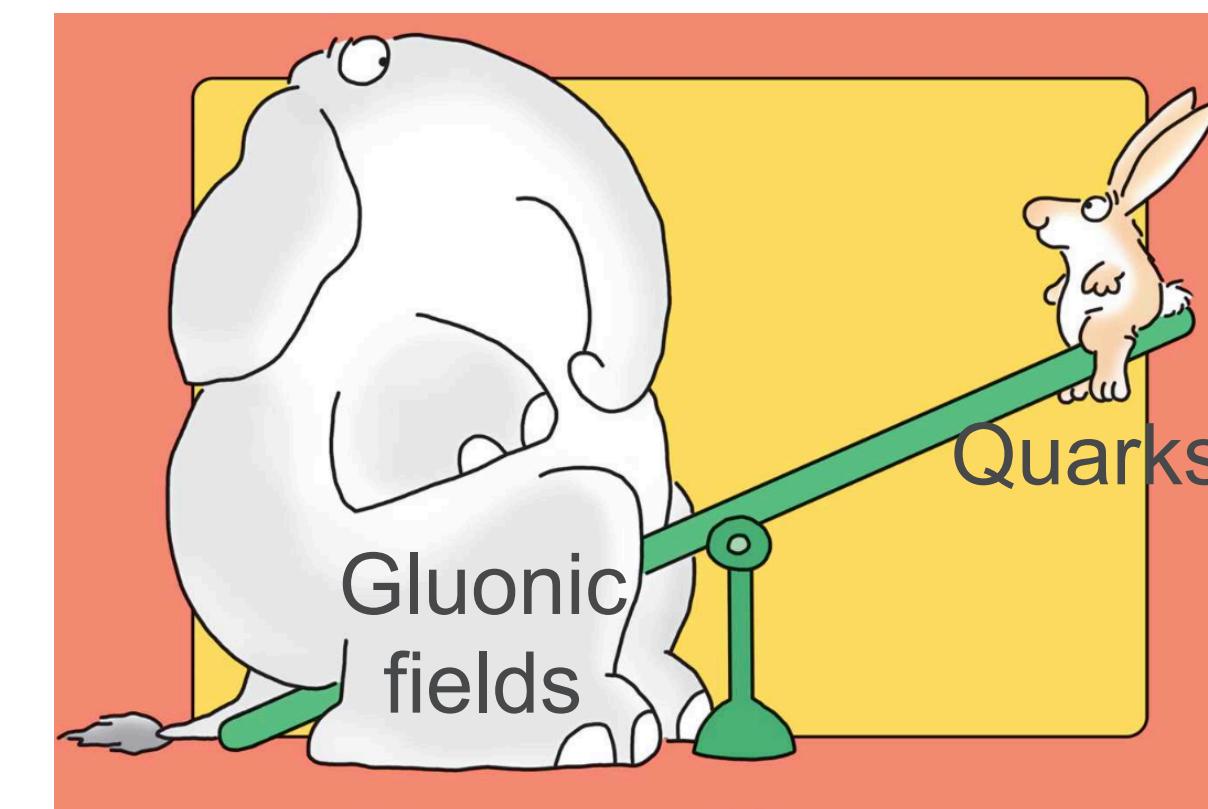


# PROTON MASS IS AN EMERGENT PHENOMENON

## QCD responsible for the proton mass



Most of the proton mass originates in the energy enclosed in the gluonic fields of the Strong Interaction itself

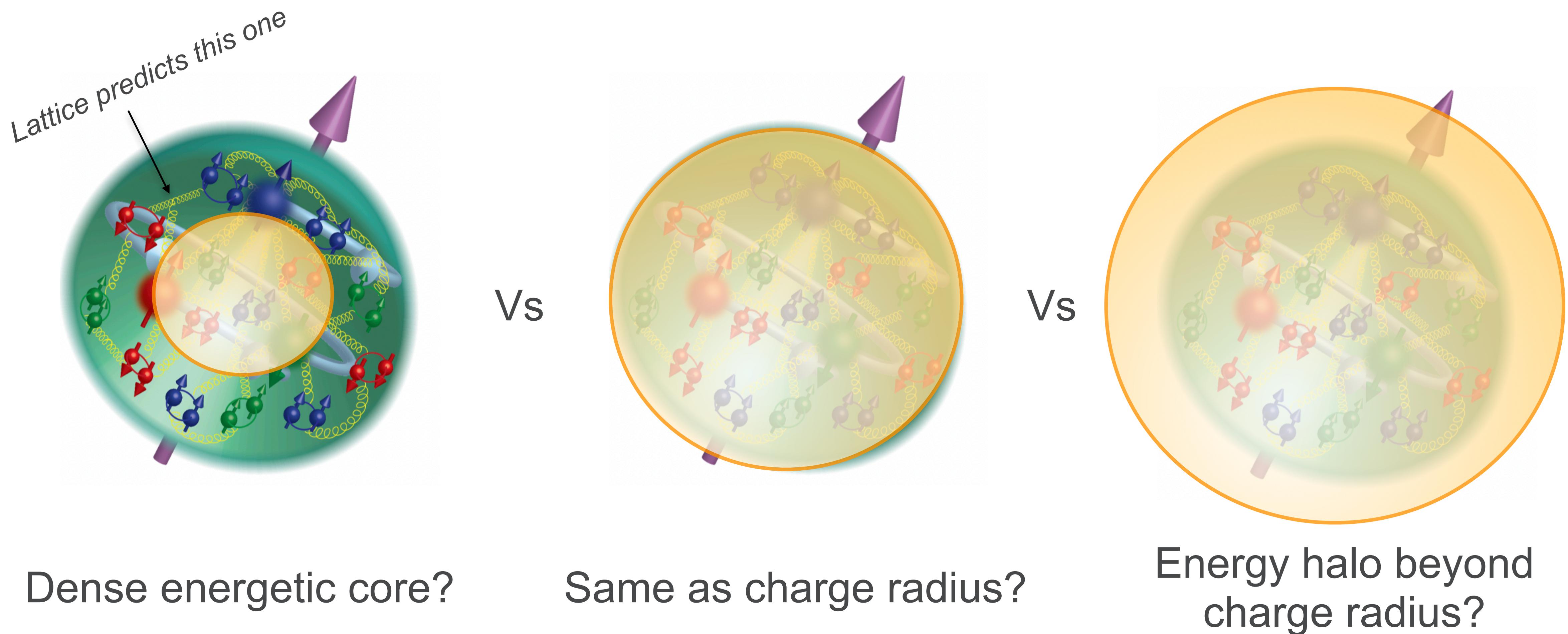


M. S. Bhagwat et al., Phys. Rev. C 68, 015203 (2003)  
I. C. Cloet et al., Prog. Part. Nucl. Phys. 77, 1-69 (2014)

**Bottom line:** The Higgs mechanism is largely irrelevant for most of “normal” visible matter!

# WHERE IS THE ENERGY INSIDE THE PROTON?

## How does the mass radius compare to the charge radius?



Dense energetic core?

Same as charge radius?

Energy halo beyond  
charge radius?

# GRAVITATIONAL FORM FACTORS (GFFS)

## Towards observables for the matter structure of the proton

GFFs are the form factors of the QCD energy-momentum tensor (EMT) for quarks and gluons

$$\langle N' | T_{q,g}^{\mu,\nu} | N \rangle = \bar{u}(N') \left( A_{g,q}(t) \gamma^{\{\mu} P^{\nu\}} + B_{g,q}(t) \frac{i P^{\{\mu} \sigma^{\nu\}} \rho \Delta_\rho}{2M} + C_{g,q}(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{g,q}(t) M g^{\mu\nu} \right) u(N)$$

GFFs encode mechanical properties of the proton:

- $A_{g,q}(t)$ : Related to quark and gluon momenta,  $A_{g,q}(0) = \langle x_{q,g} \rangle$
- $J_{g,q}(t) = 1/2 (A_{g,q}(t) + B_{g,q}(t))$ : Related to angular momentum,  $J_{\text{tot}}(0) = 1/2$
- $D_{g,q}(t) = 4C_{g,q}(t)$ : Related to pressure and shear forces

Tensor (2++ graviton-like, mass) radius

$$\langle r_m^2 \rangle_g = \frac{6}{A_g(0)} \frac{dA_g(t)}{dt} \Big|_{t=0} - \frac{6}{A_g(0)} \frac{C_g(0)}{M_N^2}$$

Scalar (0++ glueball-like) radius

$$\langle r_s^2 \rangle_g = \frac{6}{A_g(0)} \frac{dA_g(t)}{dt} \Big|_{t=0} - \frac{18}{A_g(0)} \frac{C_g(0)}{M_N^2}$$

Both radii depend on the functional form of the gluonic  $A$  and  $C$  form factors at the origin.

# HOW TO MEASURE THE GLUONIC GFFS

## Gluons are elusive!

- Cannot use Electromagnetic probe: primarily couples to quarks
- Cannot use Weak probe: also primarily couples to quarks
- Cannot use hadronic probe made of light quarks: primarily sensitive to quark structure
- Cannot use direct gravitational probe: interaction too weak

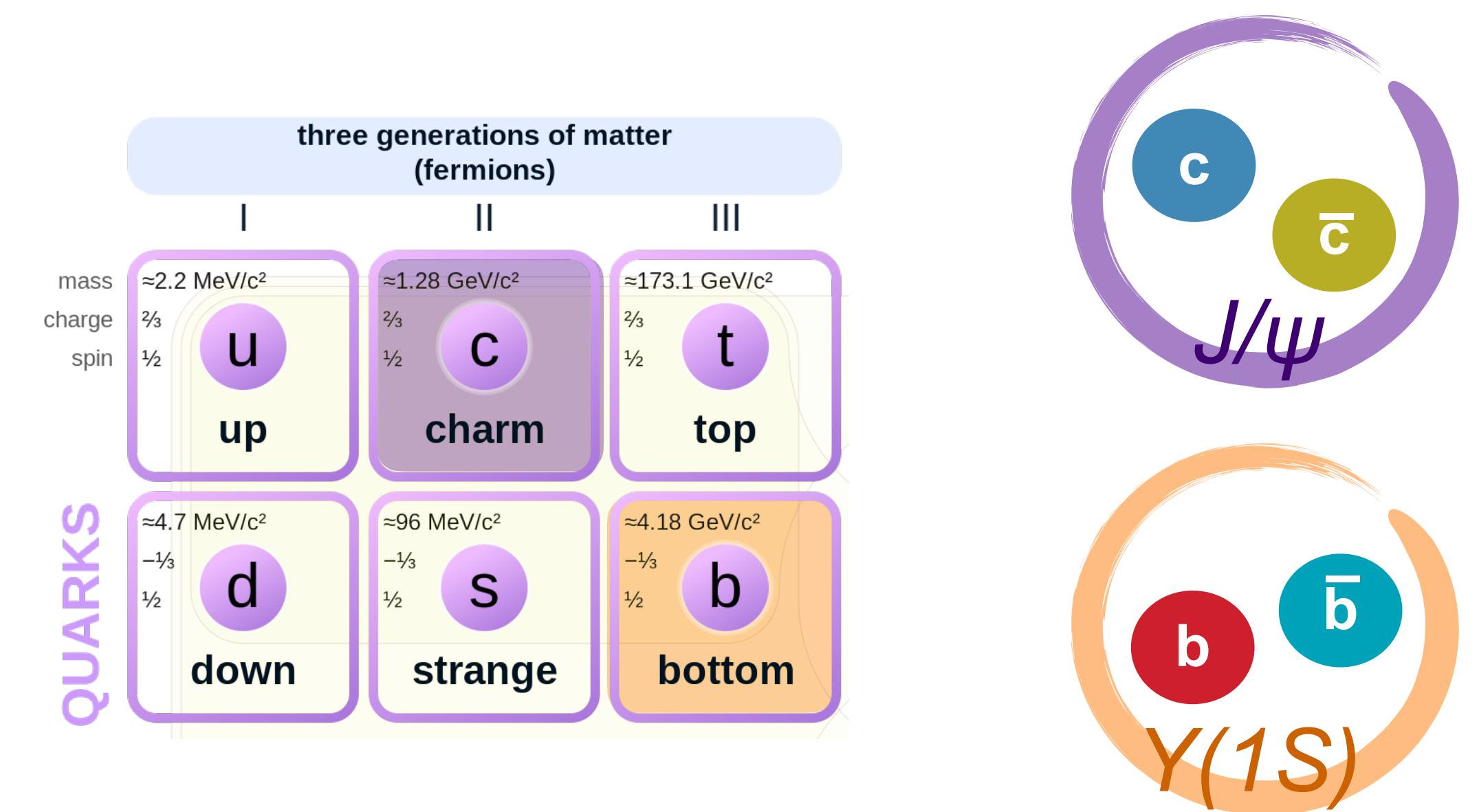
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- Small “color” dipole made of heavy quarks well-suited to study gluons

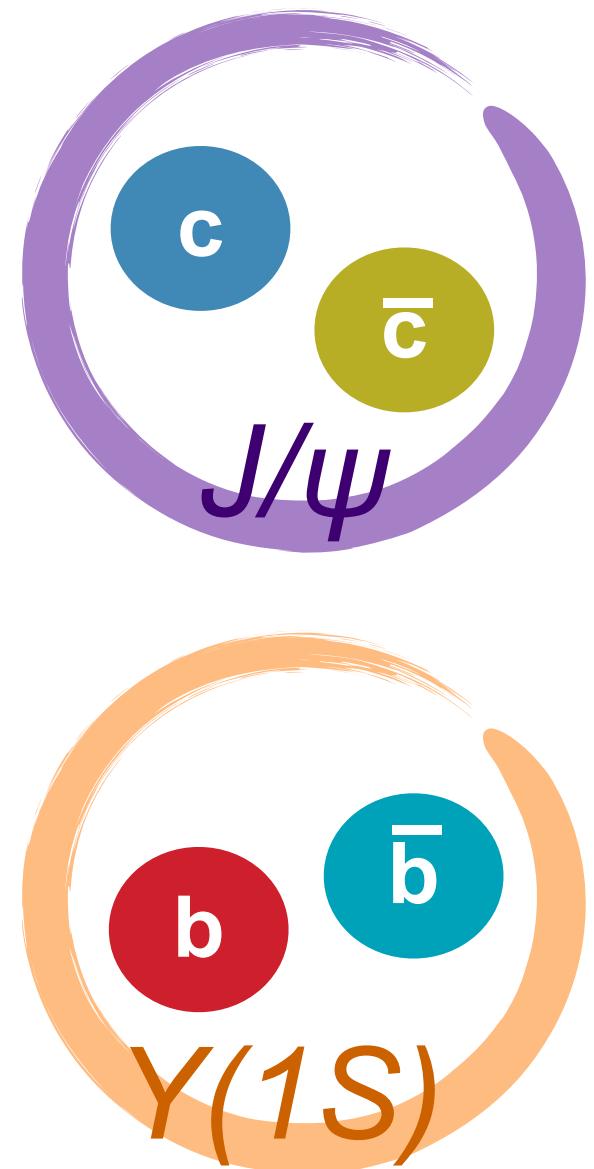
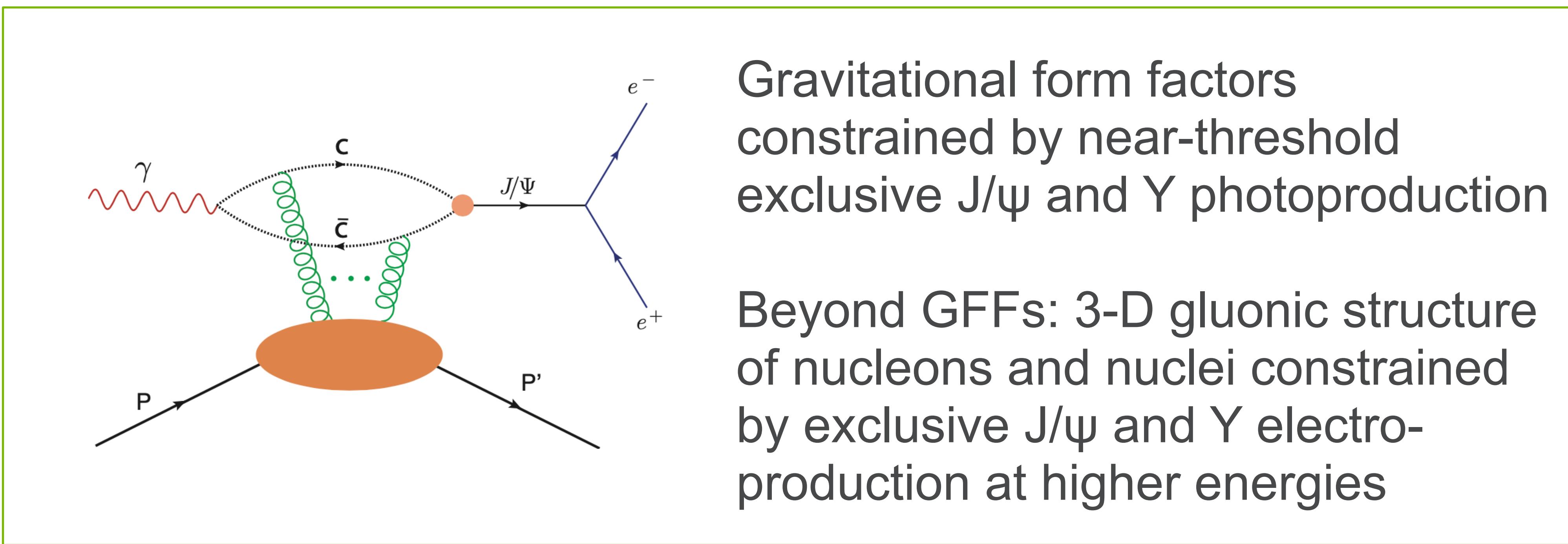
three generations of matter (fermions)			
mass	I	II	III
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	u up	c charm	t top
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	d down	s strange	b bottom
mass	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$
charge	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$



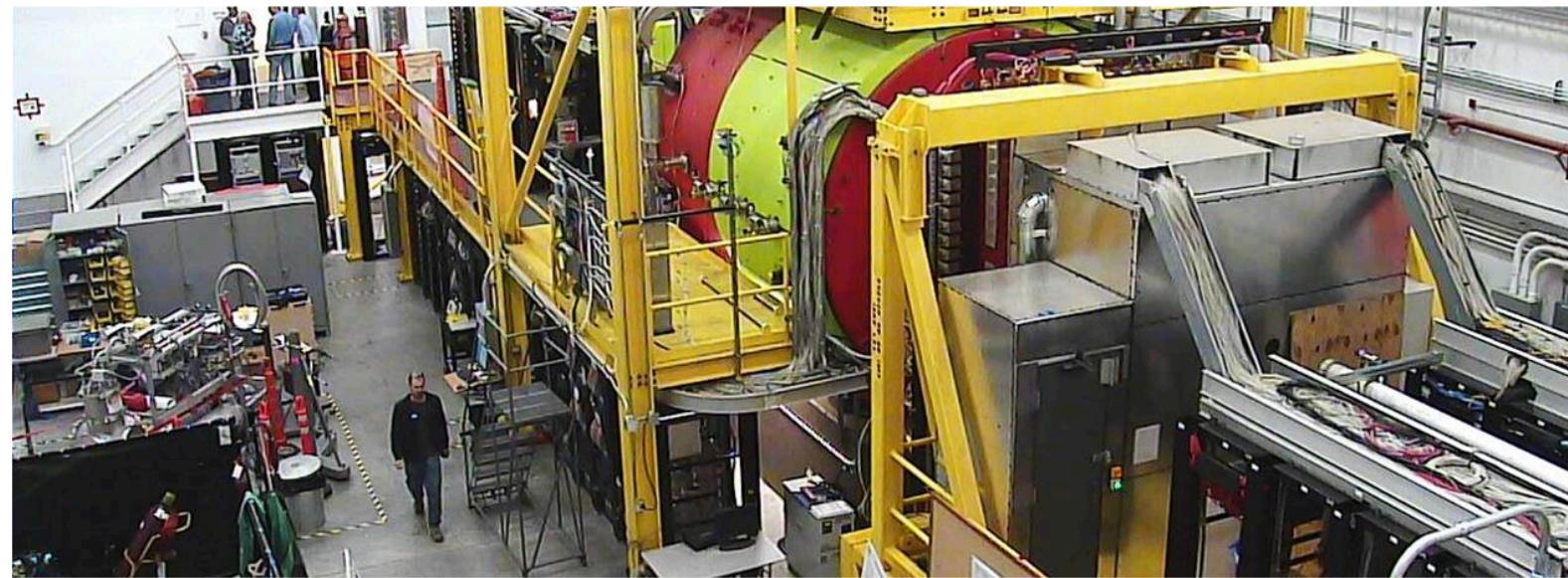
# HOW TO MEASURE THE GLUONIC GFFS

## Gluons are elusive!

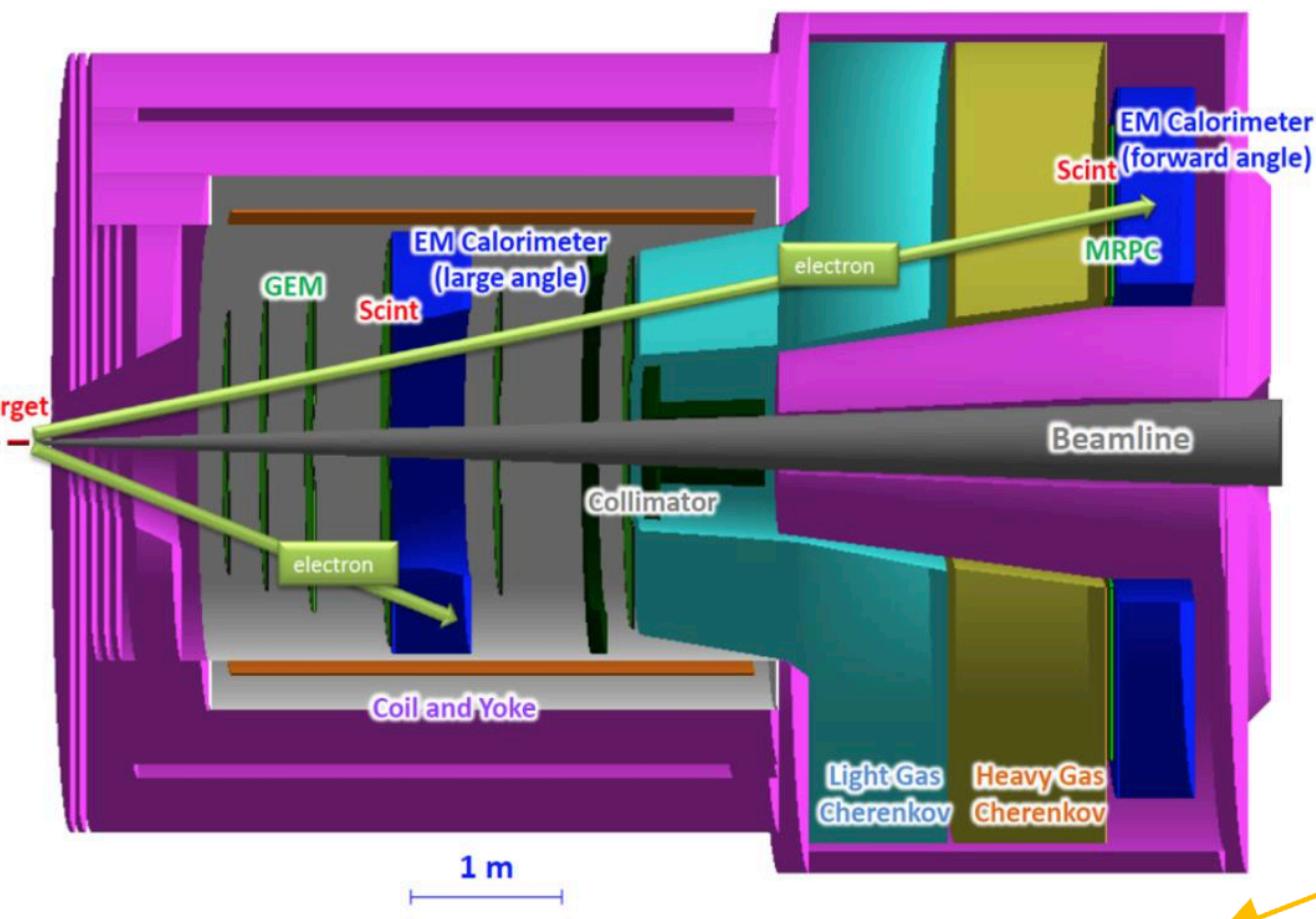
- Cannot use Electromagnetic probe: primarily couples to quarks
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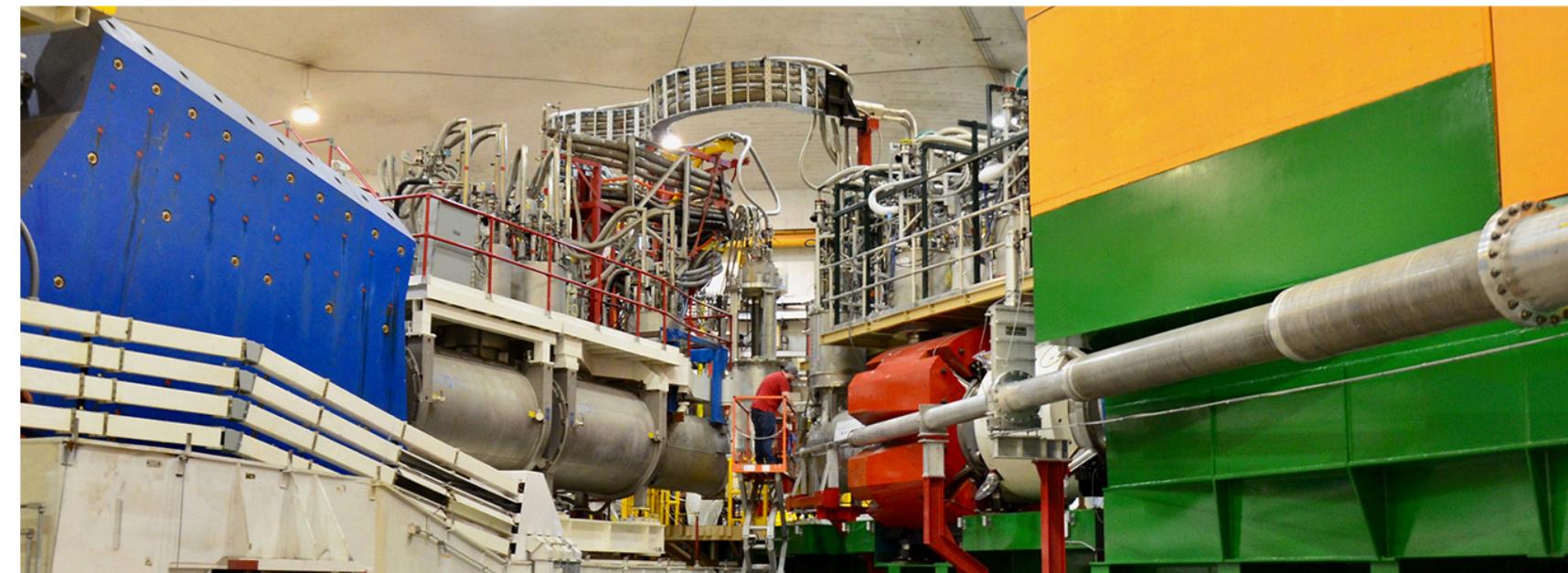
# 12 GEV J/ $\psi$ EXPERIMENTS AT JEFFERSON LAB



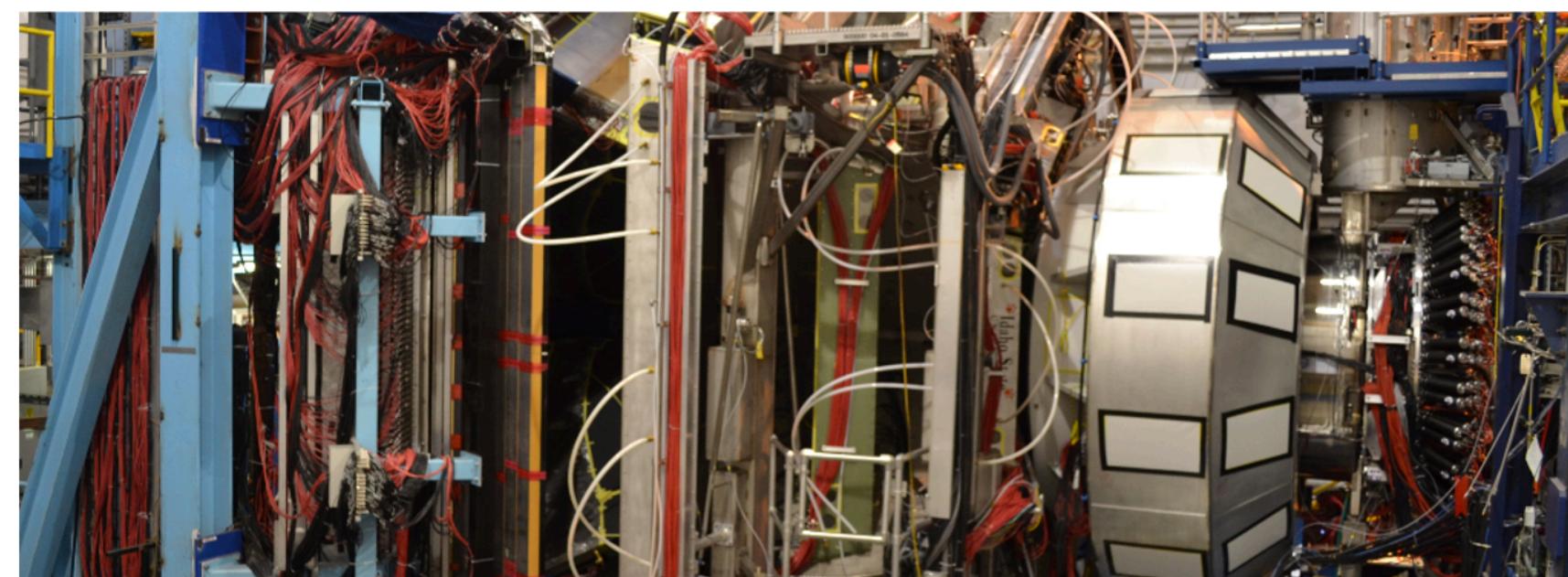
**Hall D - GlueX** observer the first J/ $\psi$  at JLab  
A. Ali *et al.*, PRL 123, 072001 (2019)



**Hall A** has experiment E12-12-006 at **SoLID** to measure J/ $\psi$  in electro- and photoproduction, and an LOI to measure double polarization using **SBS**



**Hall C** has the J/ $\psi$ -007 experiment (E12-16-007)  
LHCb hidden-charm pentaquark search



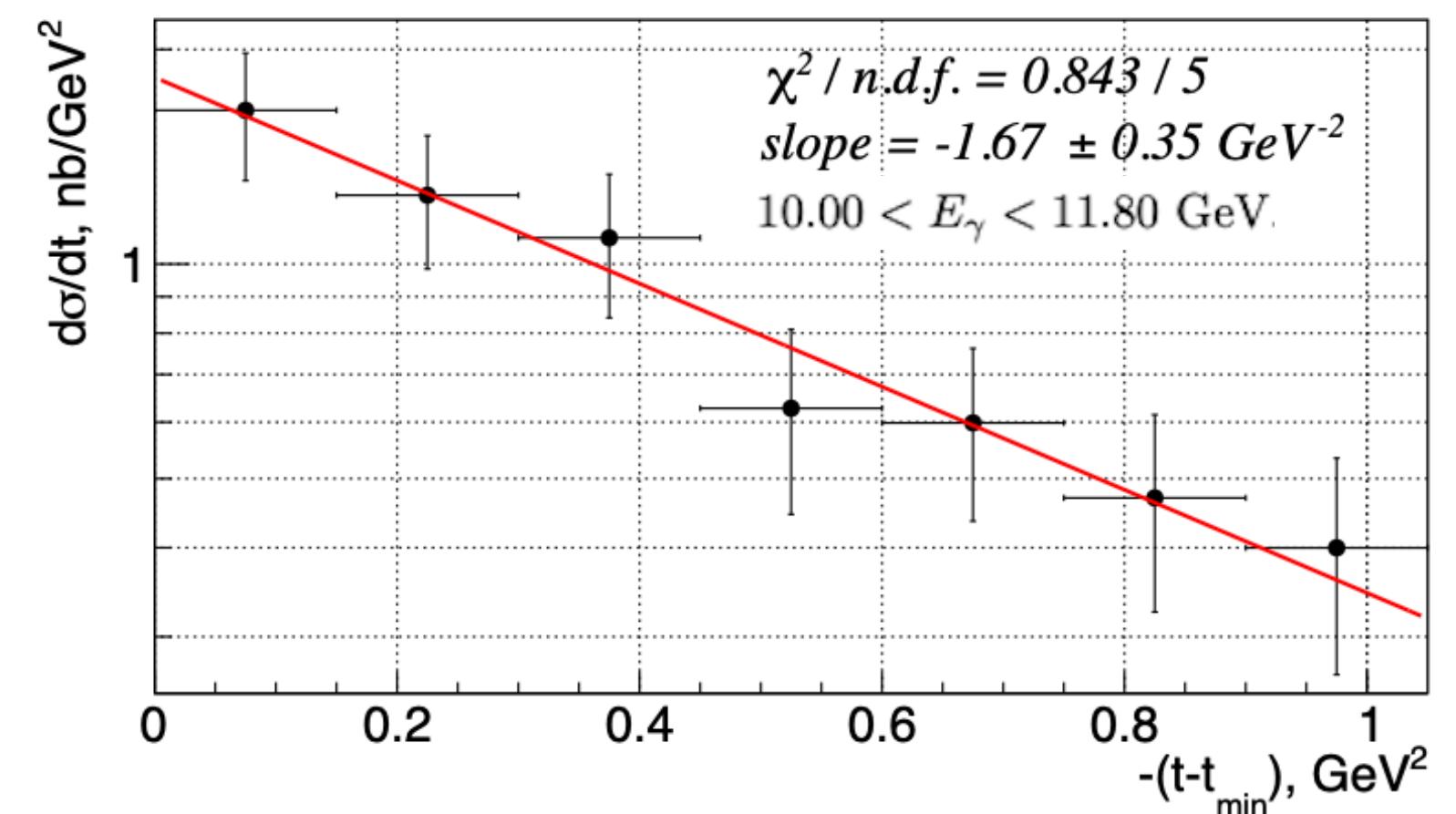
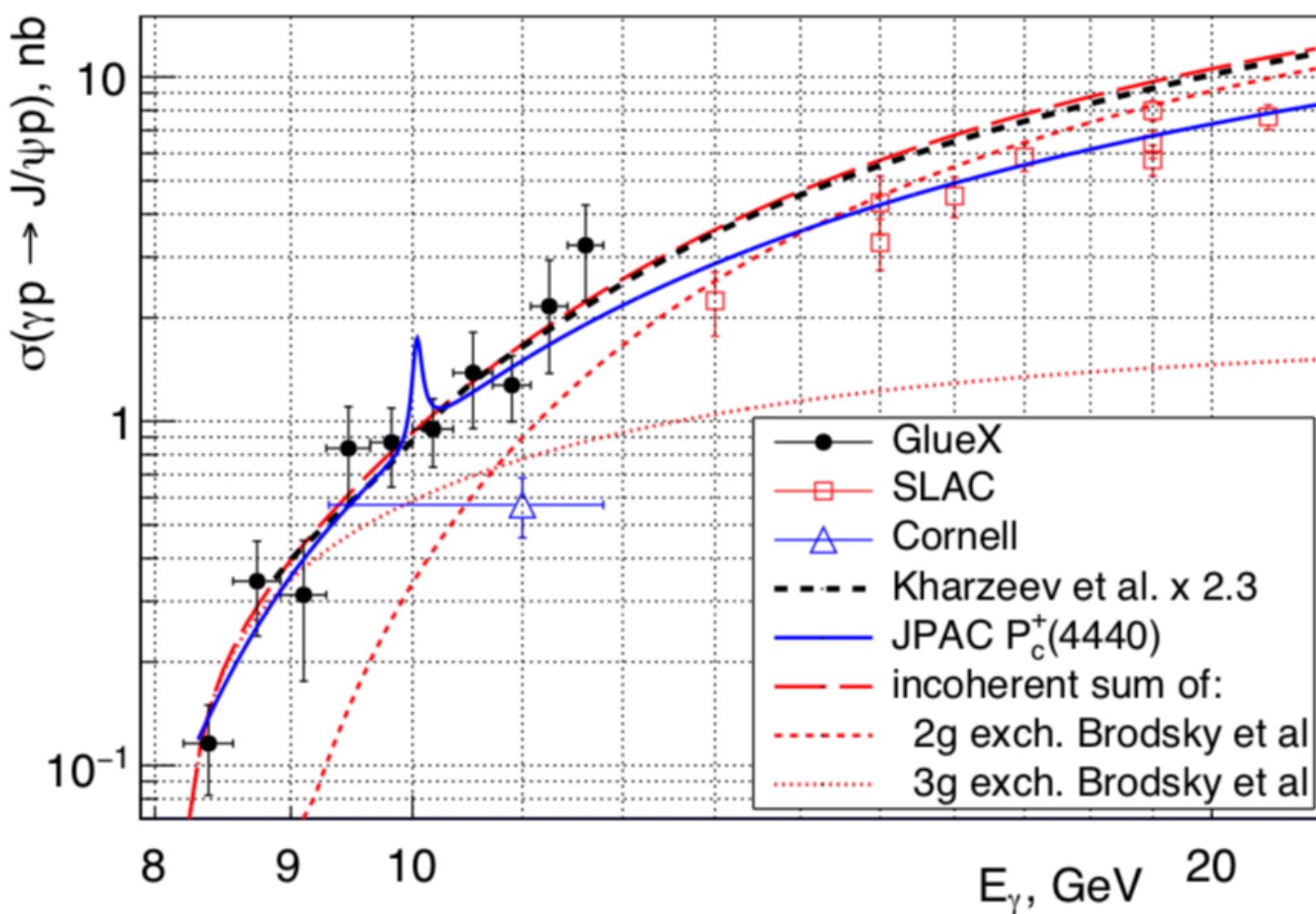
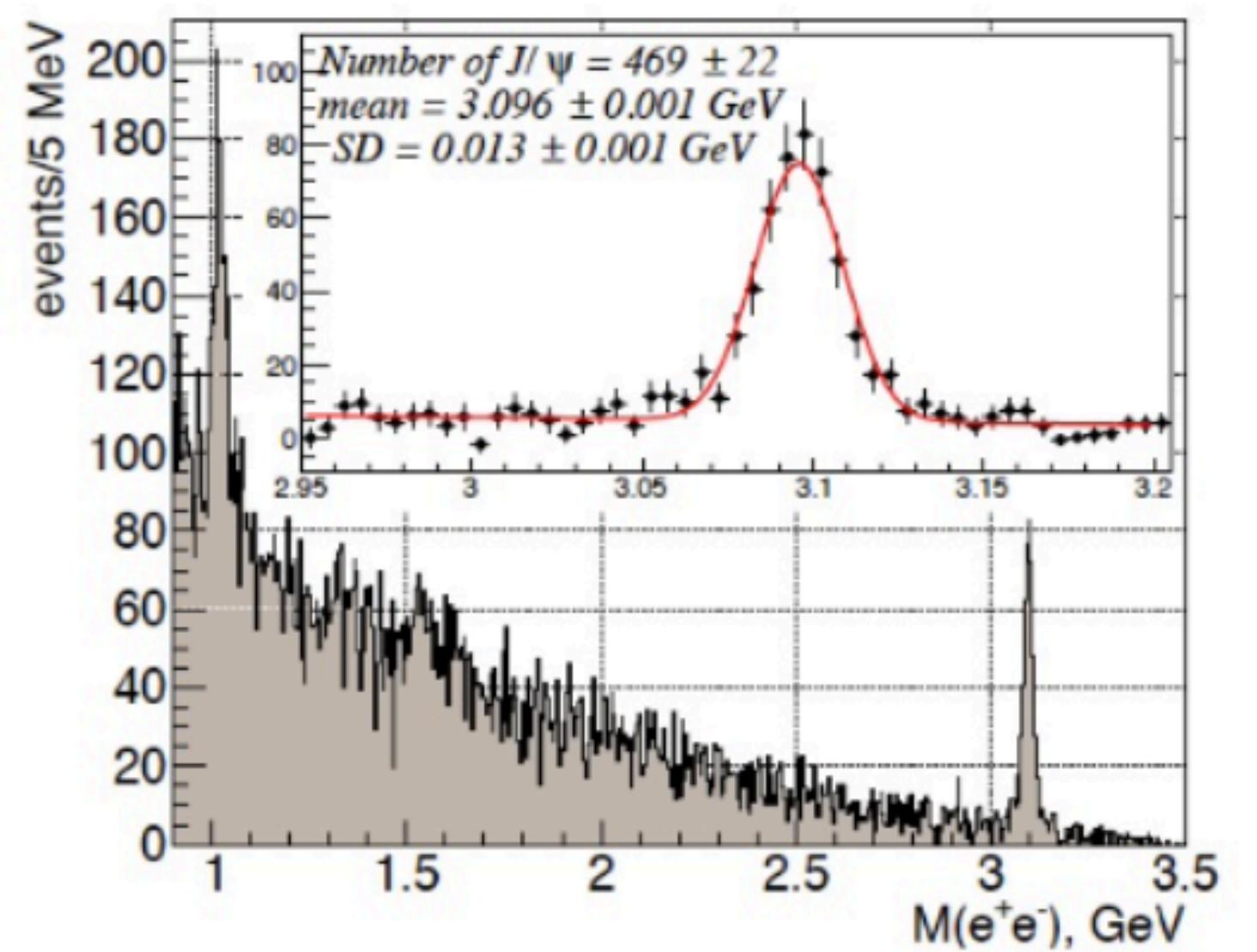
**Hall B - CLAS12** has experiments to measure TCS + J/ $\psi$  in photoproduction as part of Run Groups A (hydrogen) and B (deuterium): E12-12-001, E12-12-001A, E12-11-003B

# J/ψ NEAR THRESHOLD IN HALL D

First J/ψ results from JLab, published in PRL 123, 072001 (2019)

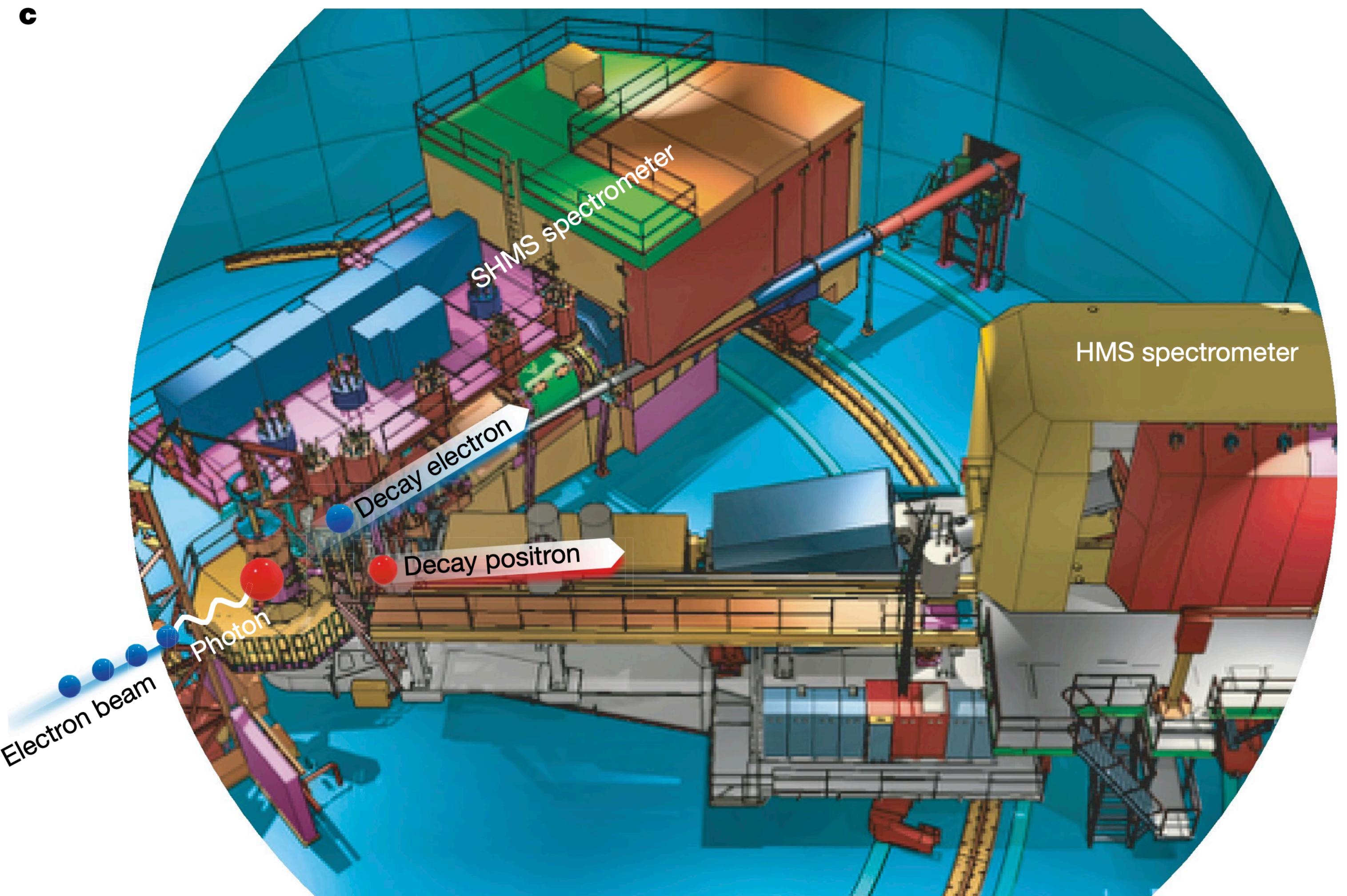
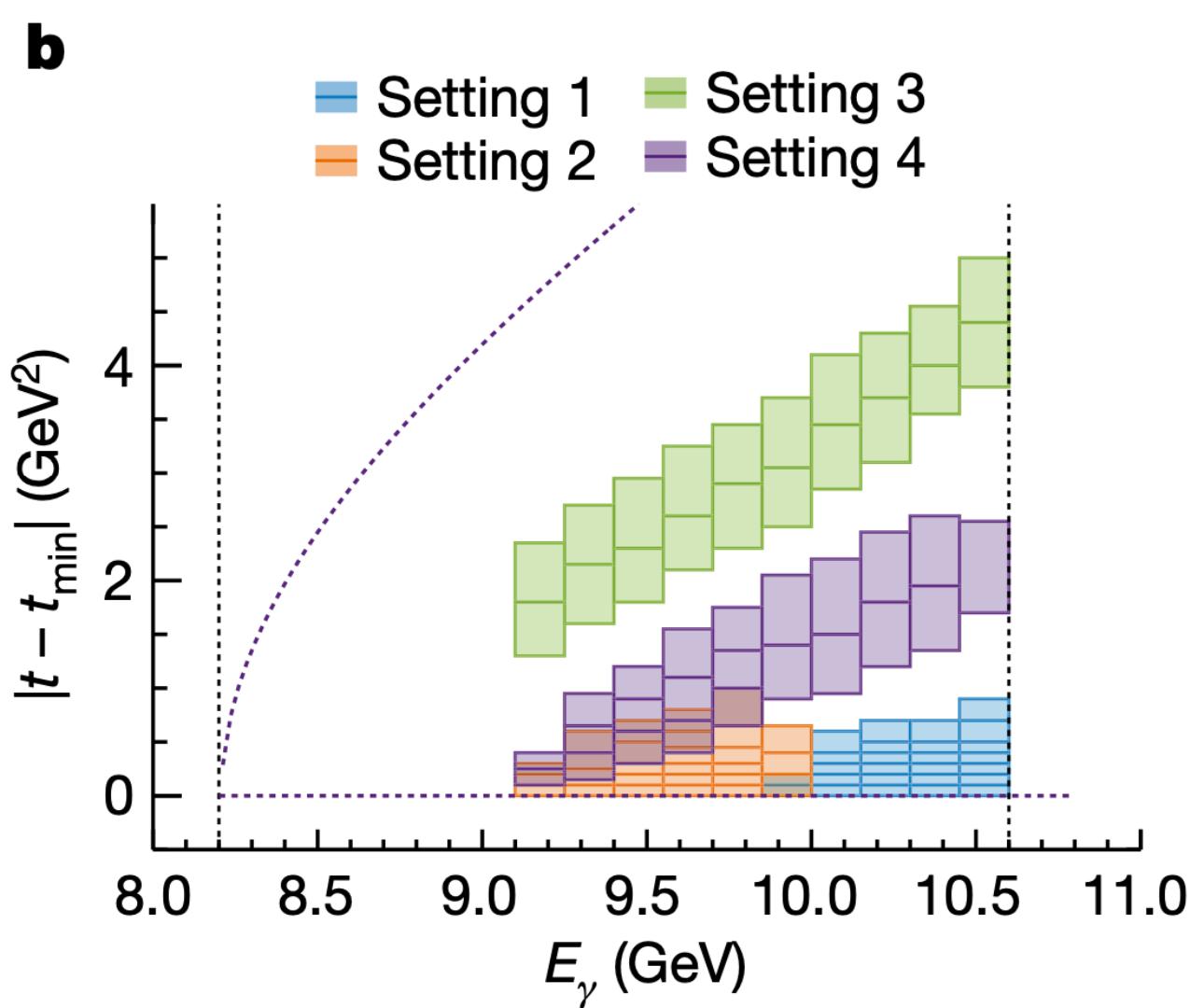
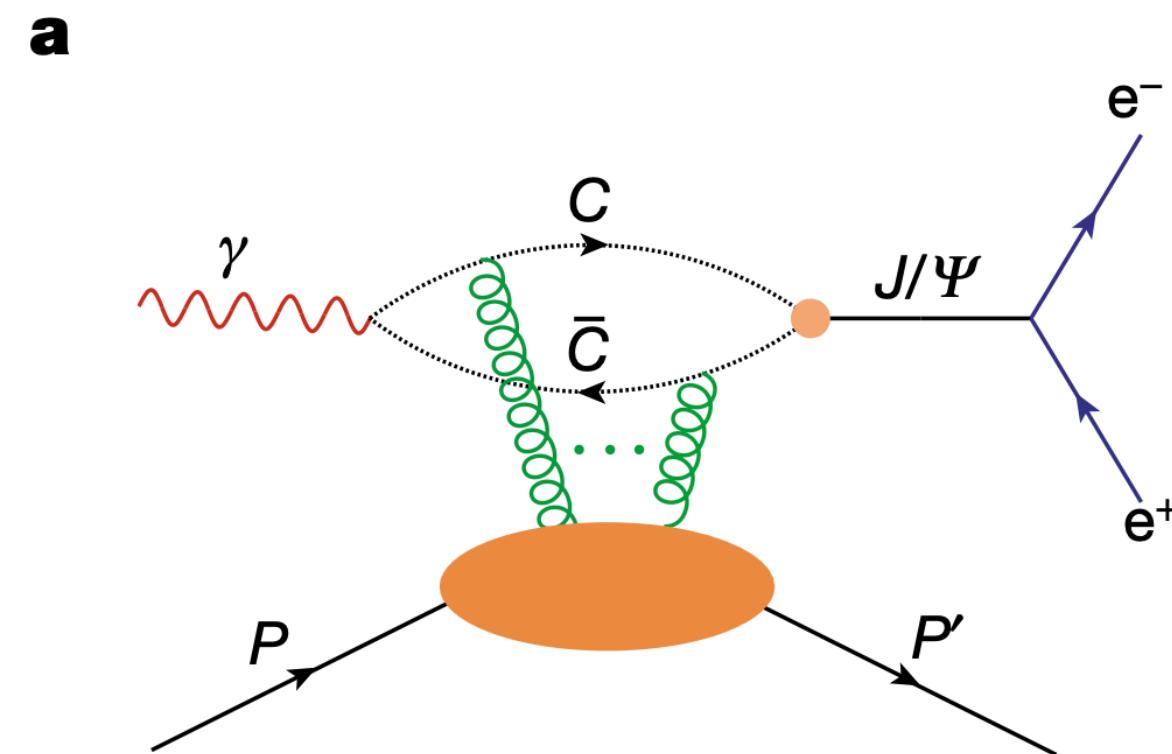
- 1D cross section (~469 counts)
- Trends significantly higher than old measurements
- Single 1D t-profile spurred on many new theoretical calculations
- Did not see evidence for hidden-charm pentaquarks

$\gamma p \rightarrow p J/\psi \rightarrow pe^+e^-$



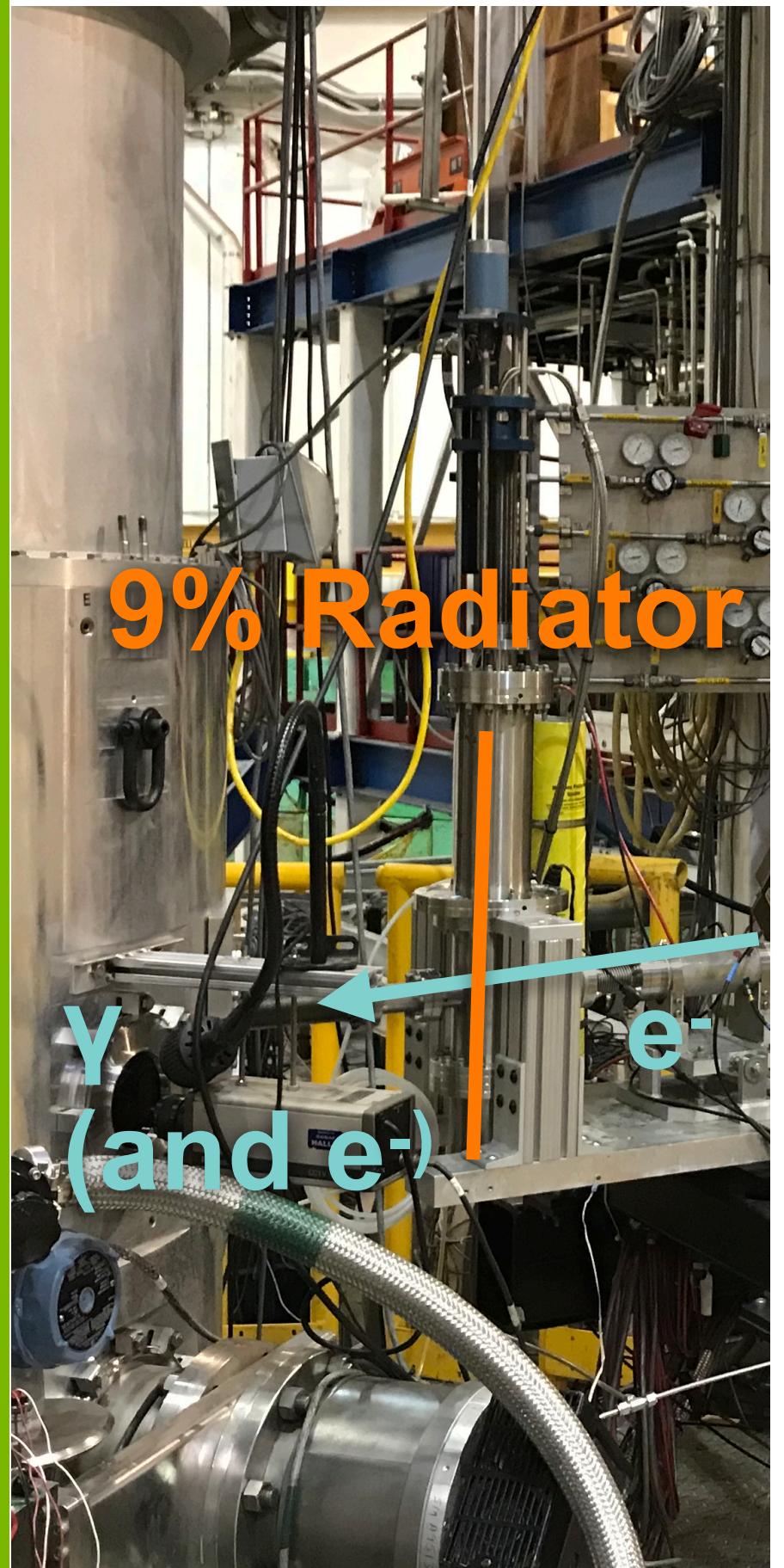
# Near-threshold J/ψ photoproduction

# HALL C EXPERIMENT E12-16-007



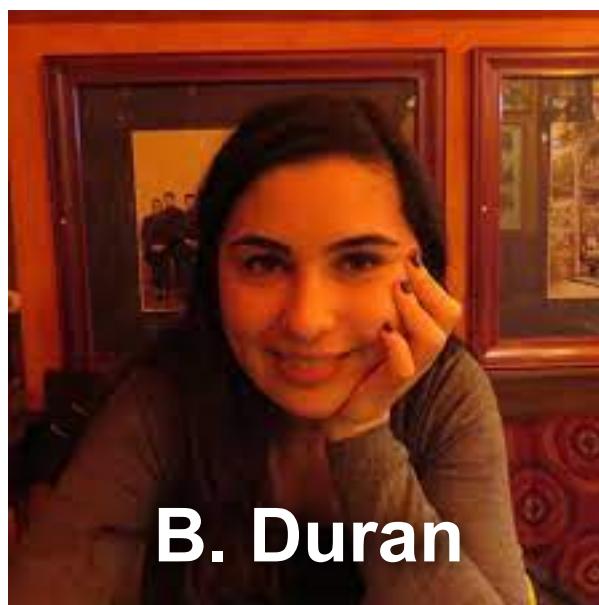
# HALL C EXPERIMENT E12-16-007

**J/ψ-007: “Search for the LHCb Pentaquark”, ran in 2019**



# THE J/Ψ-007 COLLABORATION

007<sup>J/ψ</sup>



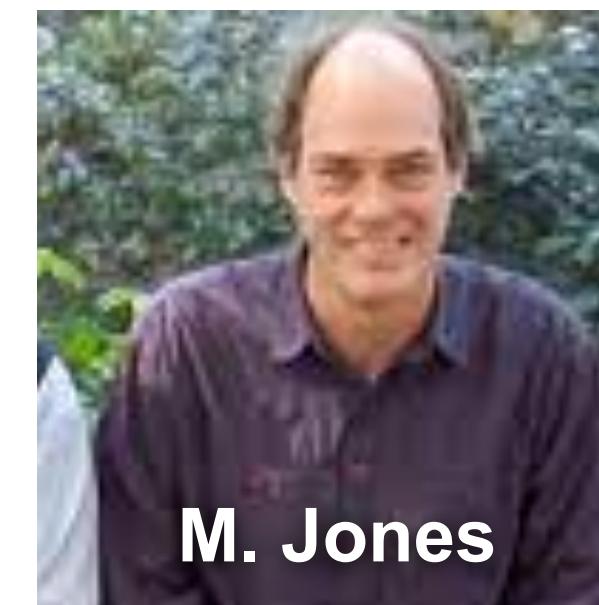
B. Duran



Z.-E. Meziani



S. Joosten



M. Jones



S. Prasad



C. Peng



W. Armstrong



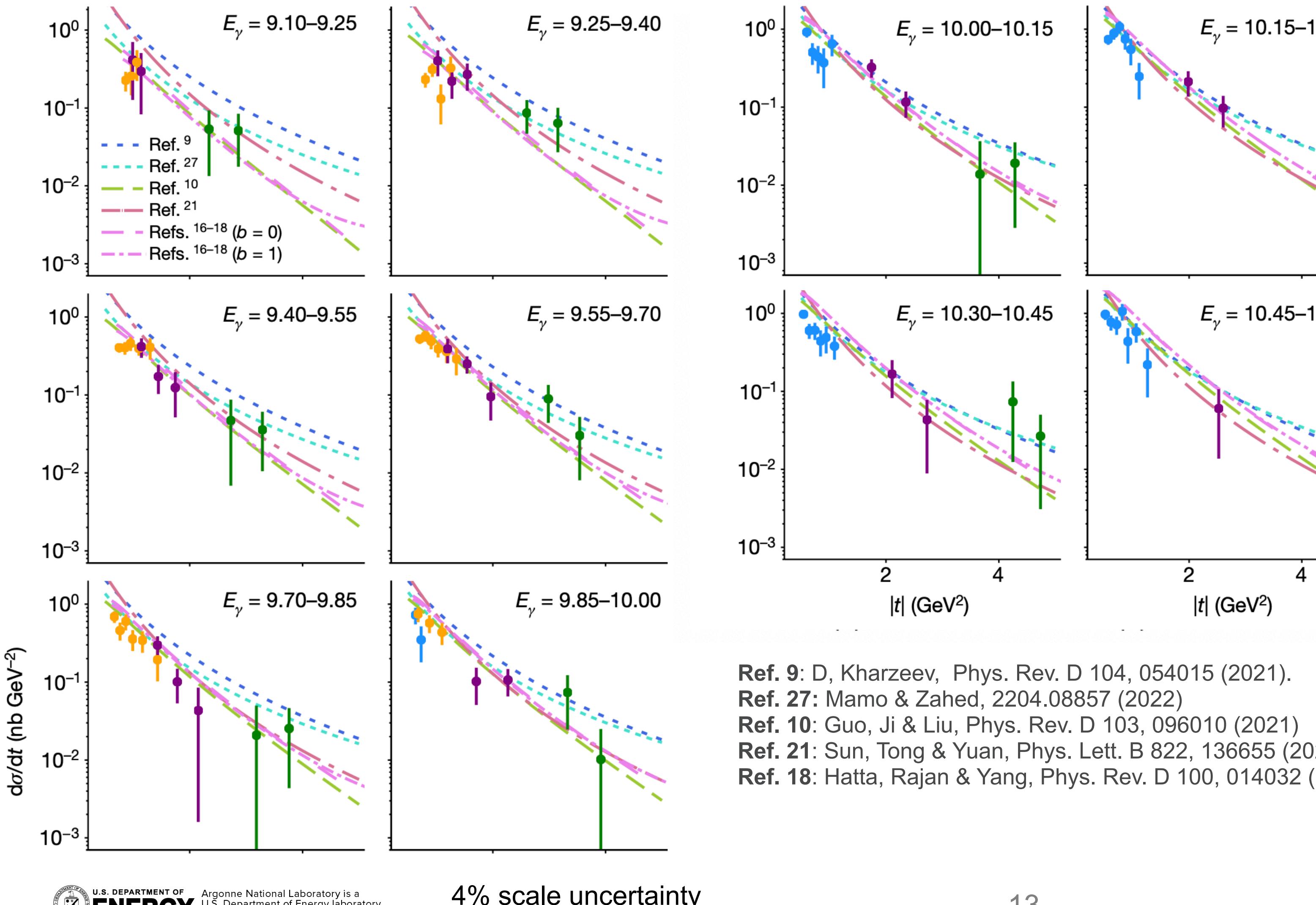
M. Paolone

*...and many others!*

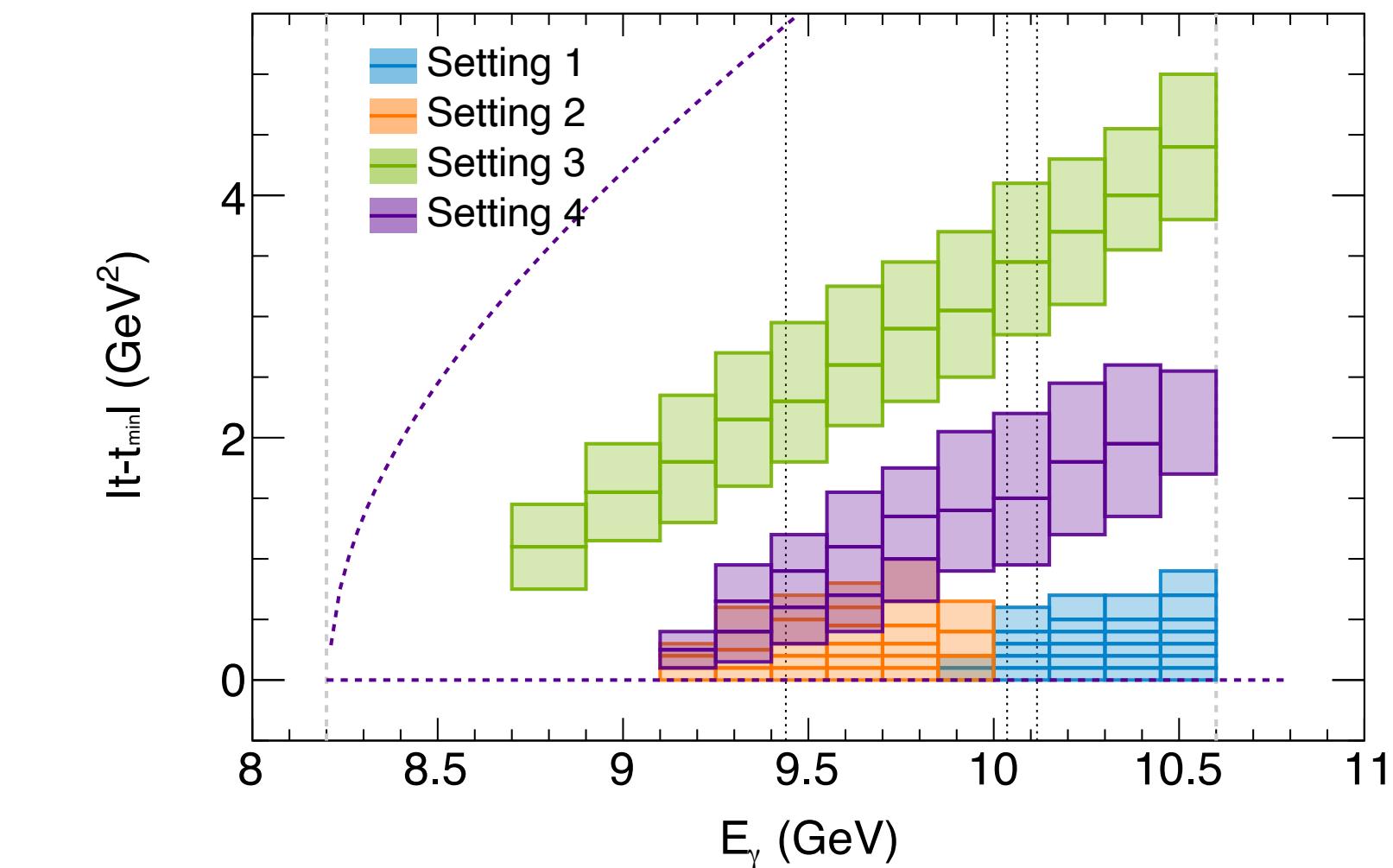
B. Duran<sup>1,2</sup>, Z.-E. Meziani<sup>1,2</sup>✉, S. Joosten<sup>1</sup>, M. K. Jones<sup>3</sup>, S. Prasad<sup>1</sup>, C. Peng<sup>1</sup>, W. Armstrong<sup>1</sup>,  
H. Atac<sup>2</sup>, E. Chudakov<sup>3</sup>, H. Bhatt<sup>4</sup>, D. Bhetuwal<sup>4</sup>, M. Boer<sup>5</sup>, A. Camsonne<sup>3</sup>, J.-P. Chen<sup>3</sup>,  
M. M. Dalton<sup>3</sup>, N. Deokar<sup>2</sup>, M. Diefenthaler<sup>3</sup>, J. Dunne<sup>4</sup>, L. El Fassi<sup>4</sup>, E. Fuchey<sup>6</sup>, H. Gao<sup>7</sup>,  
D. Gaskell<sup>3</sup>, O. Hansen<sup>3</sup>, F. Hauenstein<sup>8</sup>, D. Higinbotham<sup>3</sup>, S. Jia<sup>2</sup>, A. Karki<sup>4</sup>, C. Keppel<sup>3</sup>,  
P. King<sup>9</sup>, H. S. Ko<sup>10</sup>, X. Li<sup>7</sup>, R. Li<sup>2</sup>, D. Mack<sup>3</sup>, S. Malace<sup>3</sup>, M. McCaughan<sup>3</sup>, R. E. McClellan<sup>11</sup>,  
R. Michaels<sup>3</sup>, D. Meekins<sup>3</sup>, Michael Paolone<sup>2</sup>, L. Pentchev<sup>3</sup>, E. Pooser<sup>3</sup>, A. Puckett<sup>6</sup>,  
R. Radloff<sup>9</sup>, M. Rehfuss<sup>2</sup>, P. E. Reimer<sup>1</sup>, S. Riordan<sup>1</sup>, B. Sawatzky<sup>3</sup>, A. Smith<sup>7</sup>, N. Sparveris<sup>2</sup>,  
H. Szumila-Vance<sup>3</sup>, S. Wood<sup>3</sup>, J. Xie<sup>1</sup>, Z. Ye<sup>1</sup>, C. Yero<sup>8</sup> & Z. Zhao<sup>7</sup>

# First results published in Nature!

## 2-D J/Ψ CROSS SECTIONS NEAR THRESHOLD



Ref. 9: D. Kharzeev, Phys. Rev. D 104, 054015 (2021).  
 Ref. 27: Mamo & Zahed, 2204.08857 (2022)  
 Ref. 10: Guo, Ji & Liu, Phys. Rev. D 103, 096010 (2021)  
 Ref. 21: Sun, Tong & Yuan, Phys. Lett. B 822, 136655 (2021)  
 Ref. 18: Hatta, Rajan & Yang, Phys. Rev. D 100, 014032 (2019)



Unfolded 2D cross section results  
compared to various model predictions  
informed by the 2019 1D GlueX results  
All models work reasonably well at higher  
energies but deviate at lower energies

# MODEL ASSUMPTIONS AND CAVEATS

## First model-dependent attempt to determine the GFFs from experiment

### Assumptions

Neglect  $B(t)$  - in concordance with both models and lattice QCD

Neglect  $\bar{C}_g$  when evaluating the cross section and radii (\*)

Assume tripole shape for  $A(t)$  and  $C(t)$  (\*\*)

Fix  $A(0)$  to the average gluon PDF from CT18

Both models fit the data well ( $\chi^2 \sim 1$ )

(\*) This is appropriate for the holographic model but not the GPD model. See Hatta et al. JHEP 12 (2018) 008 & Tanaka, K. JHEP 03 (2023) 013 for a calculation of  $\bar{C}_g = -\bar{C}_q$

(\*\*) Doing the same extraction with a dipole shape, or does not impact our results

### Holographic Model

K. Mamo & I. Zahed, PRD 103, 094010 (2021) and 2204.08857 (2022)

$$\frac{d\sigma}{dt} = \mathcal{N} \times \frac{e^2}{64\pi(s - m_N^2)^2} \times \frac{A(-t, \kappa_T) + \eta^2 D(-t, \kappa_T, \kappa_S)]^2}{A^2(0)} \times F(\tilde{s}) \times 8$$

$\mathcal{N}$  normalized to the previous world data (not given by the model)

### GPD Model

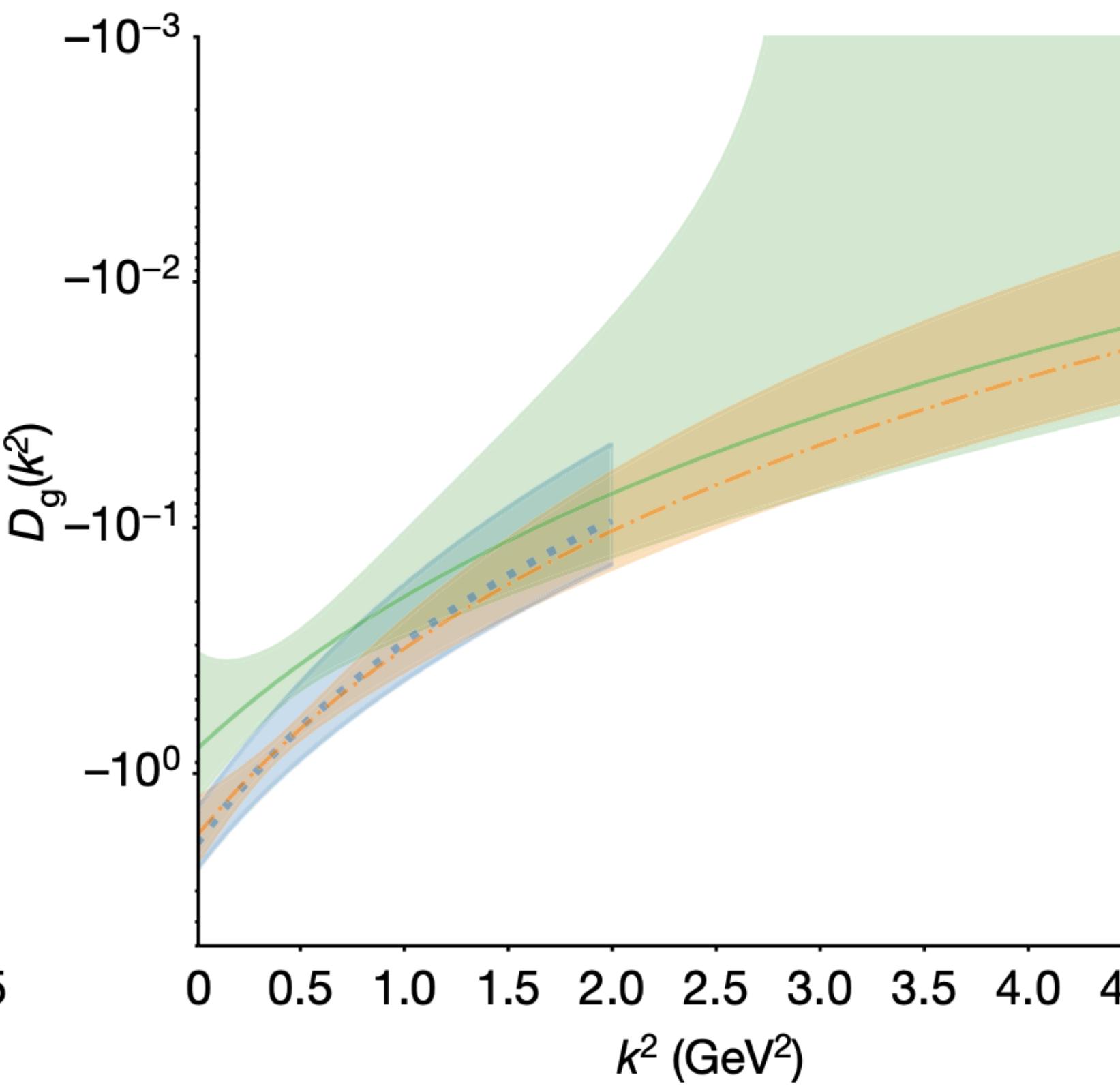
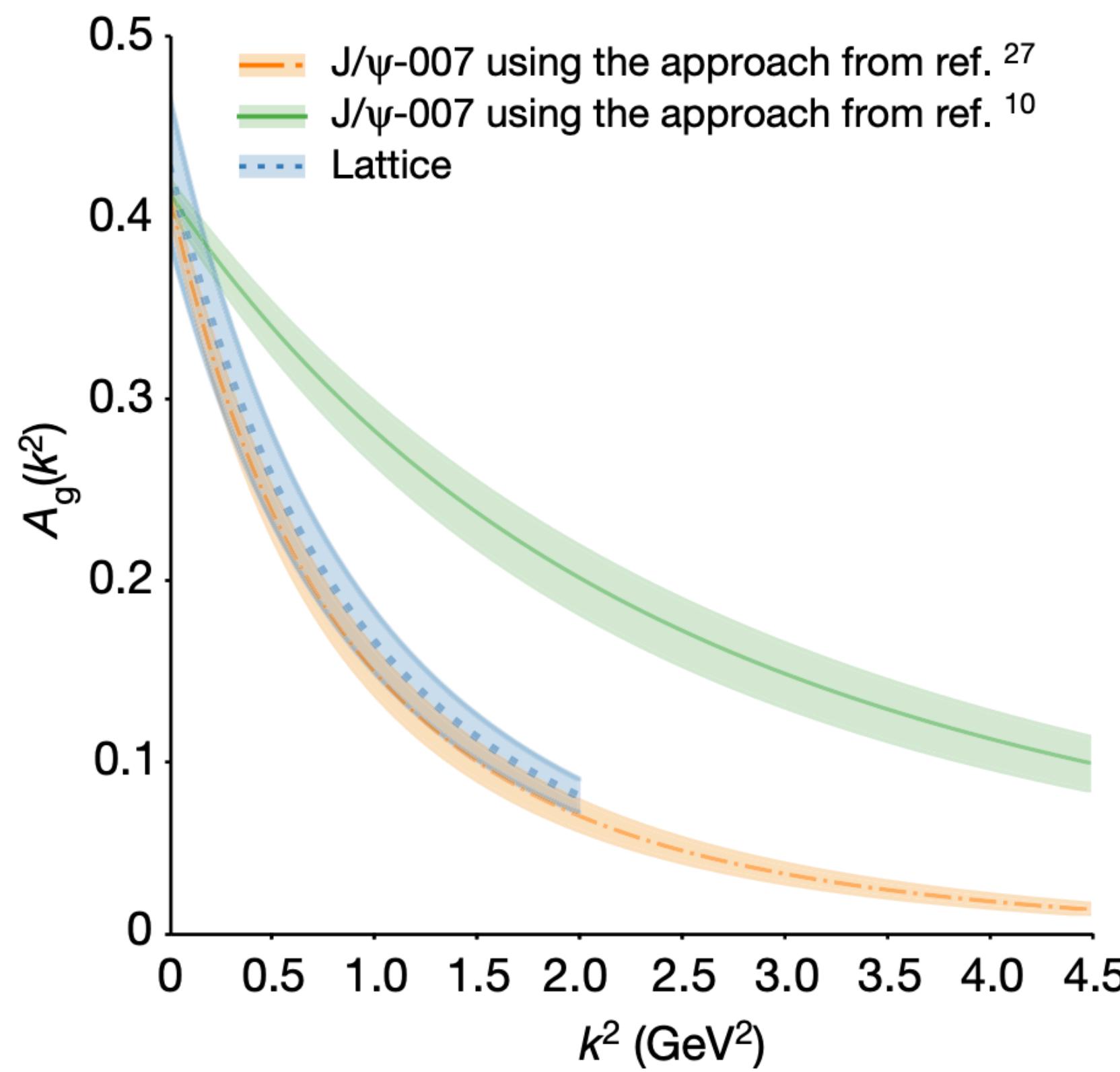
Y. Guo, X. Ji, Y. Liu, PRD 103, 096010 (2021)

$$\frac{d\sigma}{dt} = \frac{\alpha_e m e_Q^2}{4(W^2 - m_N^2)^2} \frac{(16\pi\alpha_s)^2}{3M_V^2} |\psi_{NR}|^2 |G(t, \xi)|^2$$

Assume  $\xi \sim 1$  (it is less than 0.5 for most of the experimental data)

# GLUONIC GFFS FROM EXPERIMENTAL DATA

**Remarkable agreement between GFFs determined from data using the Holographic QCD approach and the direct Lattice QCD calculation!**



**Determined from experiment**  
 J/ψ-007 using the approach from ref. <sup>27</sup>  
 J/ψ-007 using the approach from ref. <sup>10</sup>  
 Lattice

**Determined from theory**  
 Holographic QCD approach  
 GPD approach  
 Lattice QCD calculation

Ref 27 (Holographic QCD): K. Mamo & I. Zahed, PRD 103, 094010 (2021) and 2204.08857 (2022)  
 Ref 10 (GPD Formalism): Y. Guo, X. Ji, Y. Liu, PRD 103, 096010 (2021)  
 Lattice: D. Pefkou, D. Hackett, P. Shanahan, Phys. Rev. D 105, 054509 (2022).

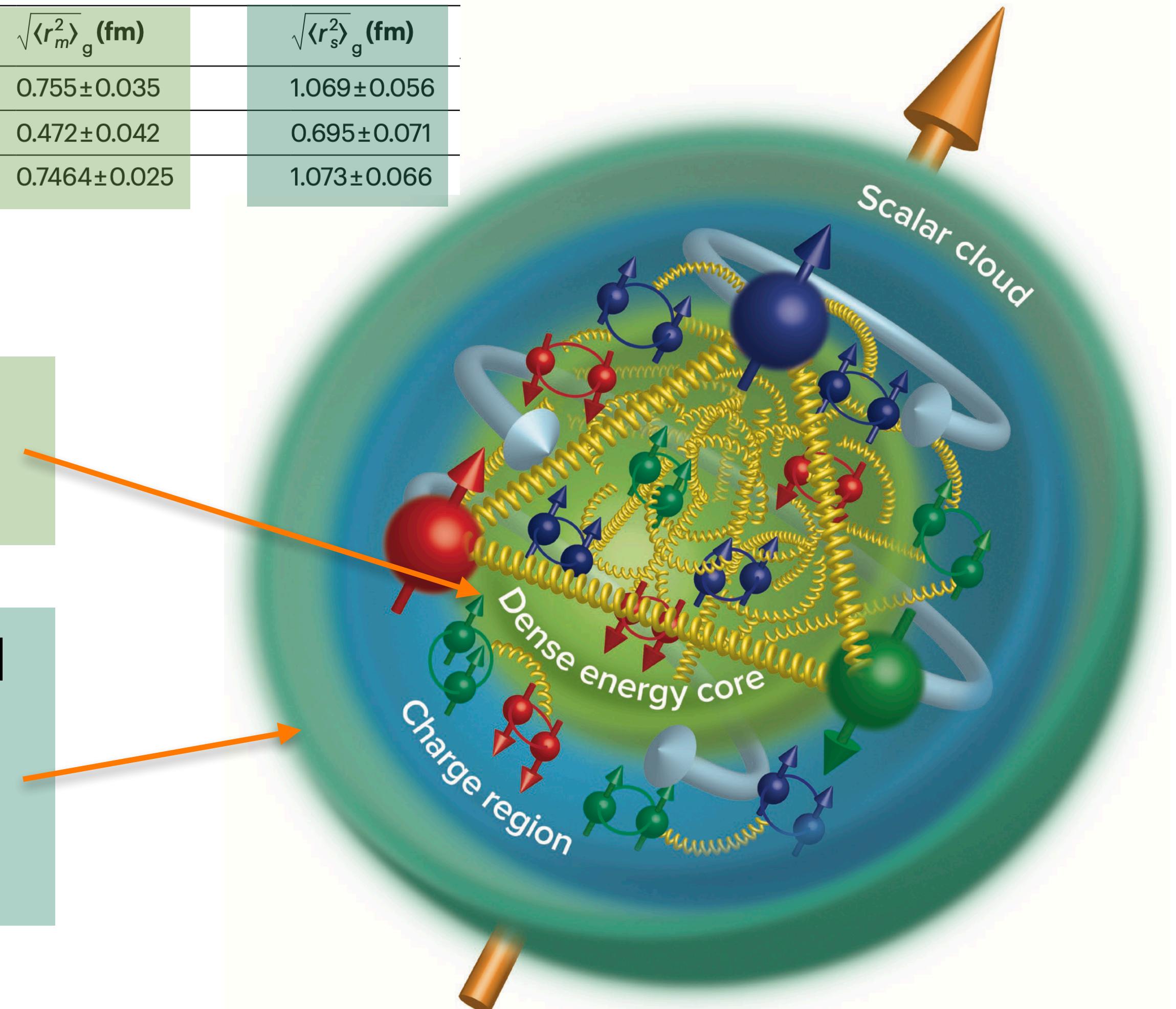
# THE PROTON IN THREE REGIONS

**Table 1 | The gluonic GFF fit parameters, proton mass radius and scalar radius**

Theoretical approach	$\chi^2/\text{n.d.f.}$	$m_A (\text{GeV})$	$m_c (\text{GeV})$	$C_g(0)$	$\sqrt{\langle r_m^2 \rangle_g} (\text{fm})$	$\sqrt{\langle r_s^2 \rangle_g} (\text{fm})$
Holographic QCD	0.925	$1.575 \pm 0.059$	$1.12 \pm 0.21$	$-0.45 \pm 0.132$	$0.755 \pm 0.035$	$1.069 \pm 0.056$
GPD	0.924	$2.71 \pm 0.19$	$1.28 \pm 0.5$	$-0.20 \pm 0.11$	$0.472 \pm 0.042$	$0.695 \pm 0.071$
Lattice		$1.641 \pm 0.043$	$1.07 \pm 0.12$	$-0.483 \pm 0.133$	$0.7464 \pm 0.025$	$1.073 \pm 0.066$

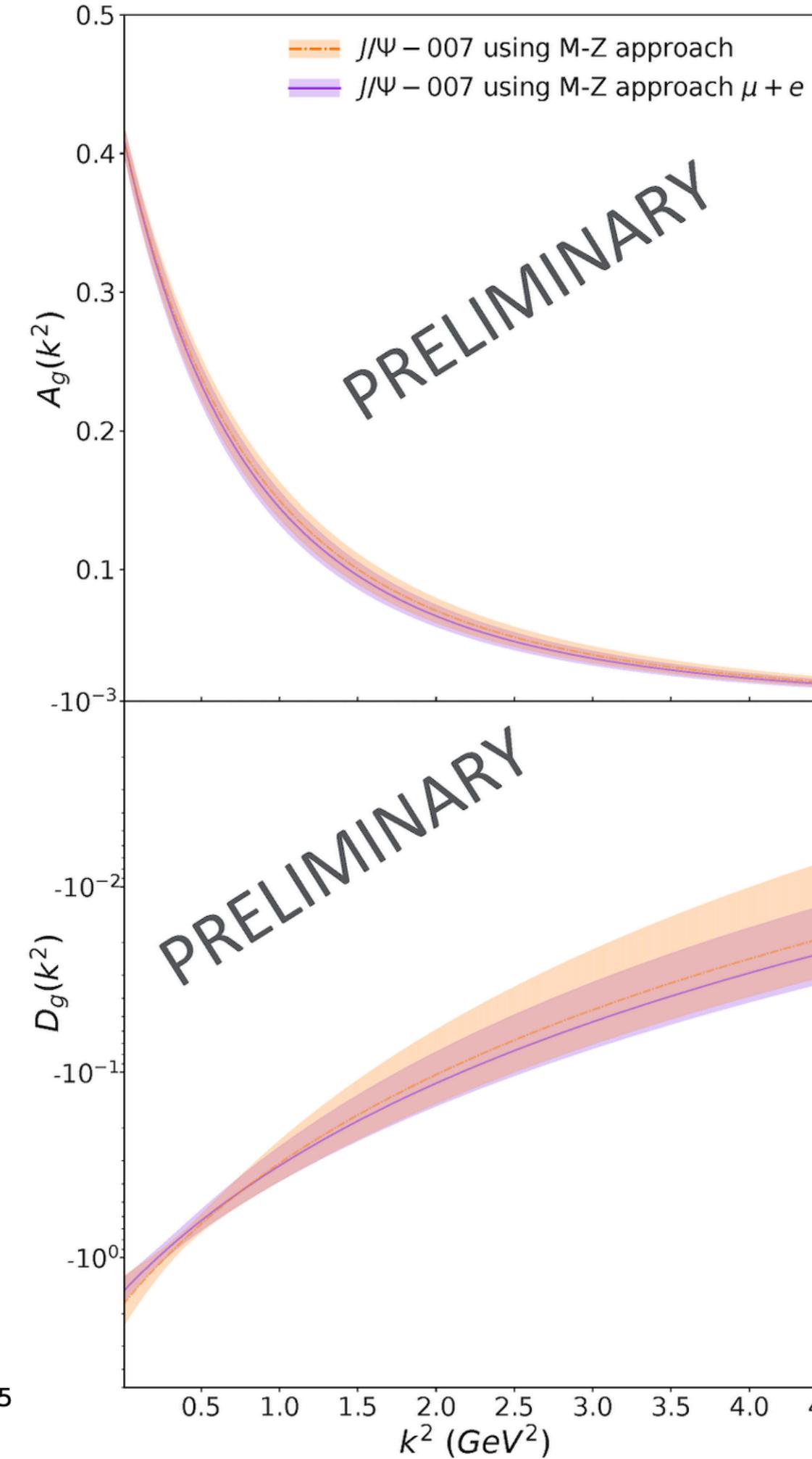
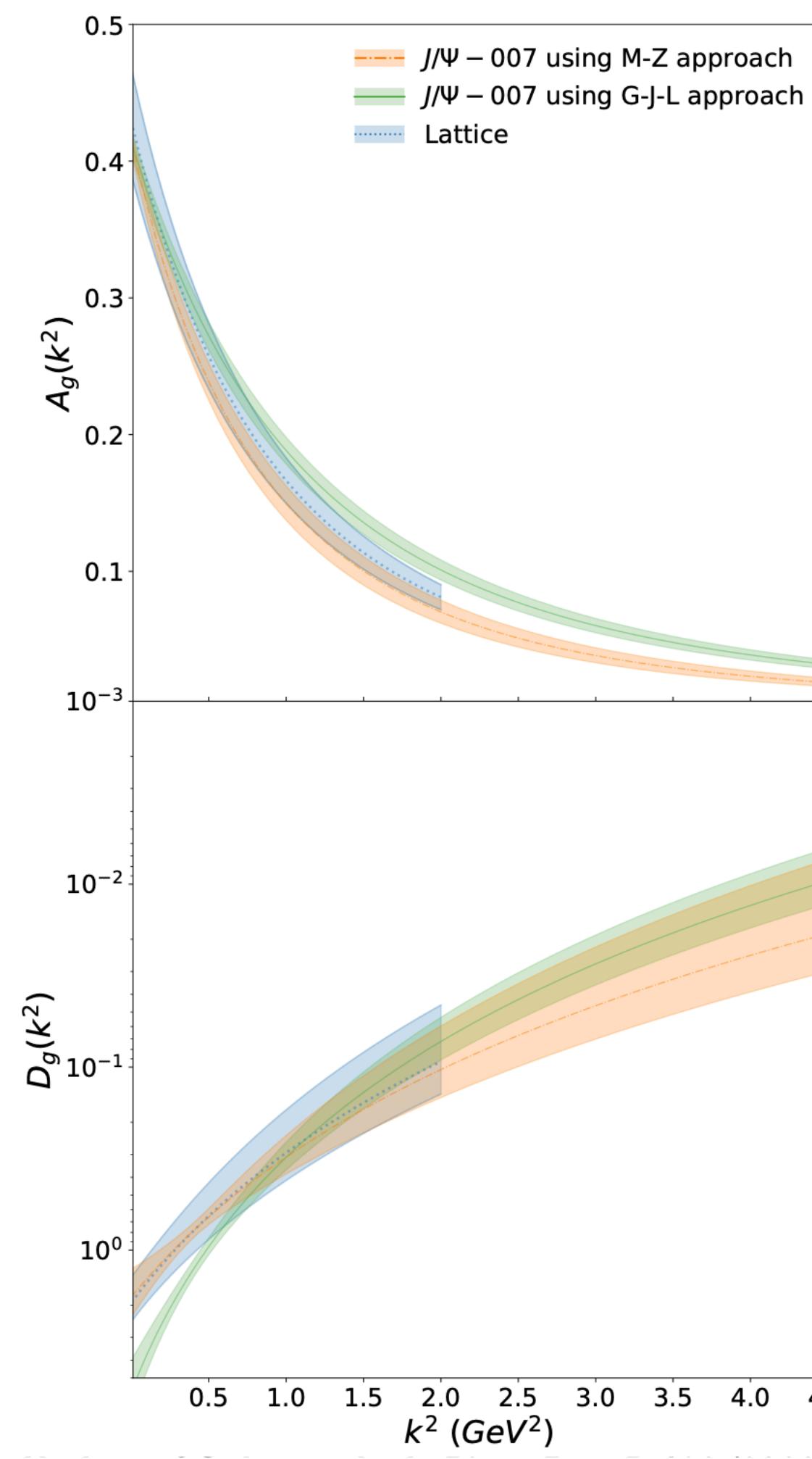
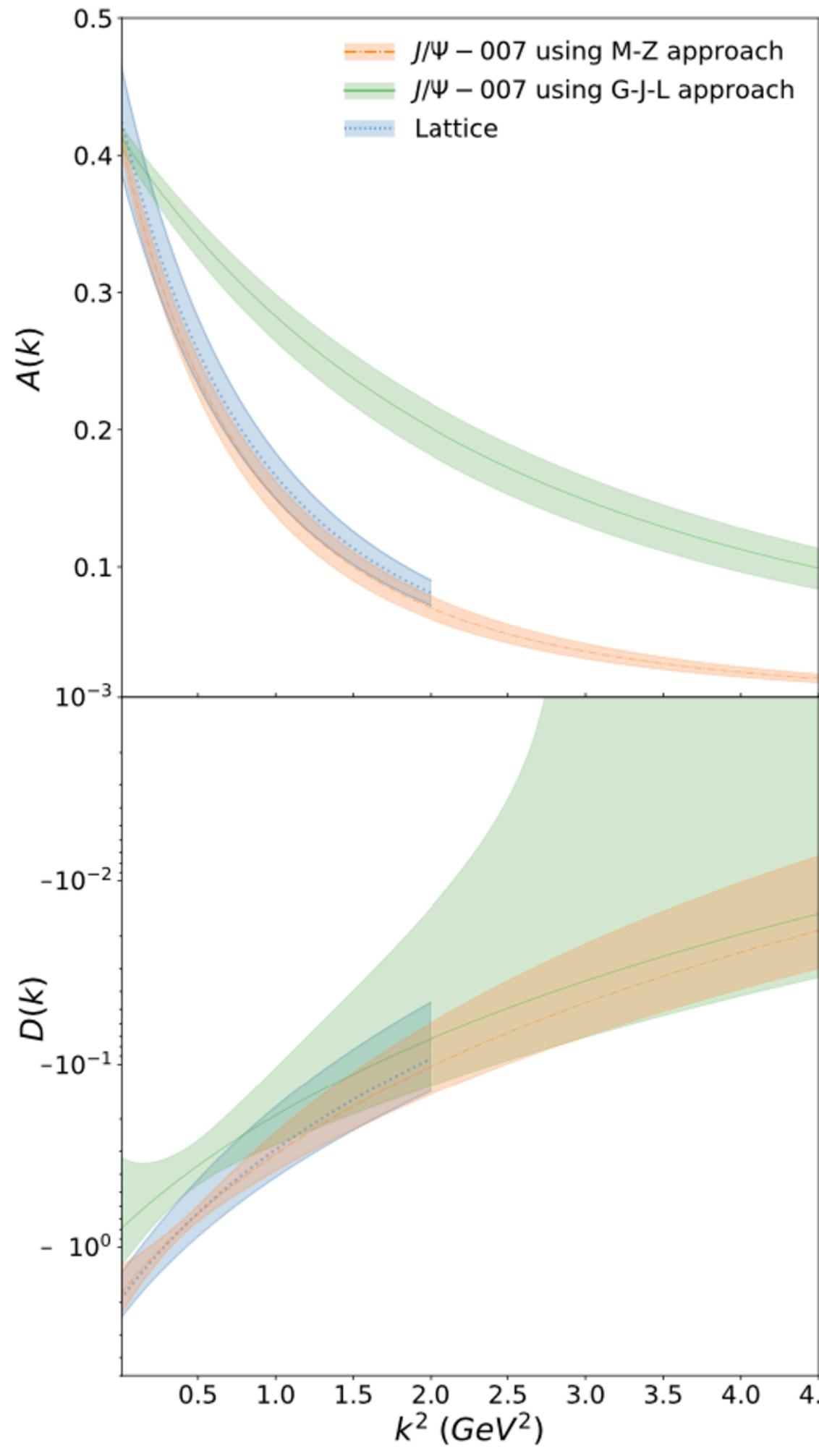
The proton's mass radius seems substantially smaller than its charge

The holographic QCD fit to our data and the latest Lattice calculations find a scalar gluonic cloud surrounding the charge region at about 1 fermi



# UPDATES SINCE THE J/ $\psi$ -007 PAPER

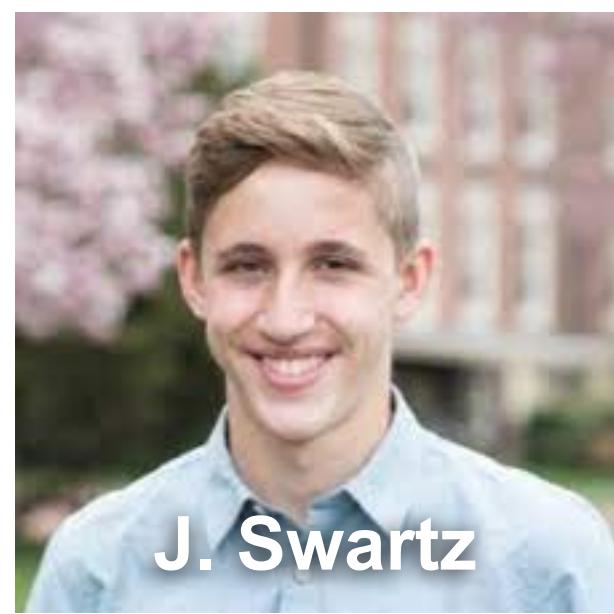
Update of G-J-L analysis PRD 108  
 (2023) no.3, 034003 arXiv:2305.06992



- Analysis with the muons decay channel results , doubling the statistics

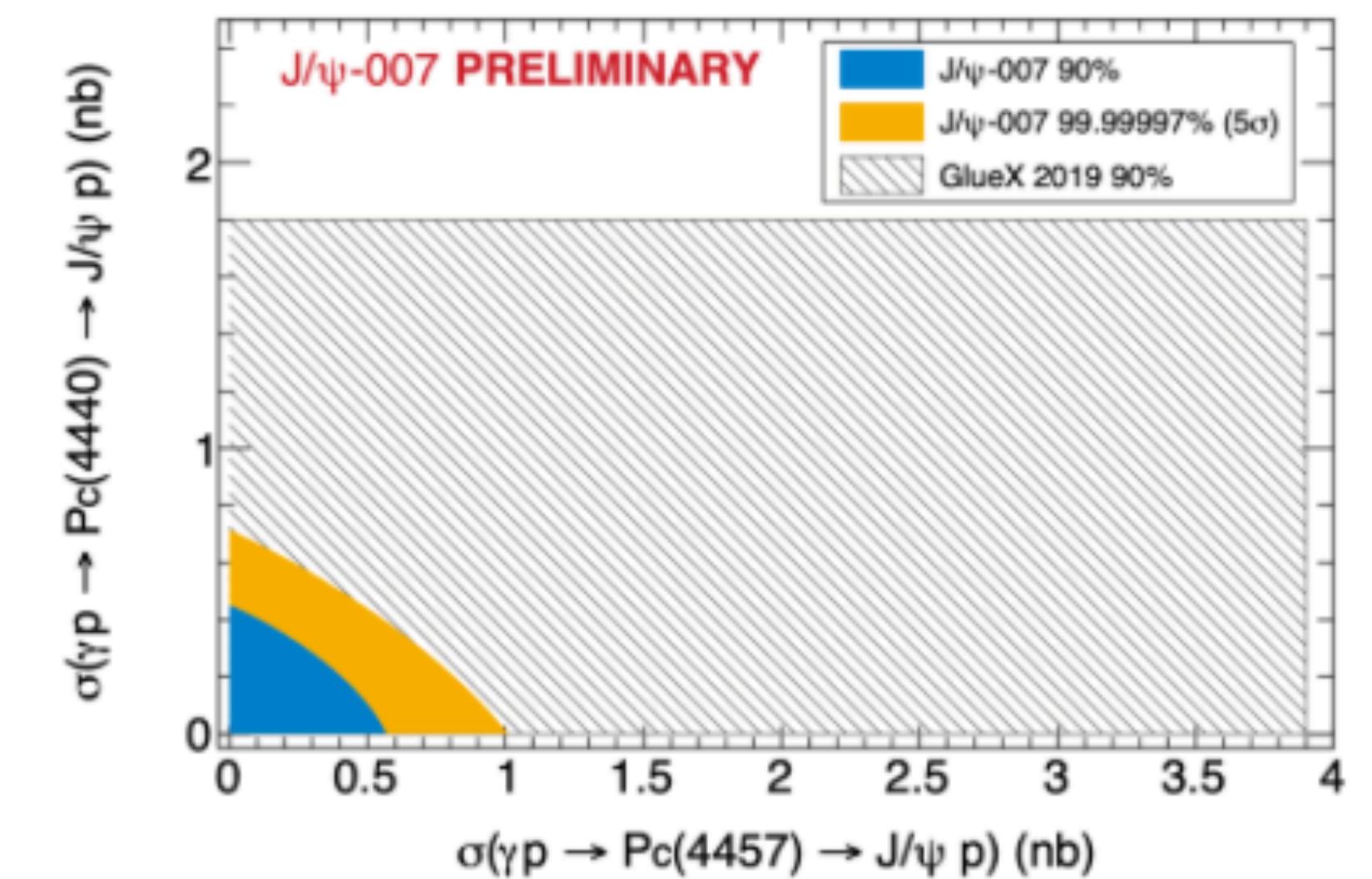
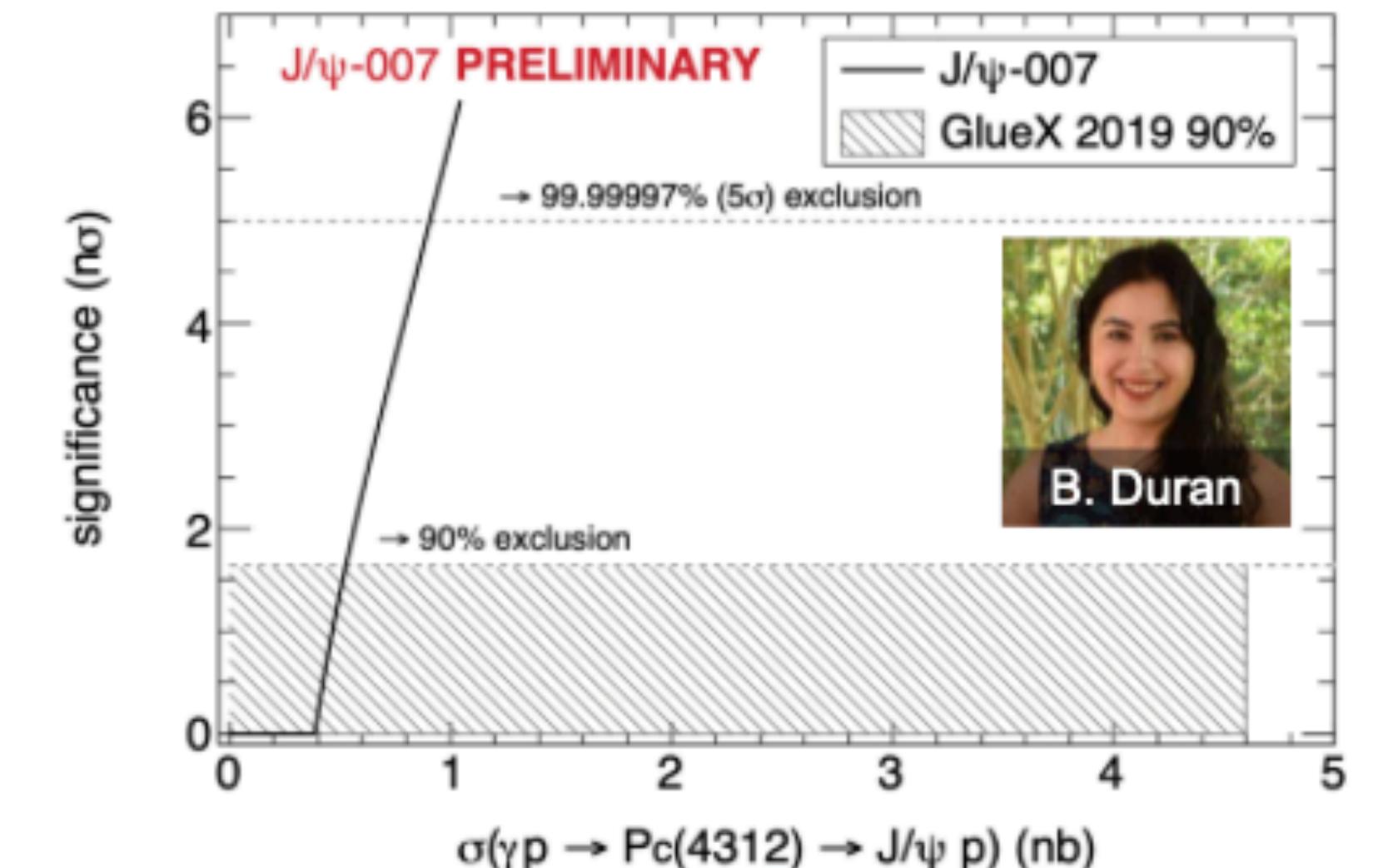
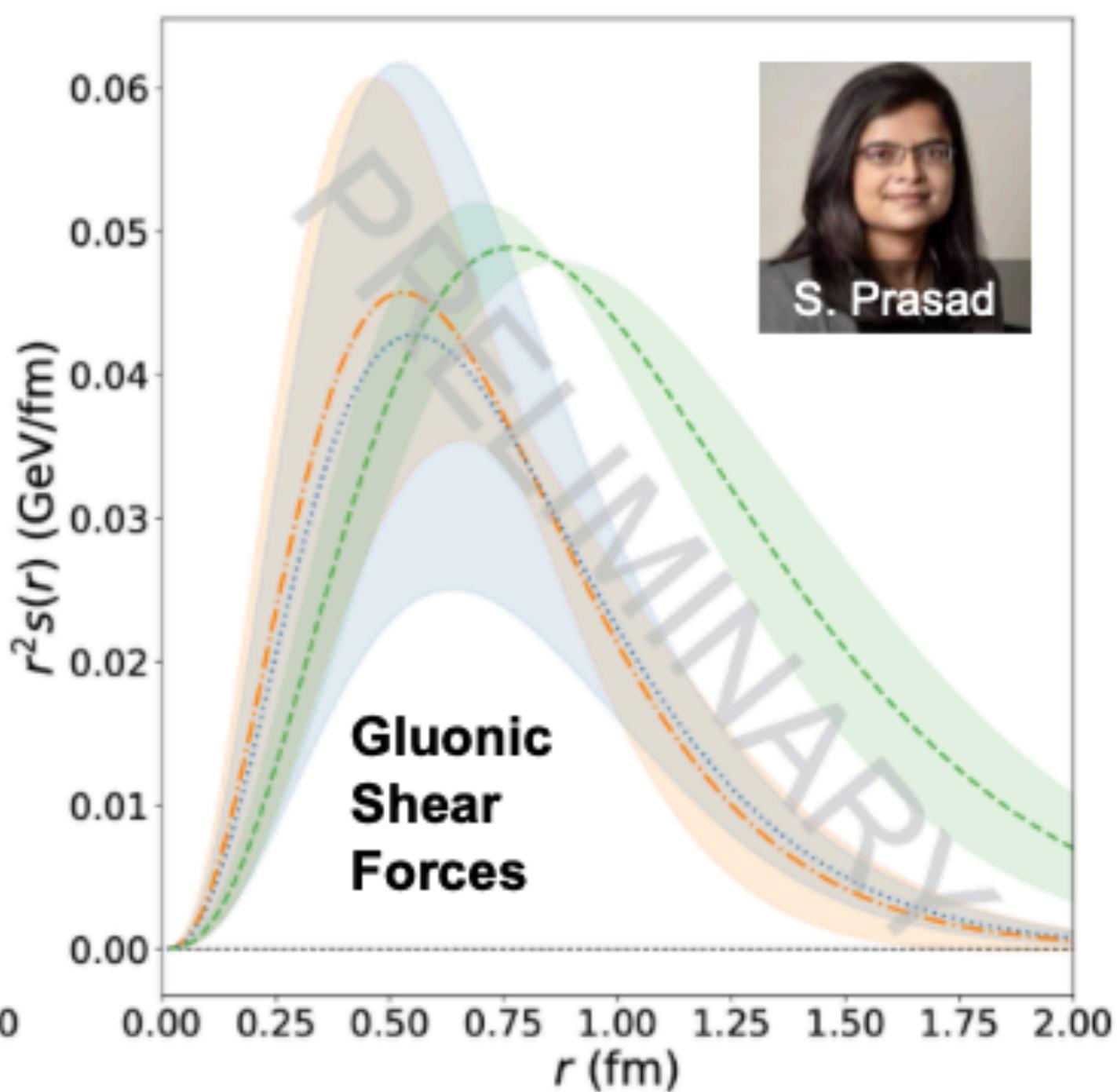
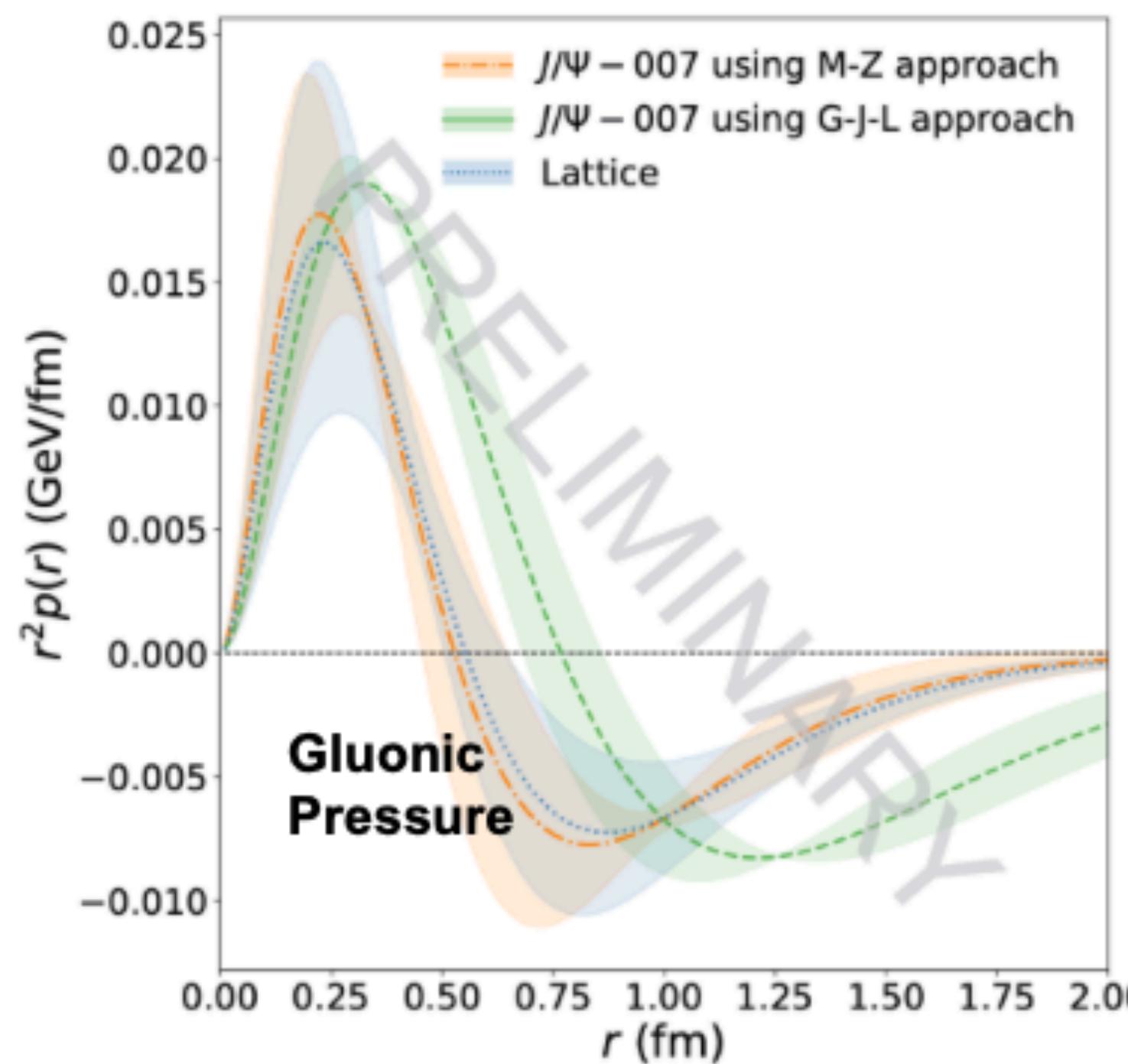
- Consistent with the electron results.

- Largest impact on the  $C(t)$  form factor with improved precision



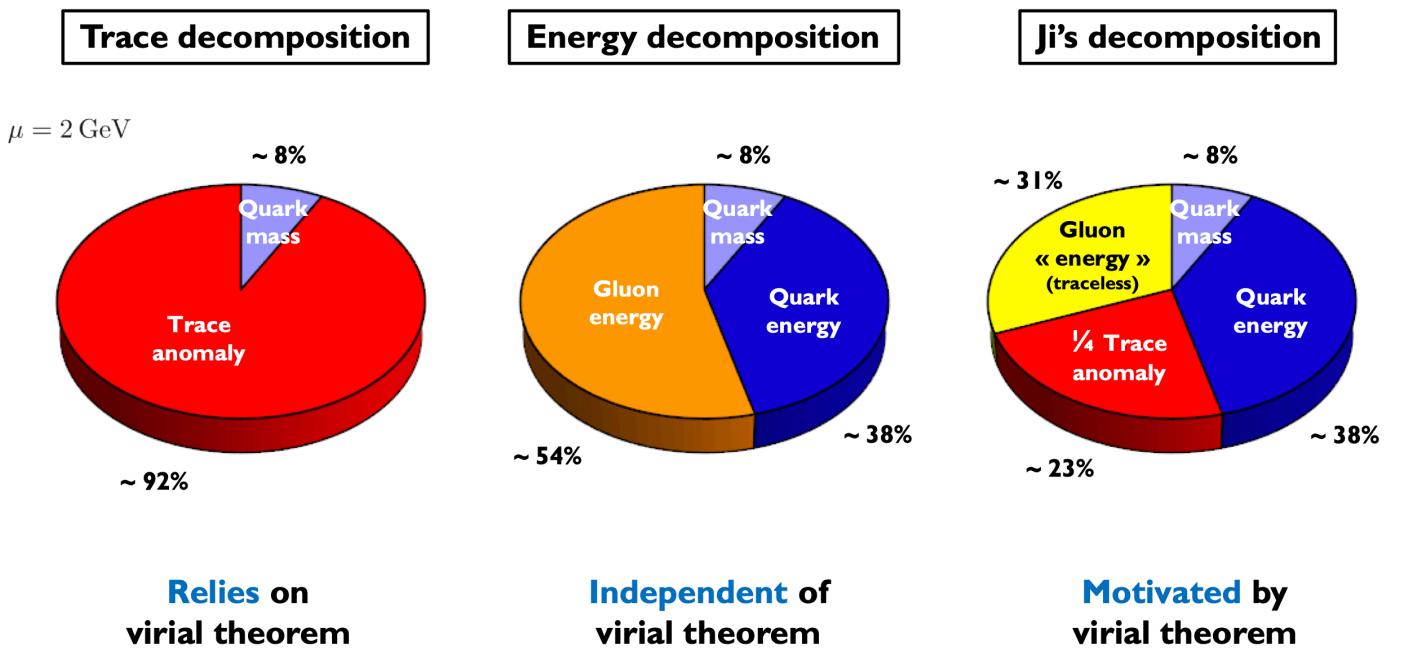
# Upcoming J/ψ-007 Results

## From near-threshold J/ψ in 2-D to gluonic gravitational form factors

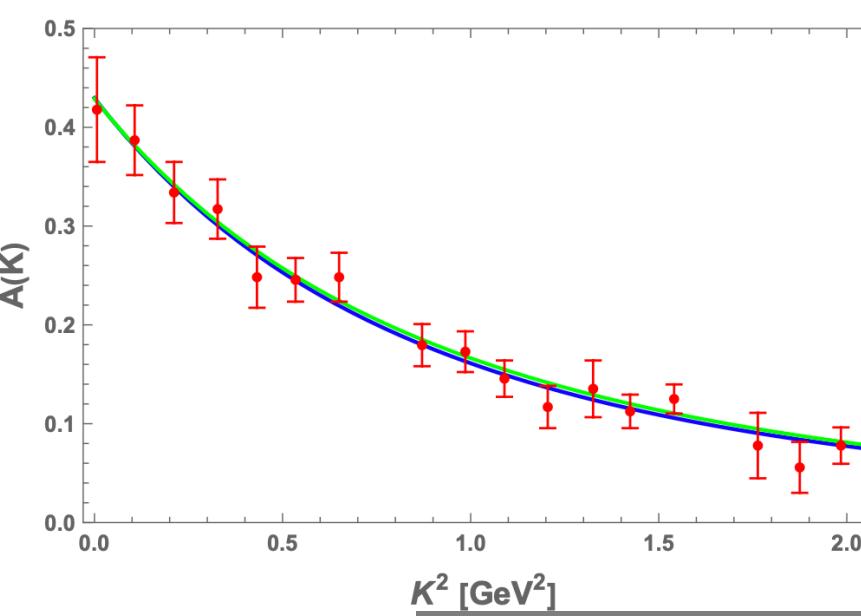


**Expect more results soon:** Finalizing systematics on the di-muon decay channel, preparing two new manuscripts: on the 2D cross section, and on the LHCb hidden-charm pentaquark

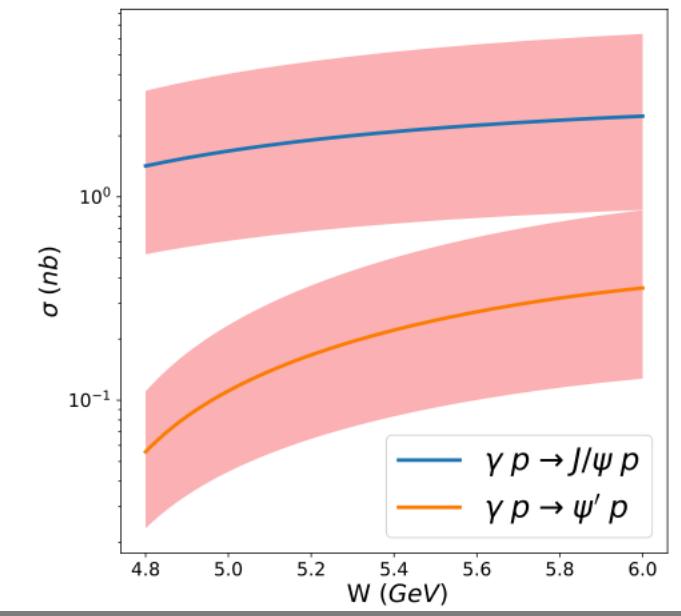
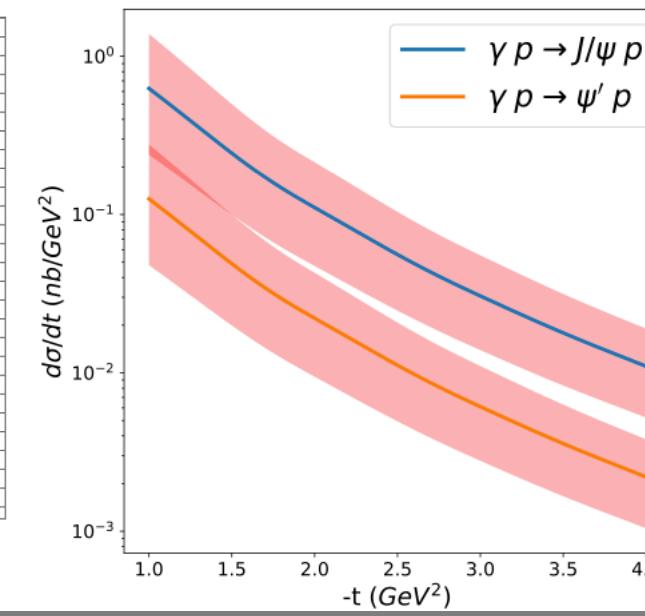
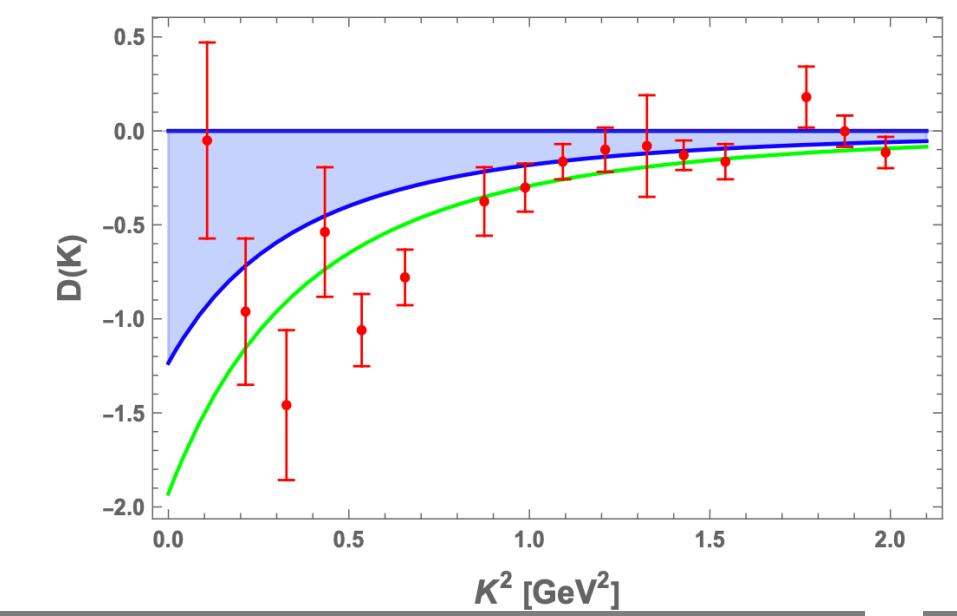
# PROMINENT RECENT DEVELOPMENTS



Proton mass budget decompositions,  
C. Lorce (from 2022 INT workshop)

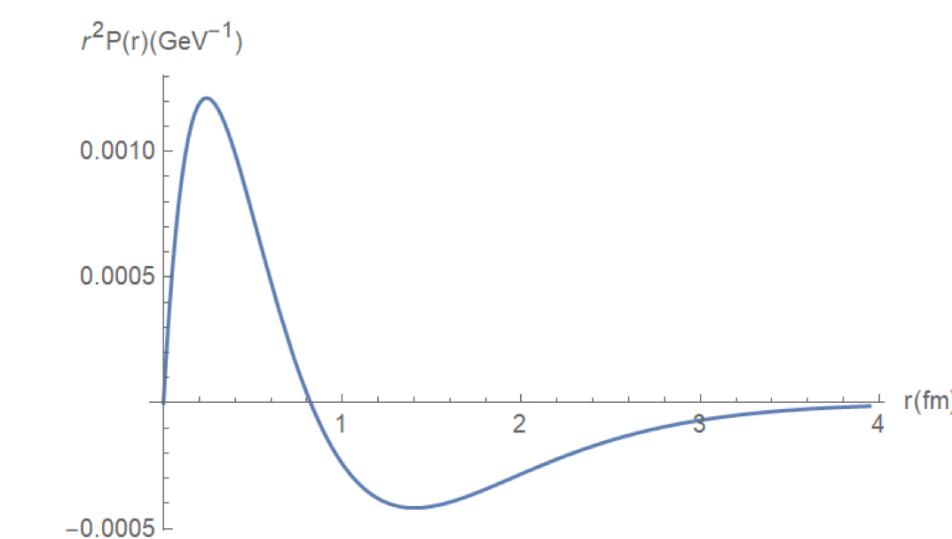


Proton gravitational form factors  
holographic QCD compared with  
Lattice, K. Mamo & I. Zahed (2022)

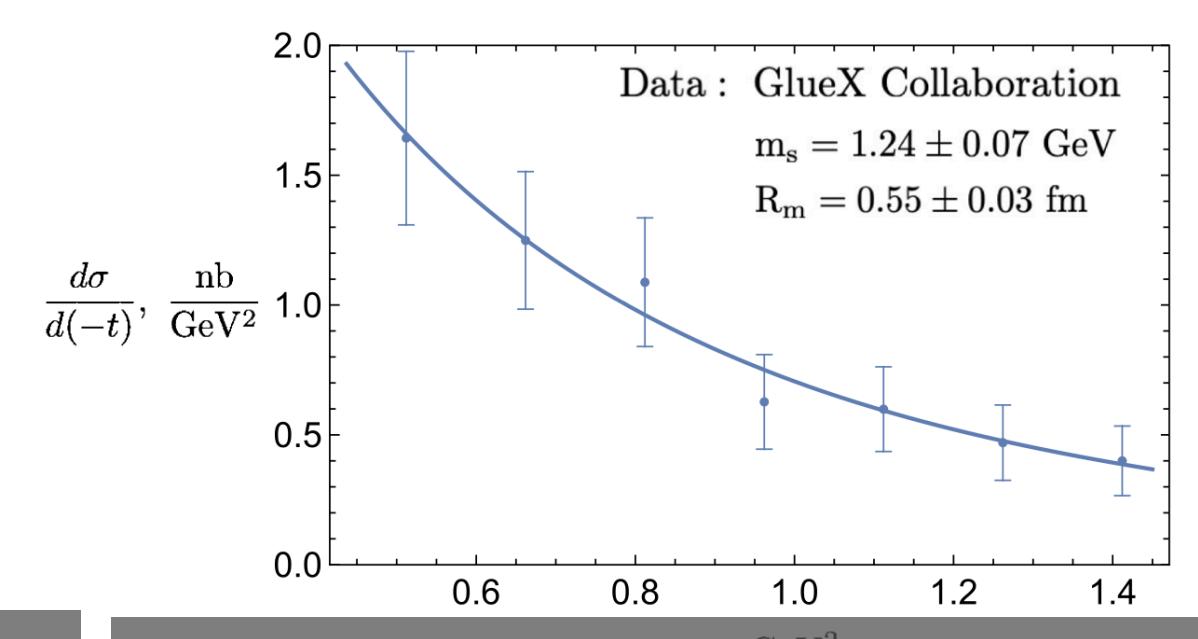


Near-threshold heavy quarkonium  
production at large momentum transfer,  
P. Sun, X-B. Tong, F. Yuan (PRD 2022)

- A hot topic: many theoretical developments, and pace of publications only speeding up!
- Many extractions depend on extrapolating to the forward limit ( $t=0$ ), which introduces theoretical systematic uncertainties. Precise high- $t$  as a function photon energy crucial.
- Other avenues for factorization include large- $t$  region, large  $Q^2$  region, or larger vector meson mass.



Gluon contribution to pressure  
in GPD formalism, Y. Guo, X. Ji,  
Y. Liu, (PRD 2021)

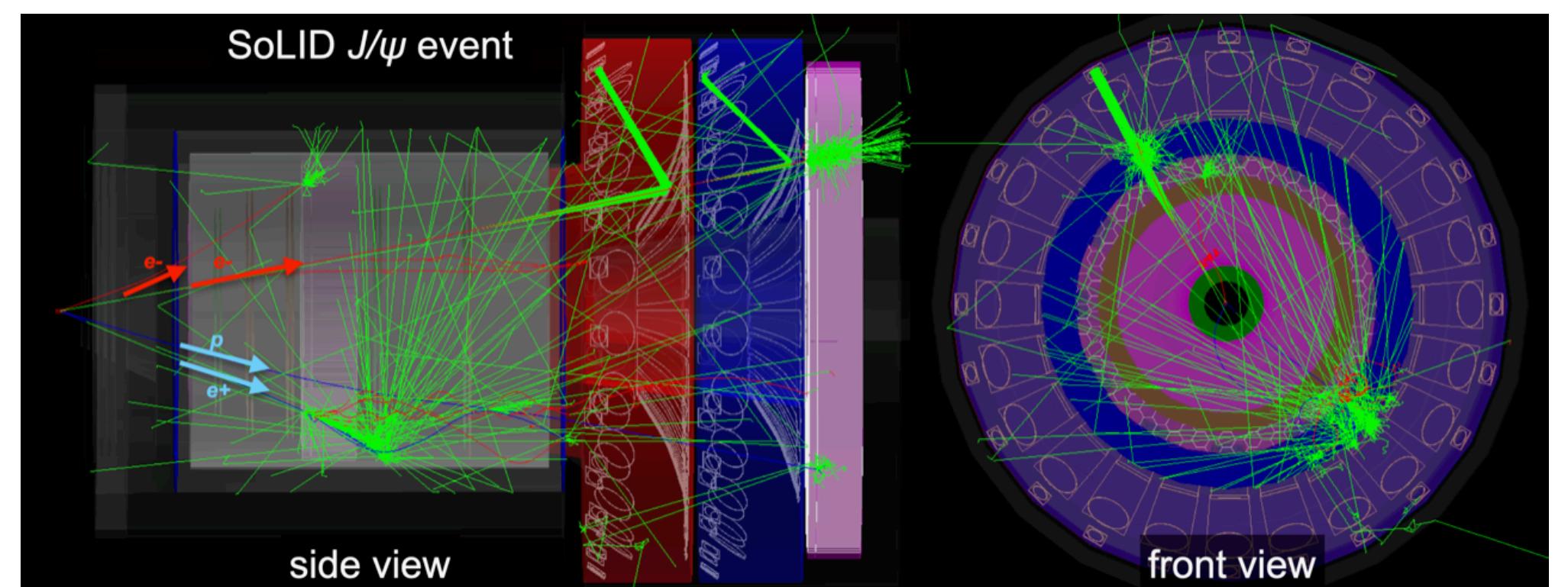
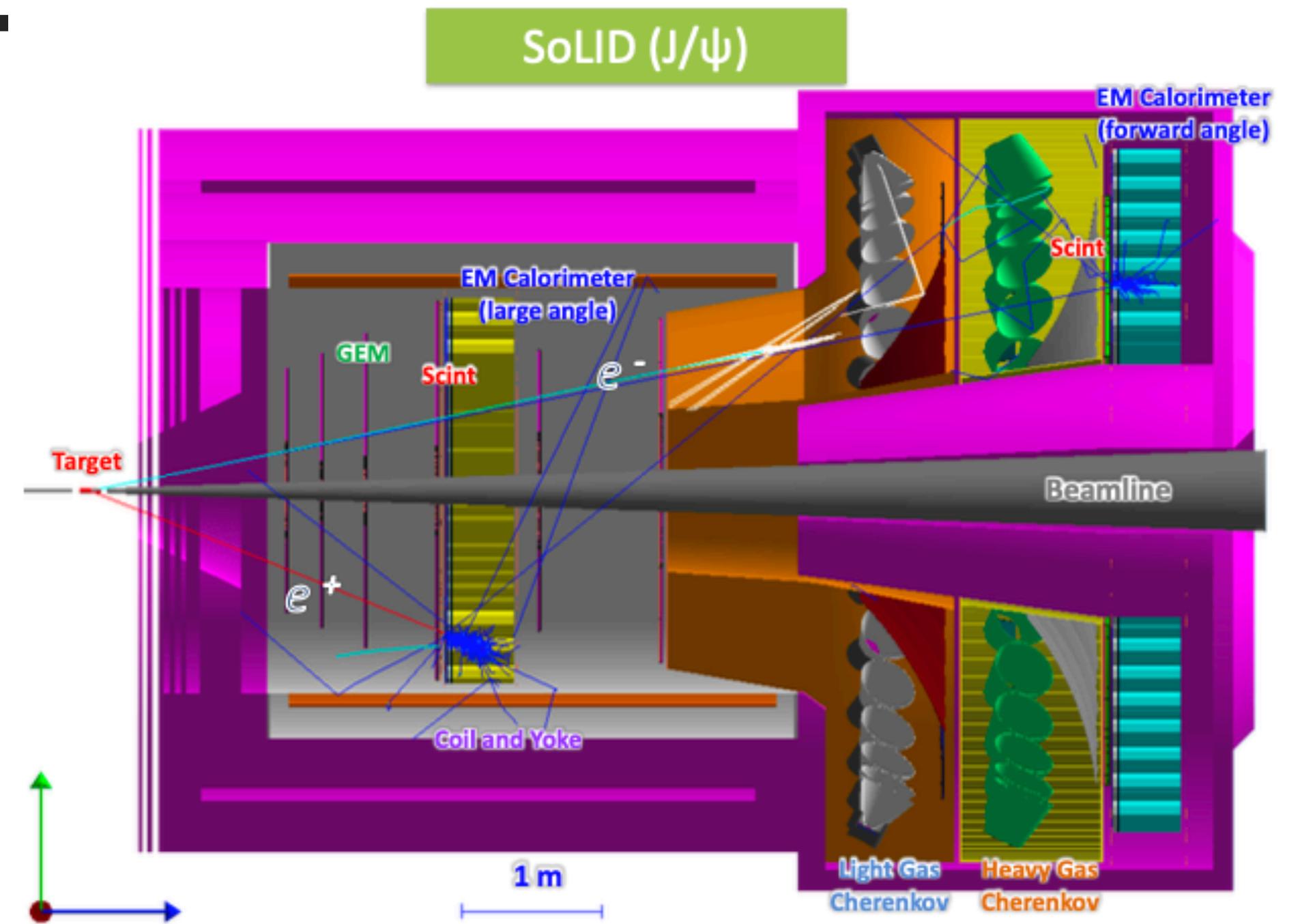


Gluonic radius of the proton  
based on 1D GlueX results, D.  
Kharzeev (PRD 2021)

# THE SoLID-J/ $\psi$ EXPERIMENT

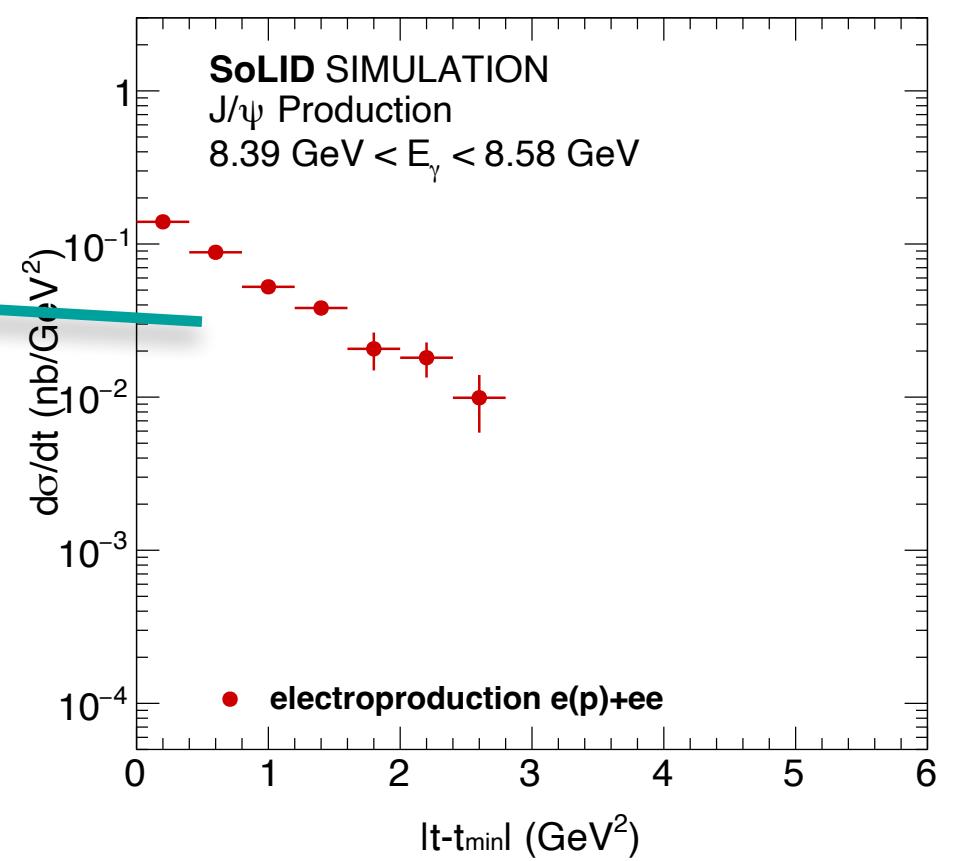
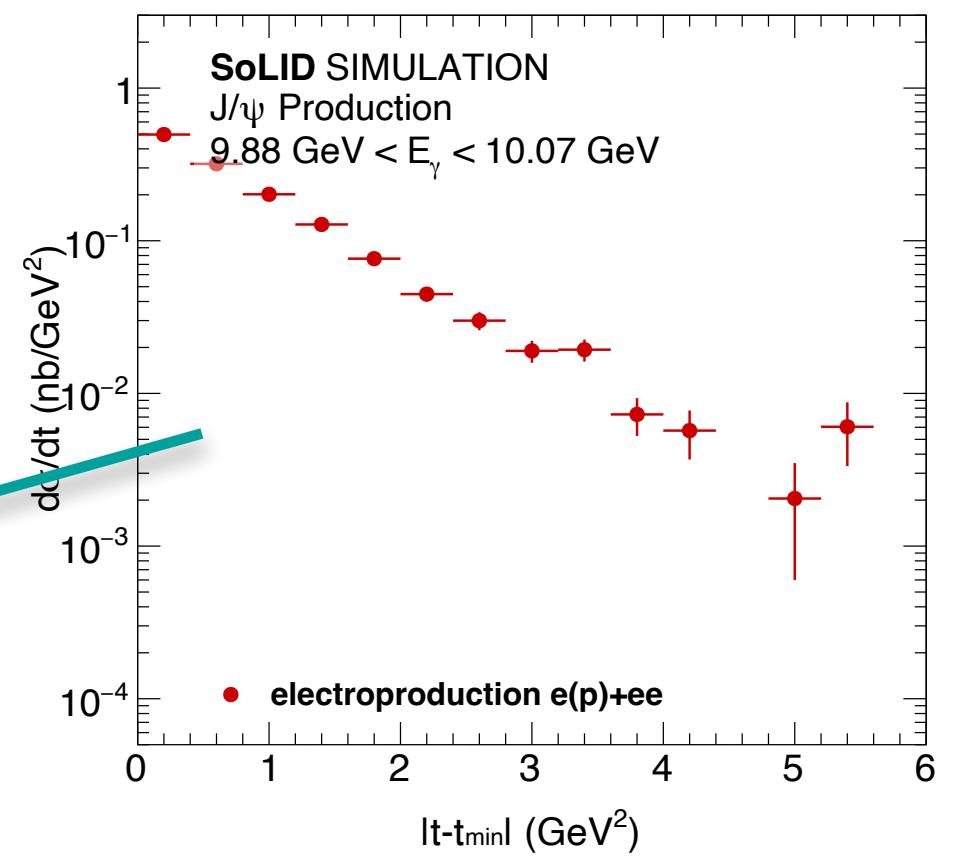
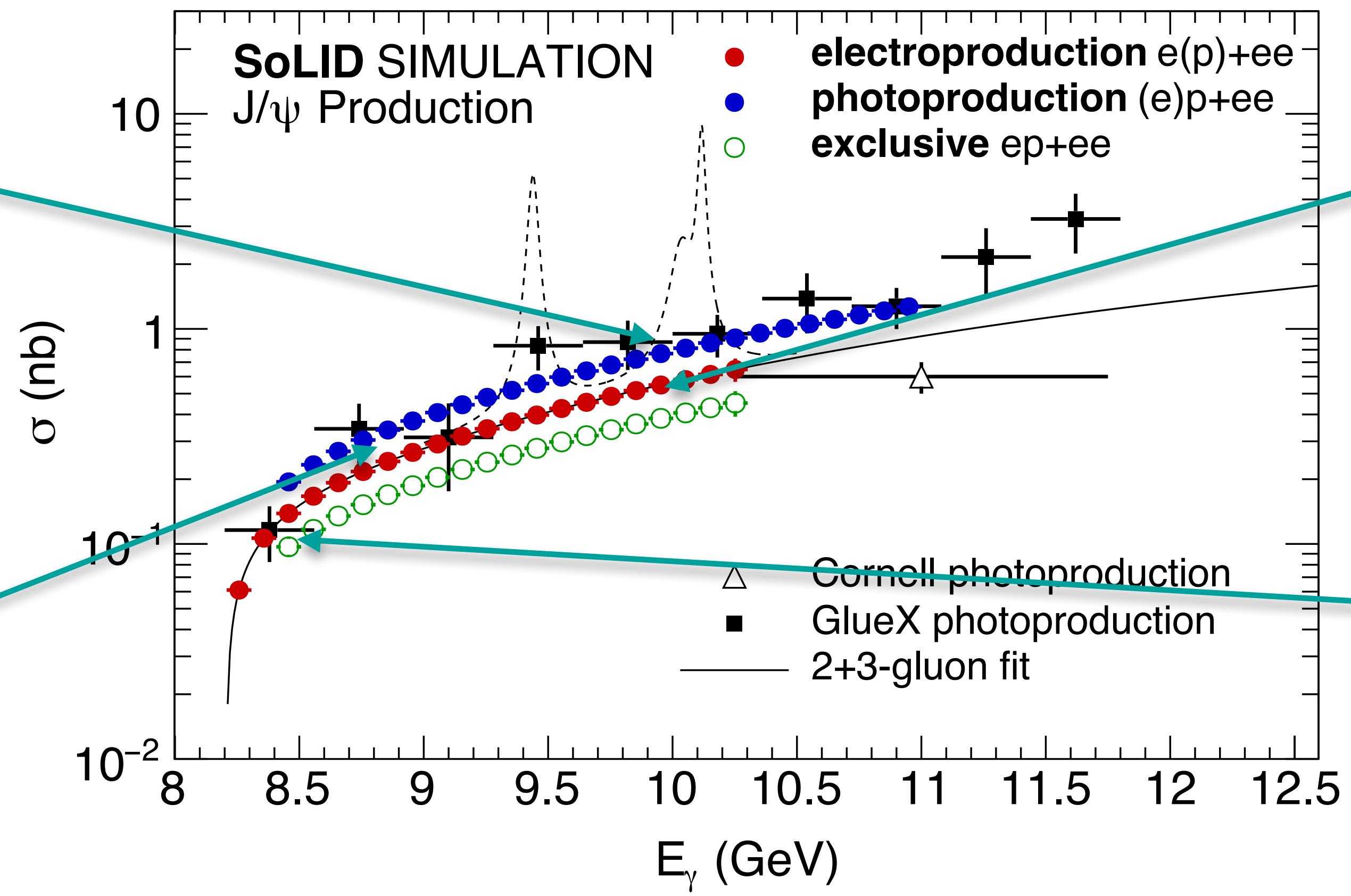
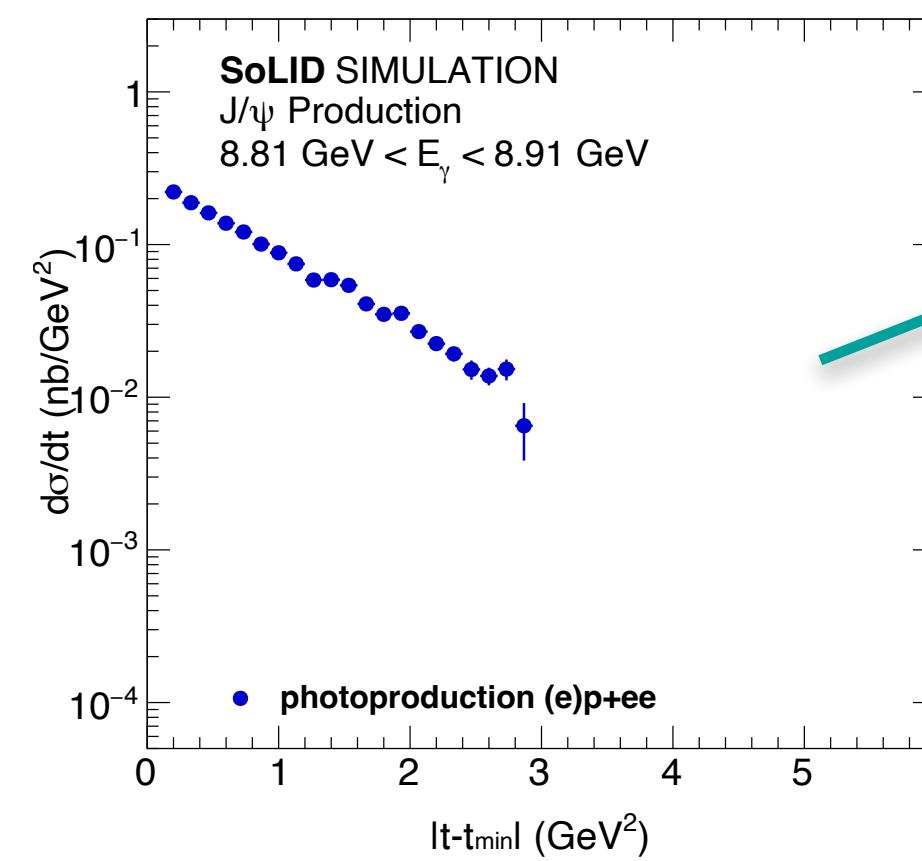
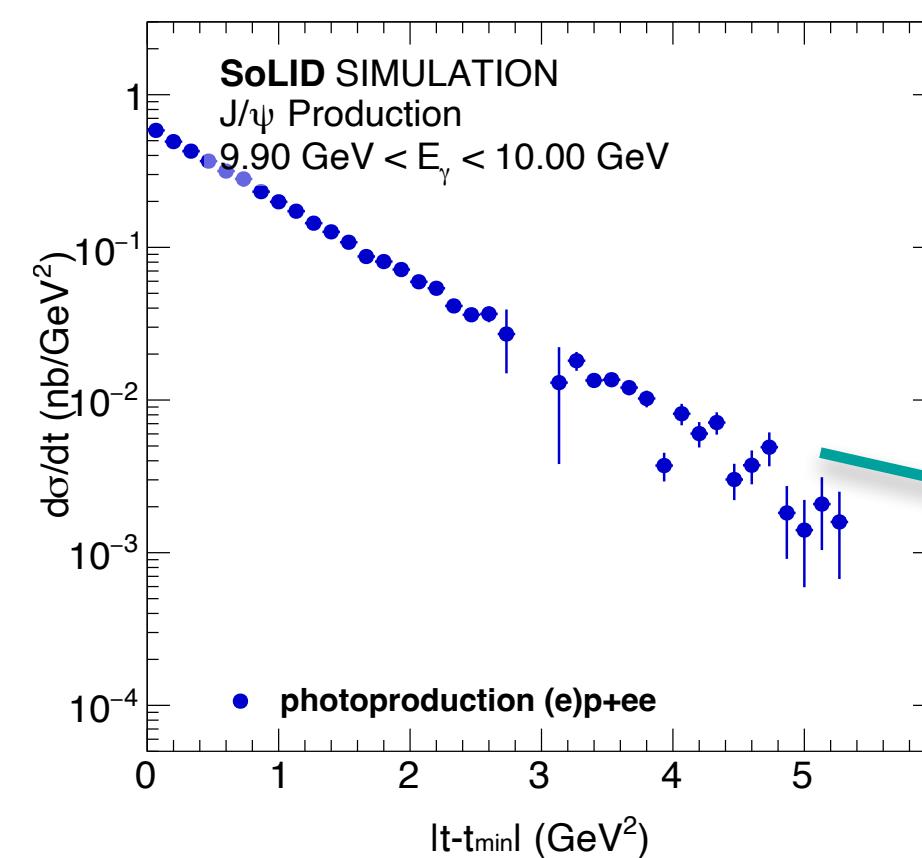
## Ultimate factory for near-threshold J/ $\psi$

- General purpose large-acceptance spectrometer
- 50+10 days of  $3\mu\text{A}$  beam on a 15cm long LH<sub>2</sub> target ( $10^{37}/\text{cm}^2/\text{s}$ )
- **Ultra-high luminosity:**  $43.2\text{ab}^{-1}$
- **Open 2-particle trigger**, covering J/ $\psi$  production in four channels:  
Electroproduction ( $e^-e^-e^+$ ), photoproduction ( $p,e^-e^+$ ),  
inclusive ( $e^-e^+$ ), exclusive ( $ep,e^-e^+$ )
- The electroproduction channel provides for a modest lever-arm in  $Q^2$  near threshold



# SOLID-J/ $\psi$ PROJECTIONS

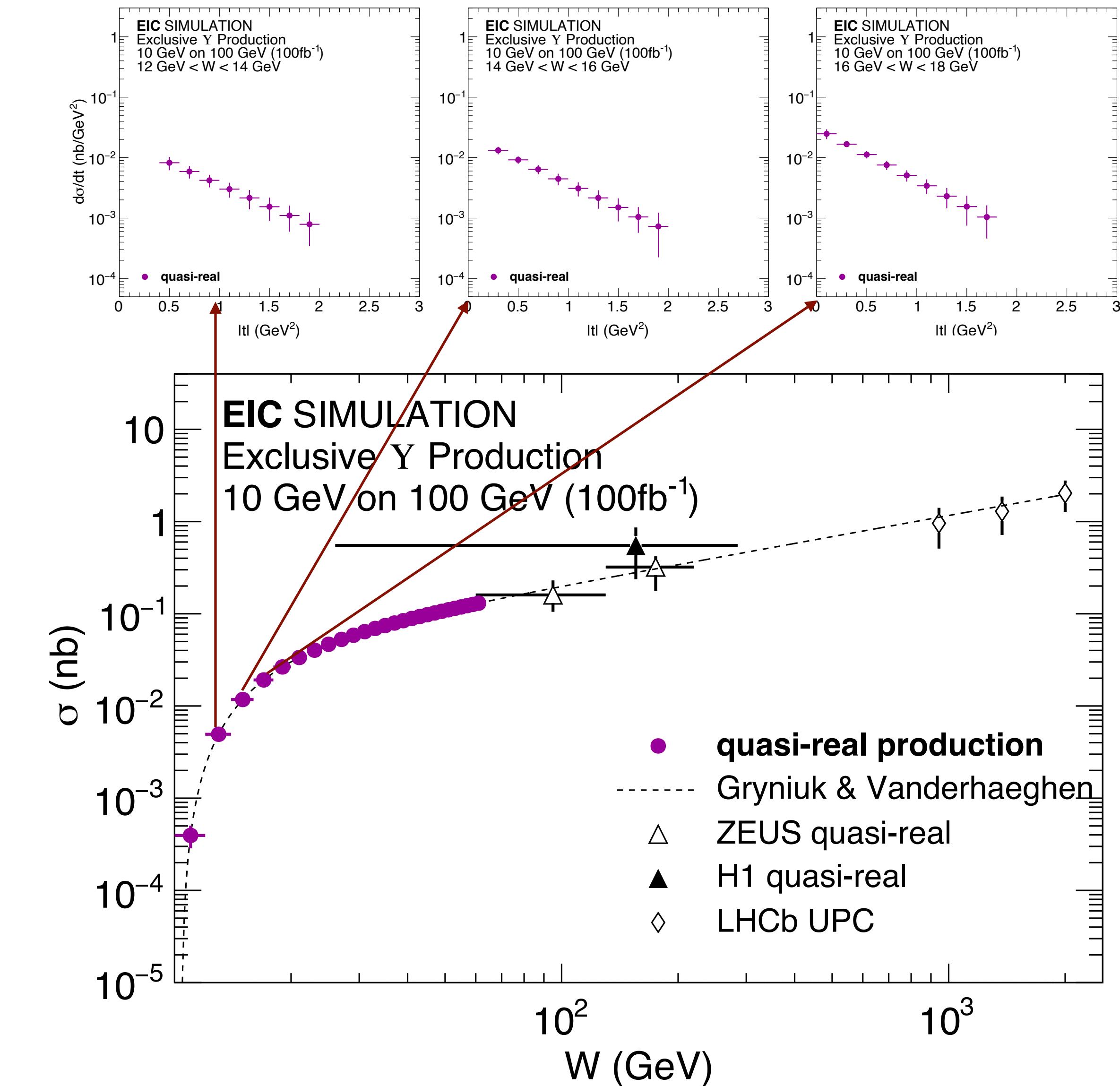
High-precision 2-D cross section crucial to really connect GFFs to data



# $\Upsilon(1S)$ NEAR THRESHOLD

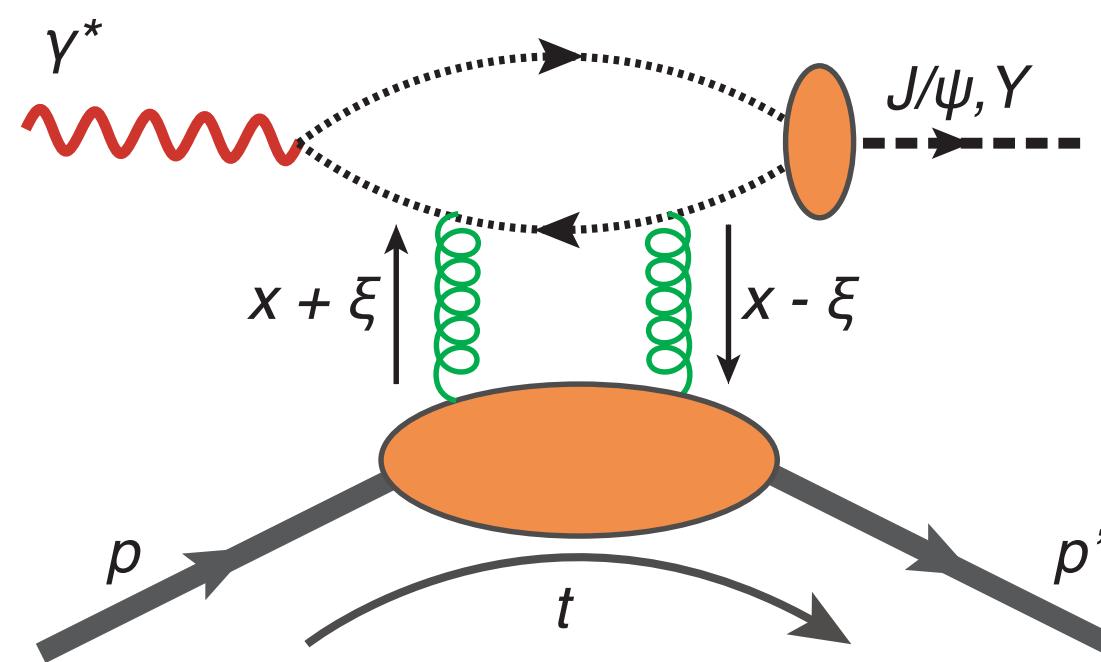
## $\Upsilon(1S)$ with the ePIC detector

- $\Upsilon(1S)$  at EIC trades statistical precision of  $J/\psi$  at SoLID for lower theoretical uncertainties and extra channel to study universality.
- Large  $Q^2$  reach at EIC an additional knob to study production



# DEEPLY-VIRTUAL QUARKONIUM PRODUCTION

## Accessing the 3-D gluon structure

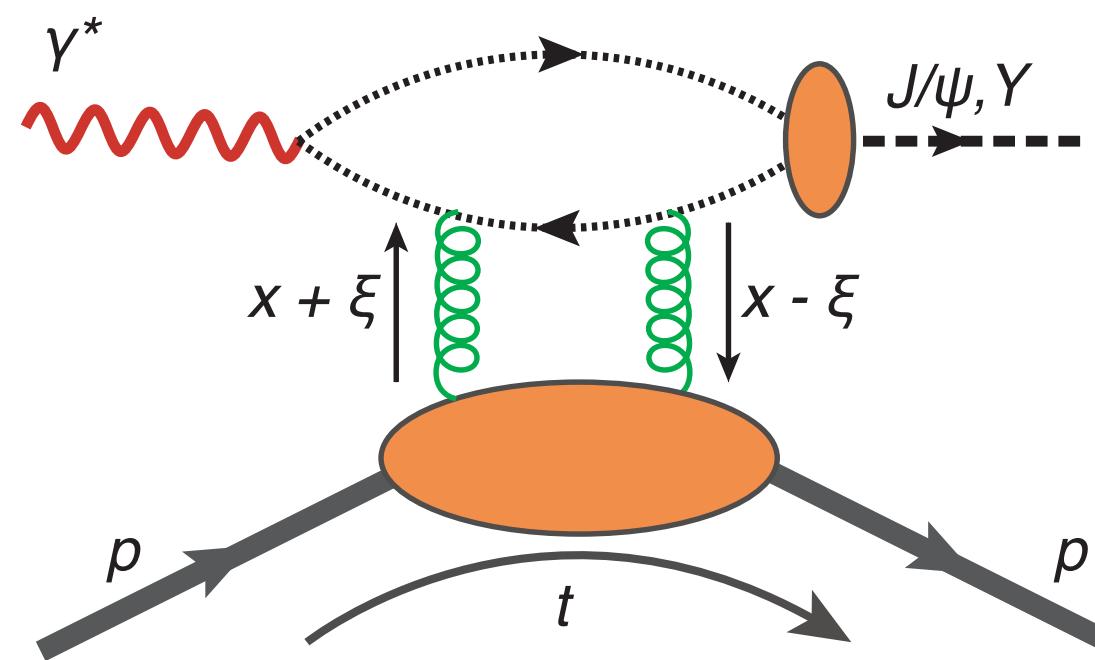


**Hard scale:**  $Q^2 + M_V^2$

**Modified Bjorken- $x$ :** 
$$x_V = \frac{Q^2 + M_V^2}{2p \cdot q}$$

# DEEPLY-VIRTUAL QUARKONIUM PRODUCTION

## Accessing the 3-D gluon structure



**average unpolarized gluon GPD related to  
 $t$ -dependent cross section (LO)**

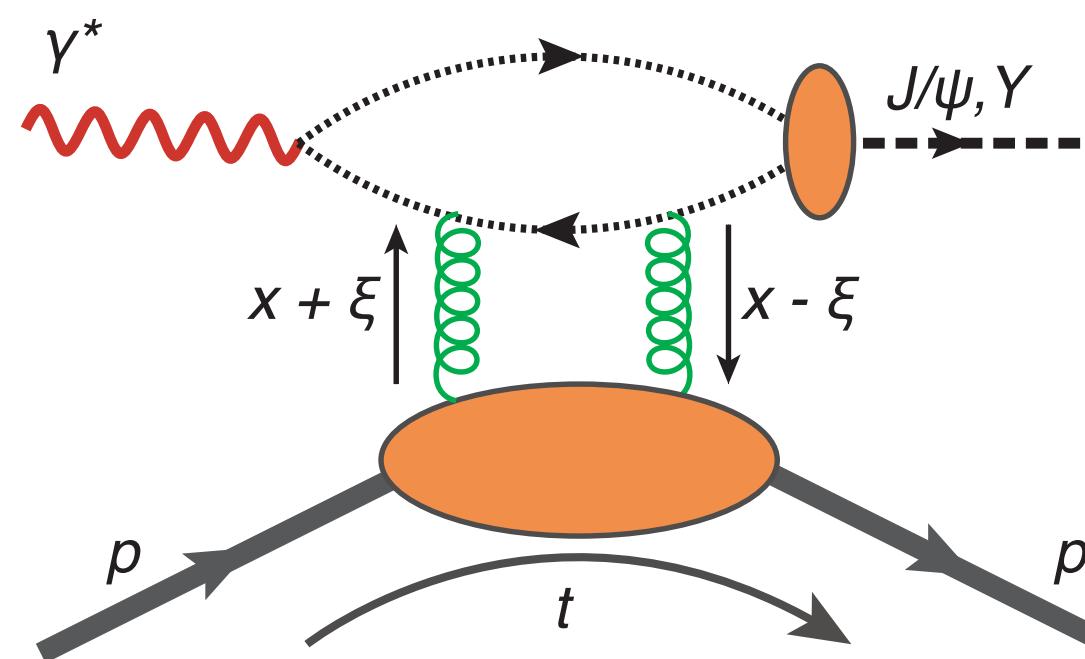
$$|\langle \mathcal{H}_g \rangle|(t) \propto \sqrt{\frac{d\sigma}{dt}(t)/\frac{d\sigma}{dt}(t=0)}$$

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**Modified Bjorken-x:**  $x_V = \frac{Q^2 + M_V^2}{2p \cdot q}$

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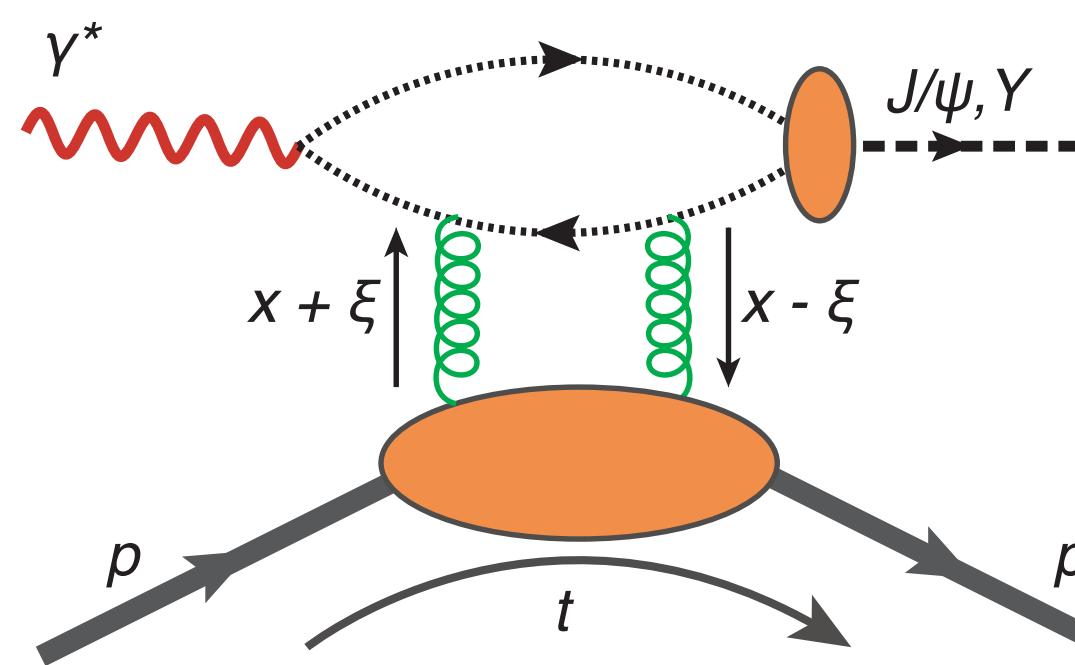
**Modified Bjorken- $x$ :**  $x_V = \frac{Q^2 + M_V^2}{2p \cdot q}$

Fourier transform:  
3-D transverse gluonic density

$$\rho(|\vec{b}_T|, x_V) = \int \frac{d^2 \vec{\Delta}_T}{(2\pi)^2} e^{i \vec{\Delta}_T \vec{b}_T} |\langle H_g \rangle|(t = -\vec{\Delta}_T^2)$$

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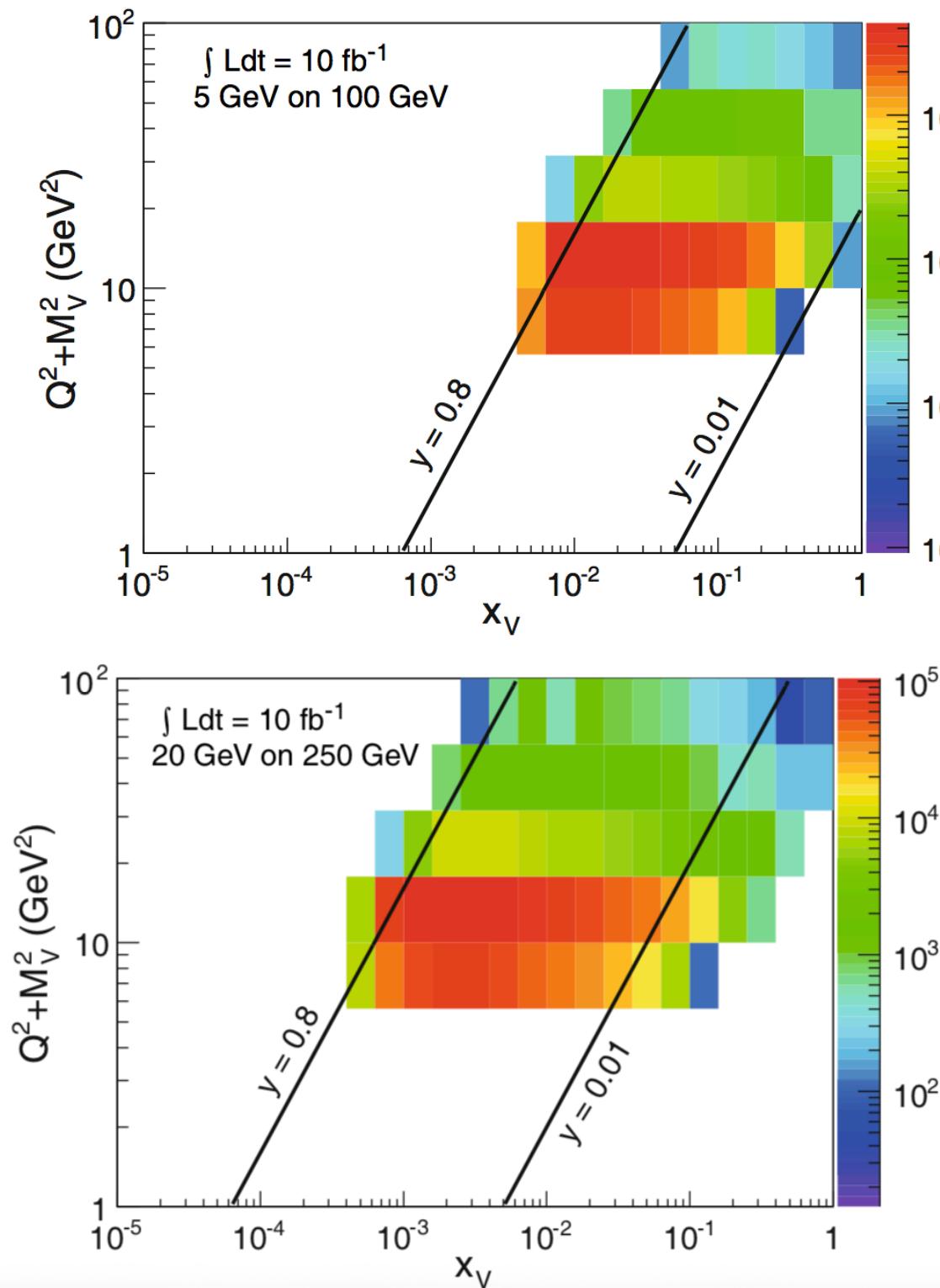
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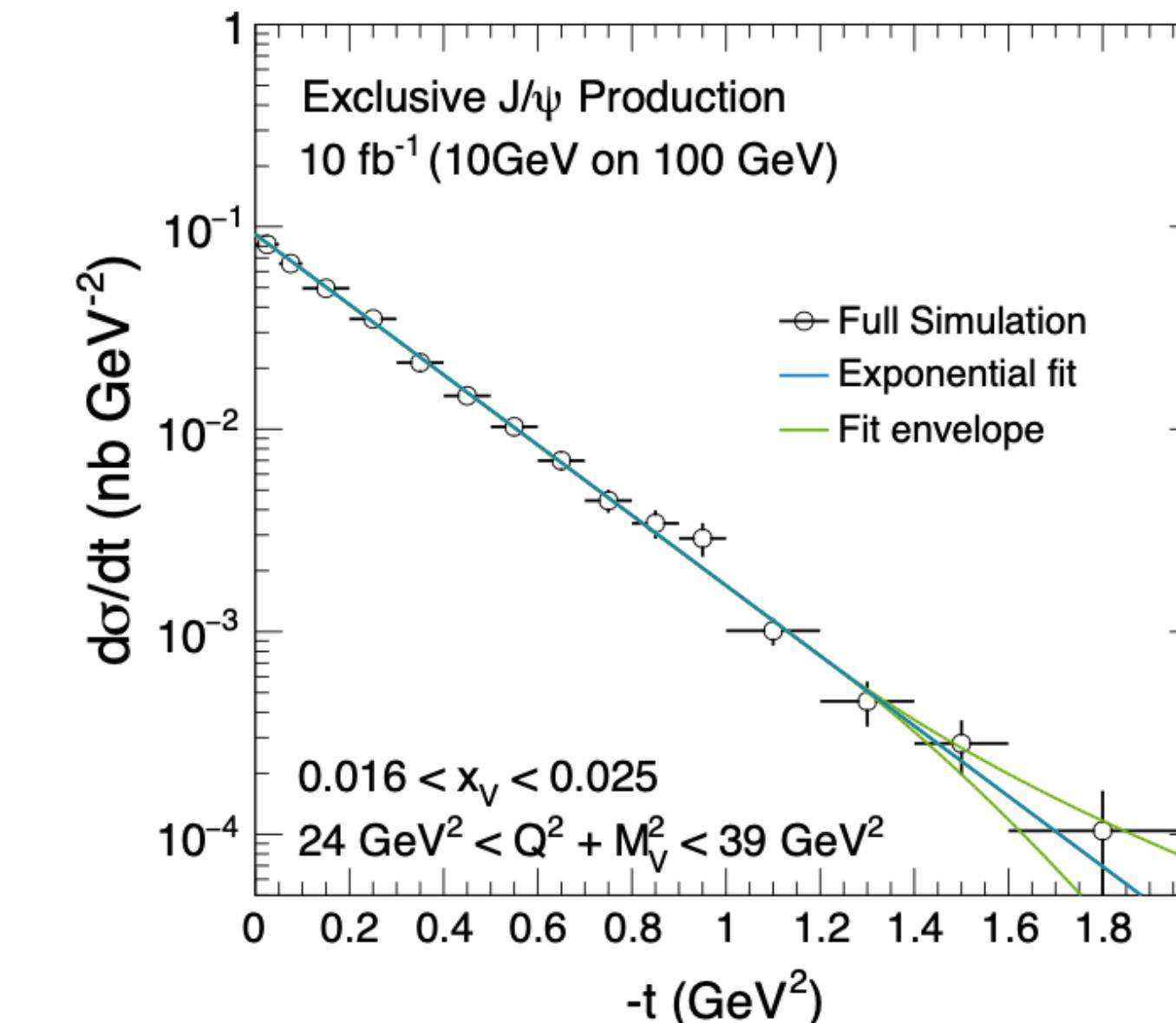
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3-D GPDs can be related to 2-D Gravitational Form Factors

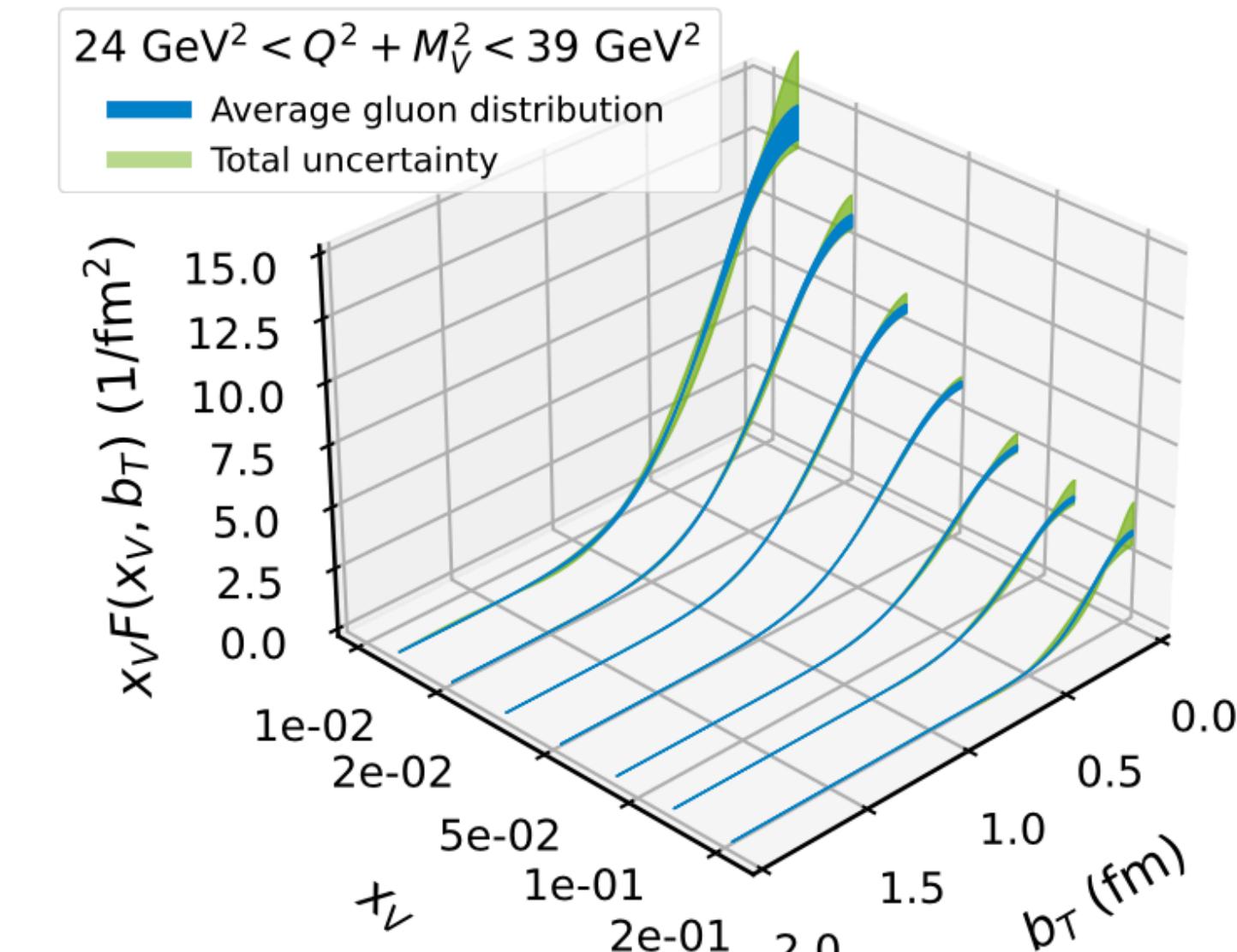
# GLUON TOMOGRAPHY WITH J/Ψ



**Only possible at an EIC:**  
 from the valence region  
 deep into the sea!



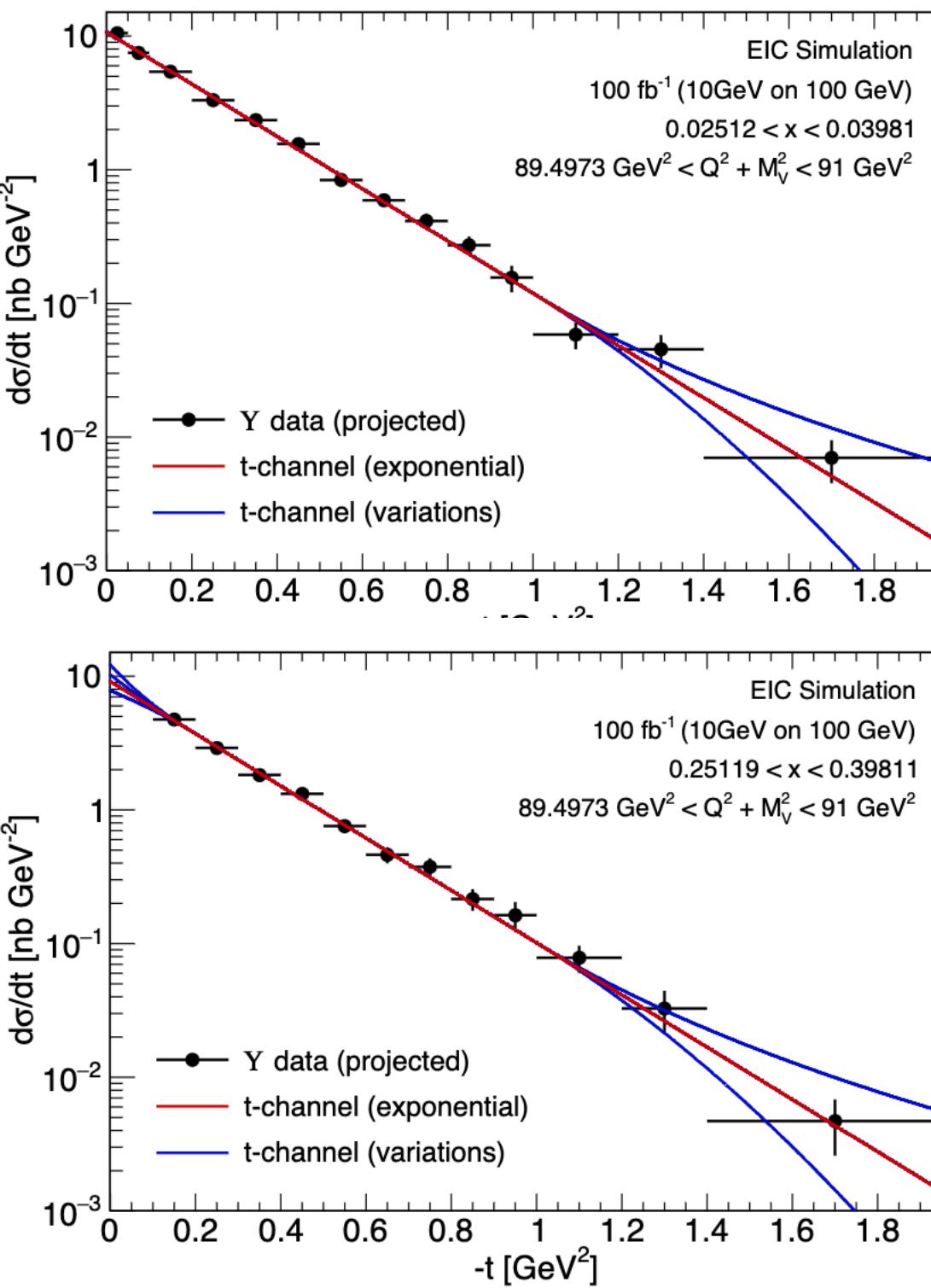
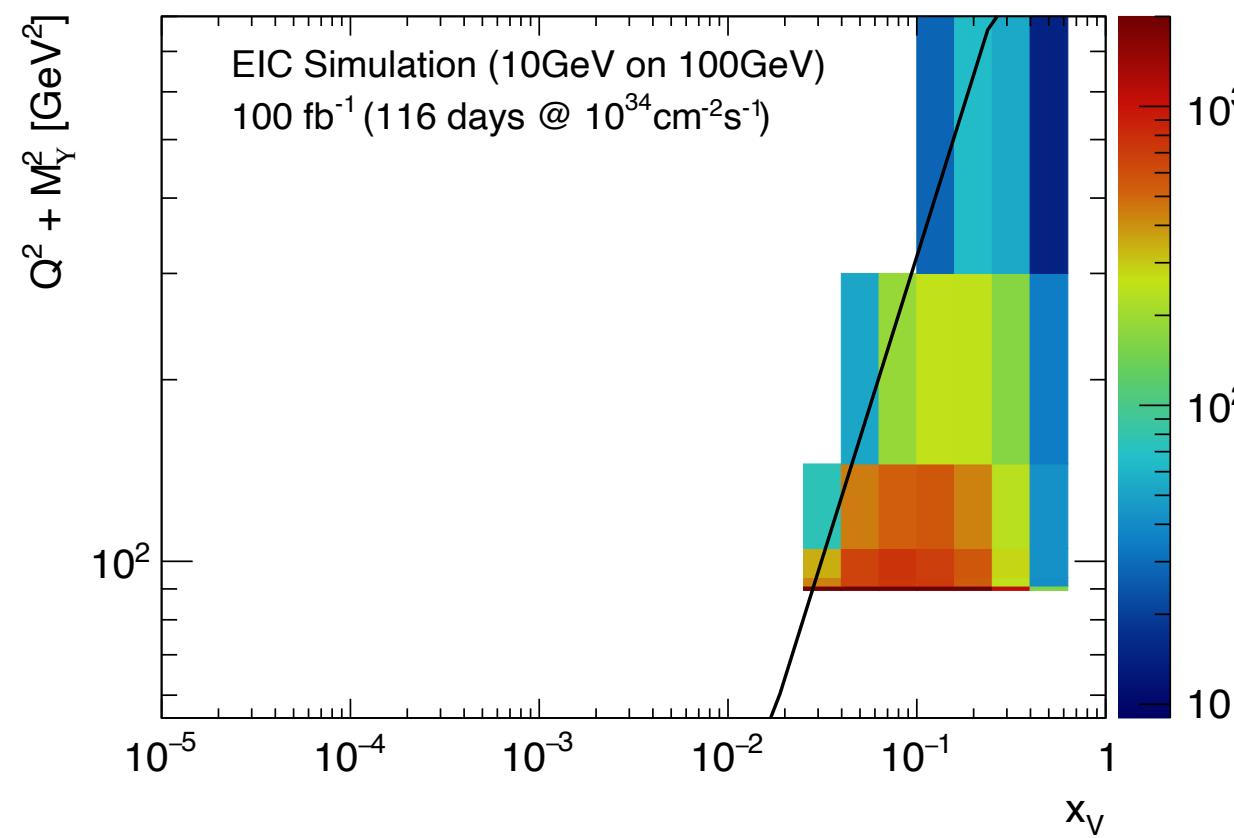
***t*-spectra**



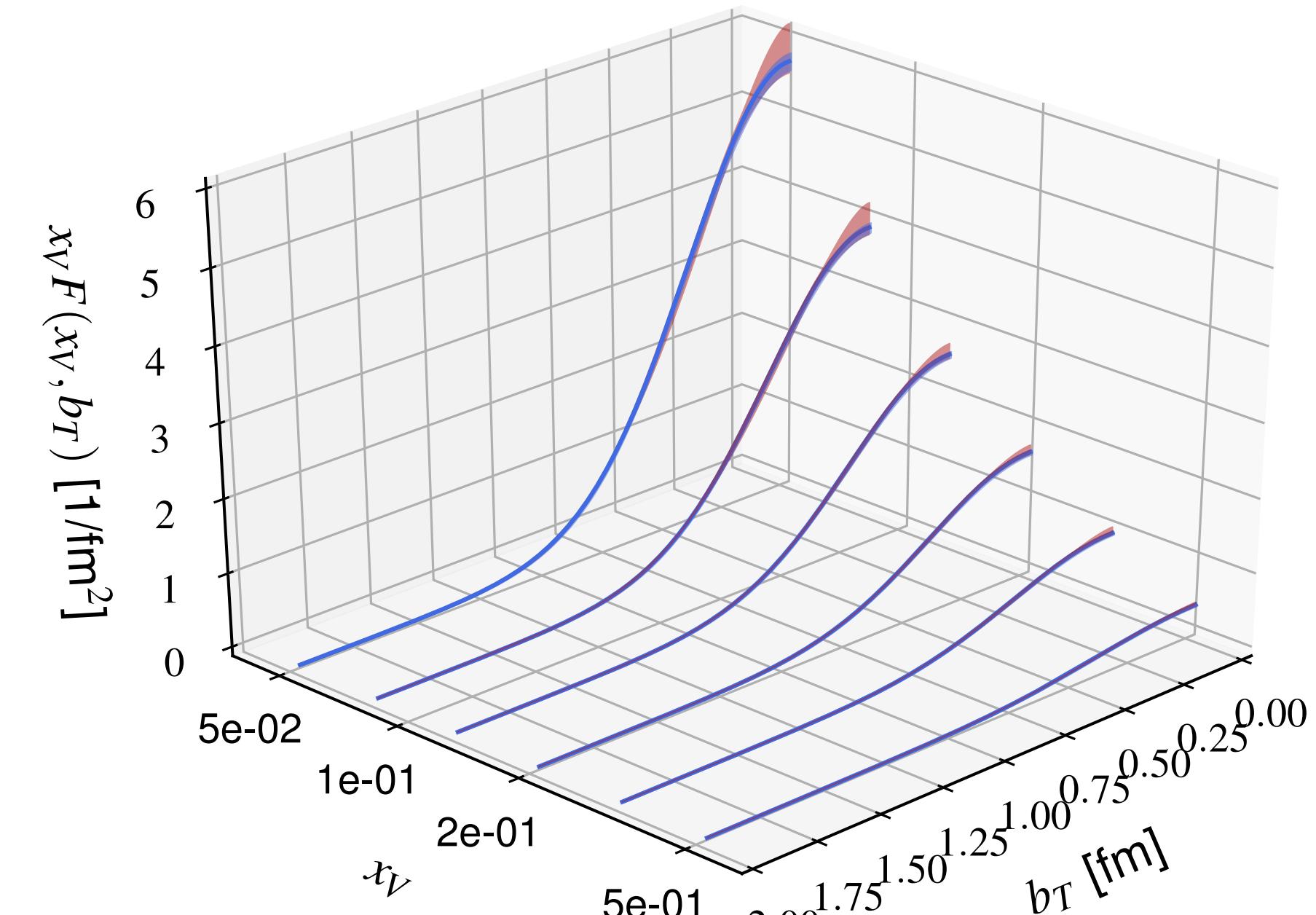
**Normalized average gluon density**

Same can be done with Y(1S) at EIC!

# GLUON TOMOGRAPHY WITH Y(1S)



*t-spectra*



Average gluon density:

1 year at peak luminosity

# CONCLUSION

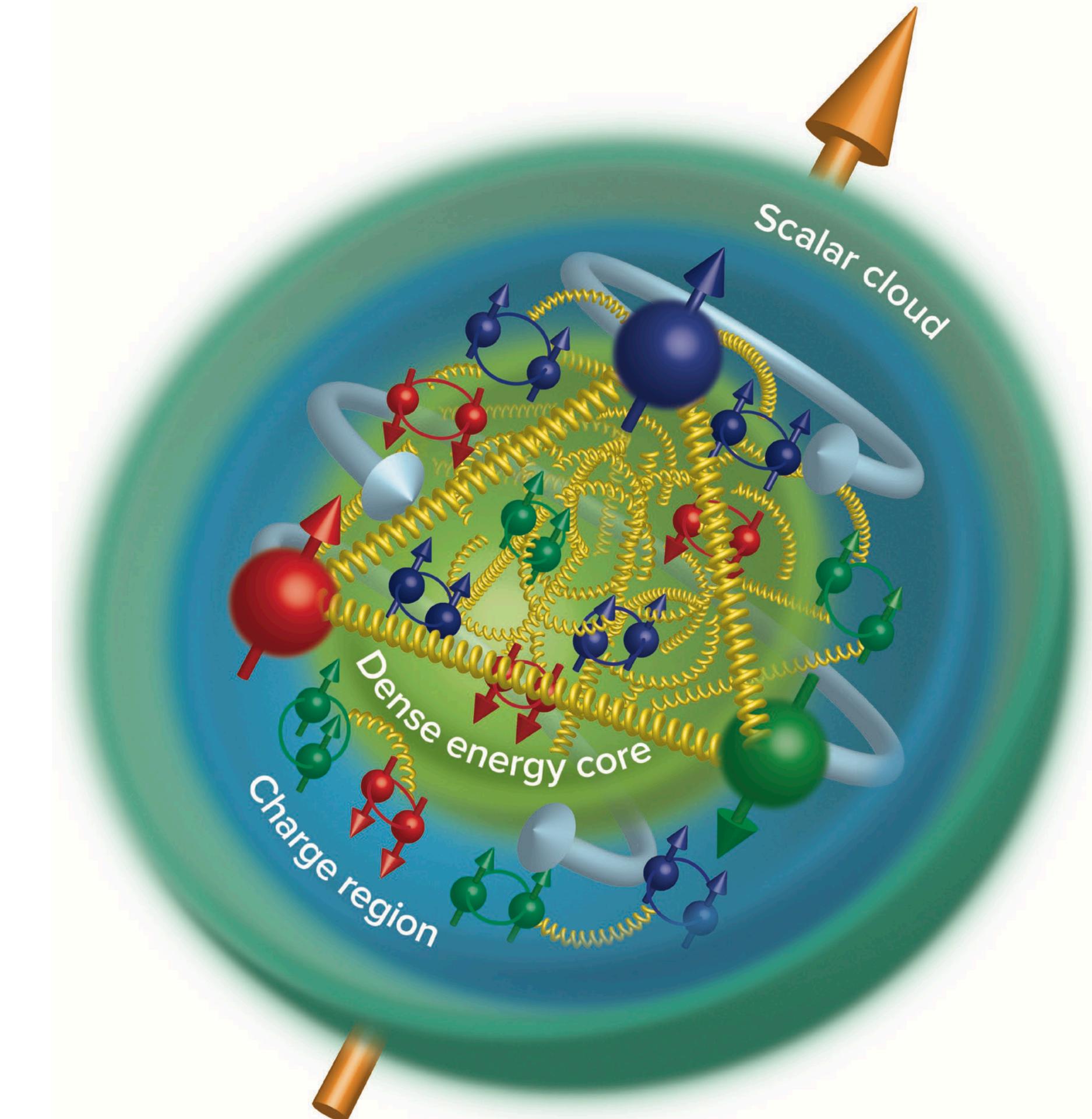
The JLab 12-GeV program has delivered important first results on near-threshold  $J/\psi$  production from GlueX and Hall C ( $J/\psi$ -007)

- A new window on the gluonic structure of the proton
- Does the proton have a dense energy core?
- What are the implications of a possible scalar gluonic cloud?

The planned near-threshold  $J/\psi$  production program at Jefferson Lab is crucial to further our understanding of the origin of mass.

- SoLID can reach  $J/\psi$  observables that cannot be achieved anywhere else, including precision measurements at high  $t$  and precision electroproduction near threshold.

The mass structure of the nucleons and nuclei is a rapidly evolving topic, reaching from Jefferson Lab to the EIC

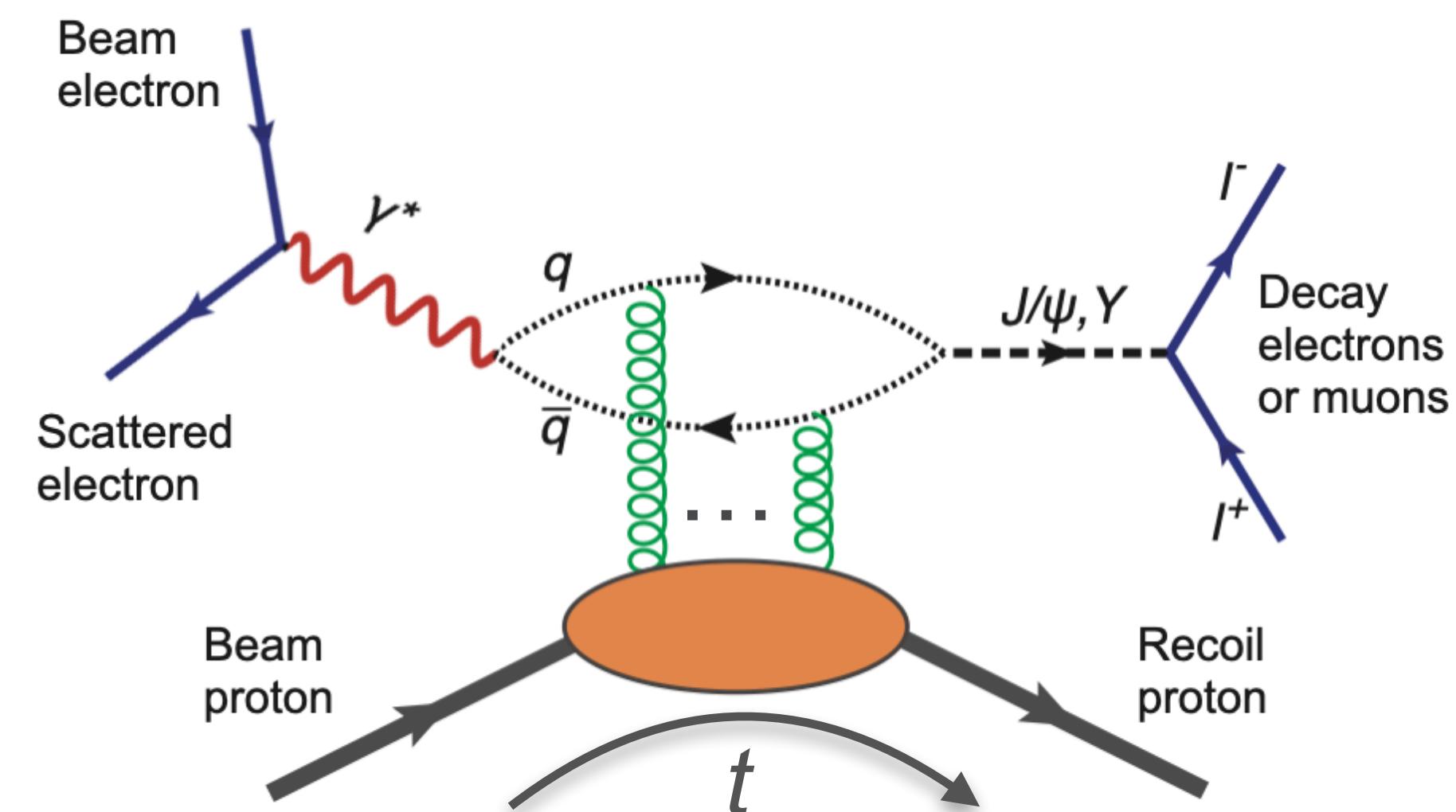


# QUESTIONS?

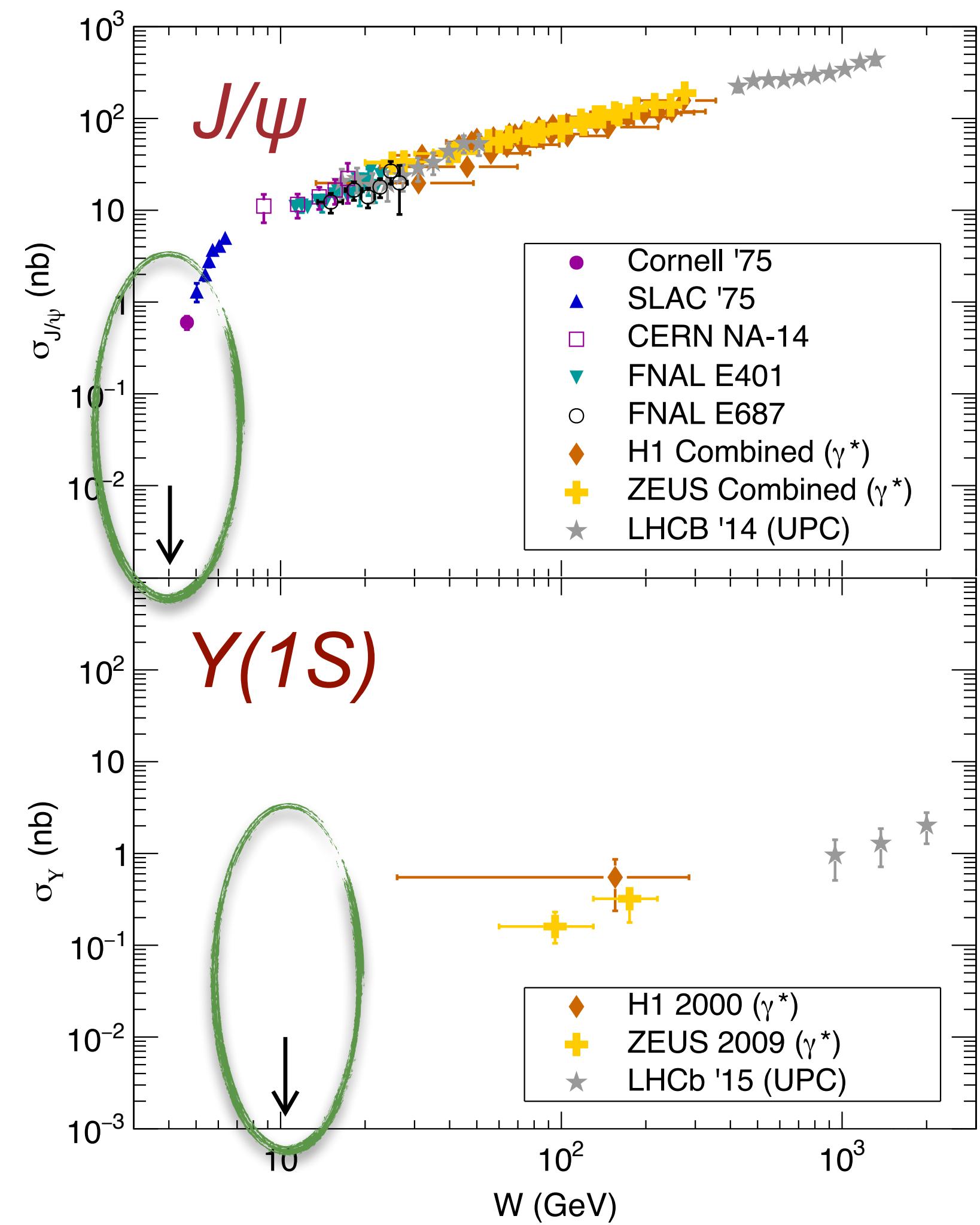


# WHAT DID WE KNOW BEFORE 12 GEV AT JLAB?

## Exclusive quarkonium photoproduction



- No near-threshold data available
- In case of  $Y(1S)$ : not much available overall
- **Almost no data near threshold before JLab 12 GeV**



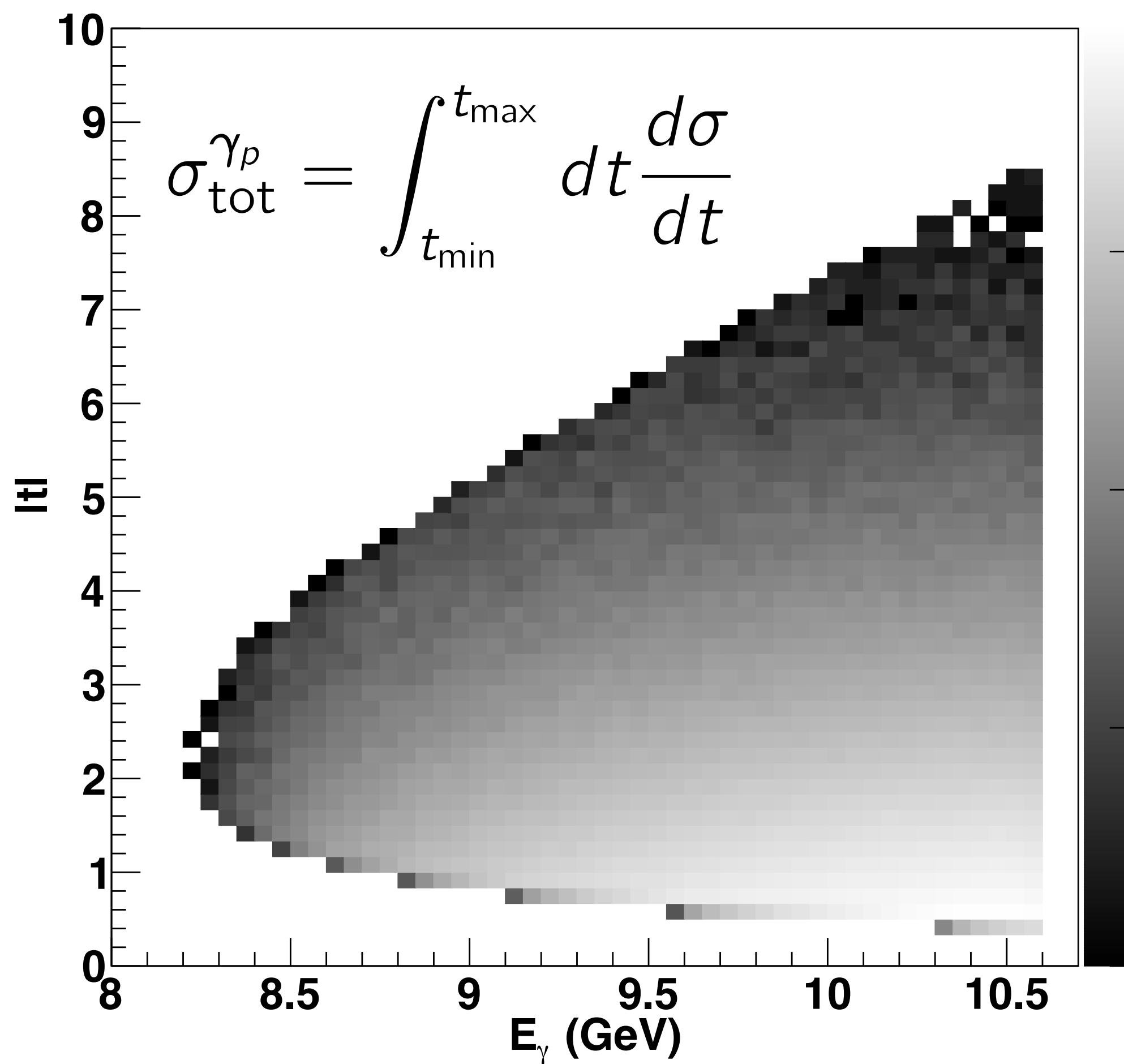
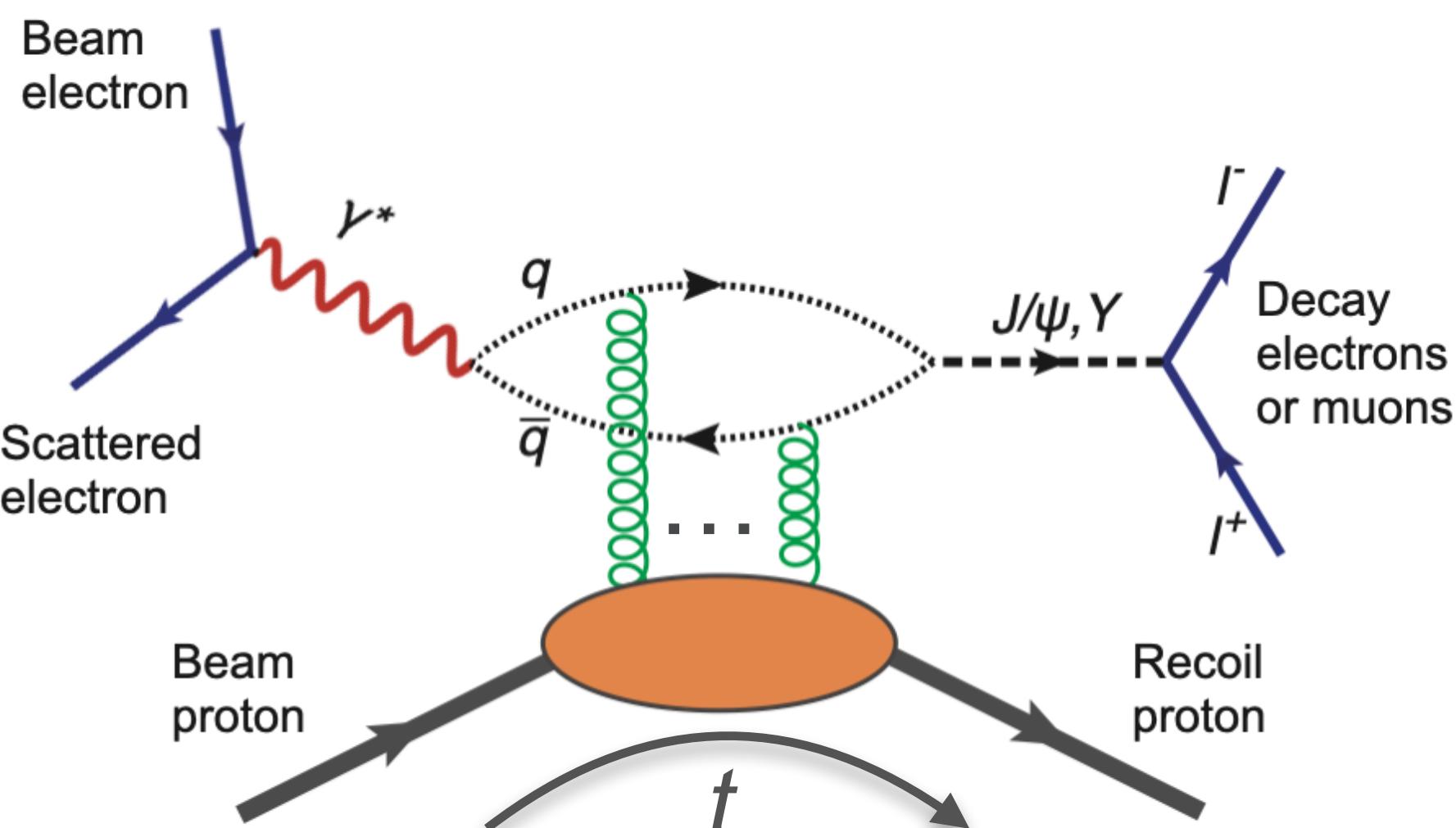
# J/Ψ EXPERIMENTS AT JLAB COMPARED

	GlueX HALL D	HMS+SHMS HALL C	CLAS 12 with upgrade <sup>1</sup> HALL B	SoLID HALL A
J/ψ counts (photo-prod.)	469 published ~10k phase I + II	2k electron channel 2k muon channel	14k	804k
<i>J/ψ</i> Rate (electro- prod.)	N/A	N/A	1k	21k
Features	Good reach to threshold. No high-t reach.	Can reach high-t only at higher energies. Low statistics.	No high-t reach. Electroproduction low statistics.	Enough luminosity to reach high t. High precision.
When?	Finished/Ongoing	Finished	Ongoing/Proposed	Future

<sup>1</sup>The CLAS12 projected count rates assume the proposed CLAS12 luminosity upgrade to  $2 \times 10^{35}/\text{cm}^2/\text{s}$

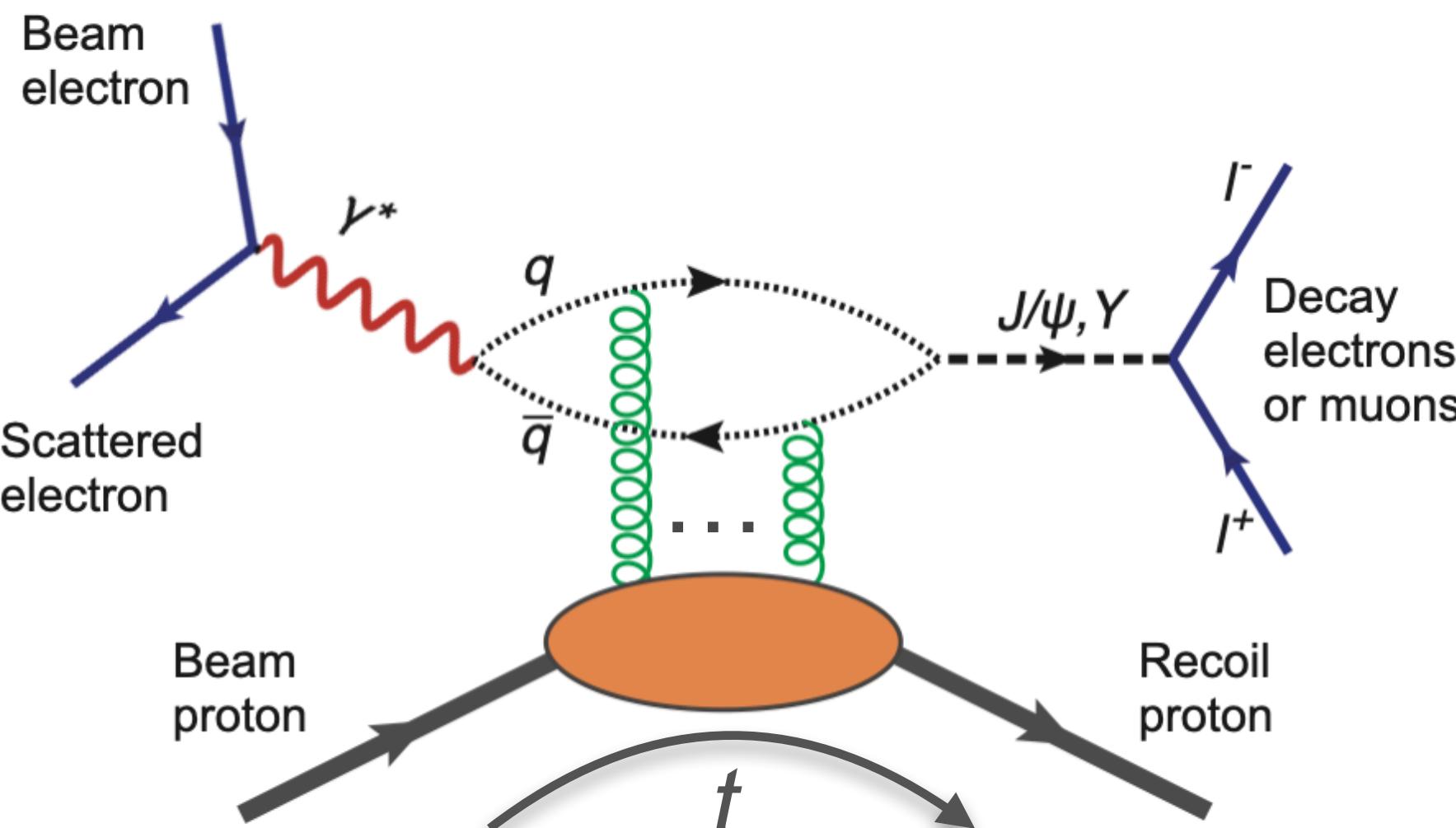
# EXCLUSIVE QUARKONIUM PRODUCTION

## The basics



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## The basics



$J/\psi$  threshold:

$$W \approx 4.04 \text{ GeV}$$

$$E_\gamma^{\text{lab}} \approx 8.2 \text{ GeV}$$

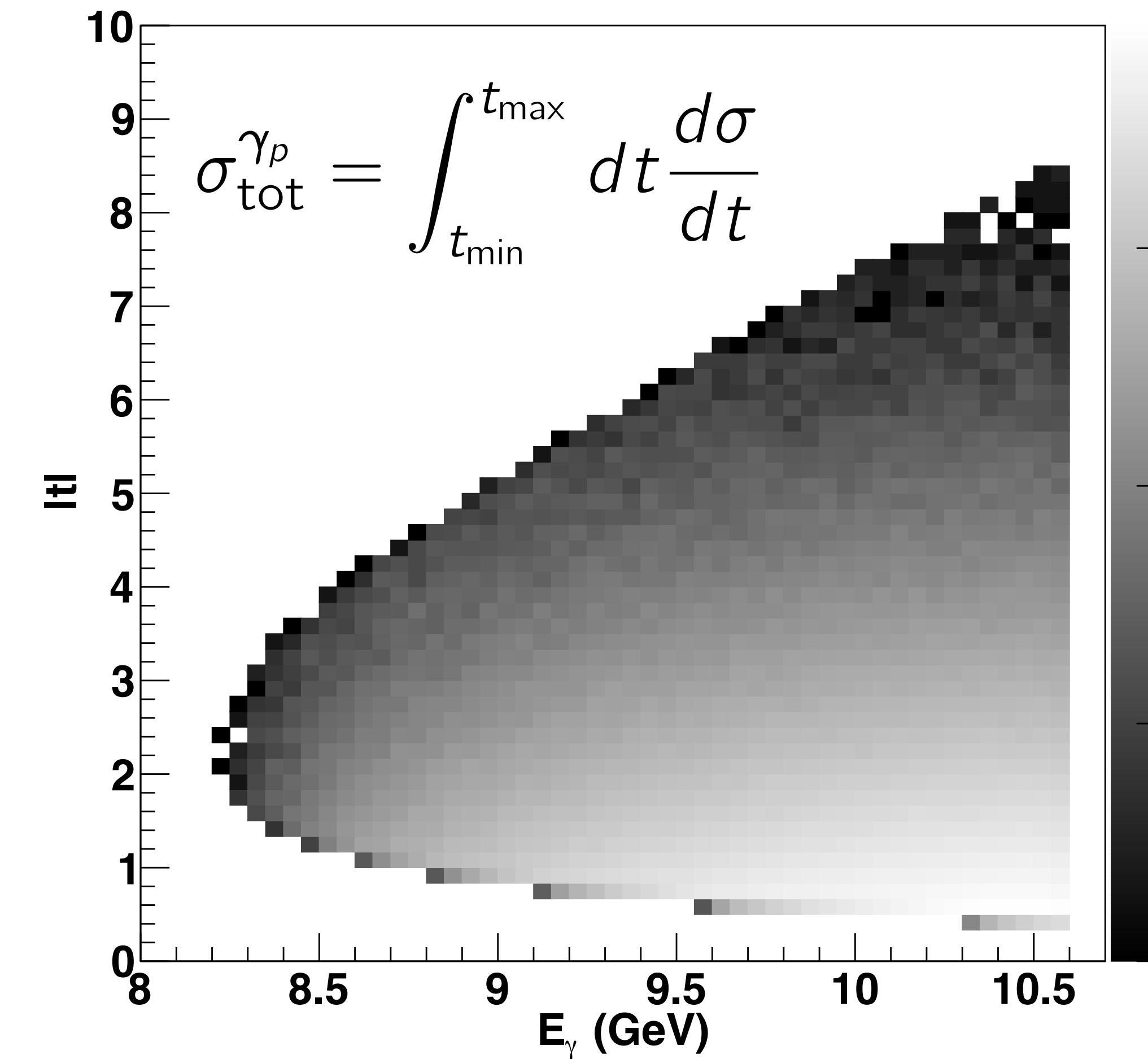
$$t \approx -1.5 \text{ GeV}^2$$

$Y(1S)$  threshold:

$$W \approx 10.4 \text{ GeV}$$

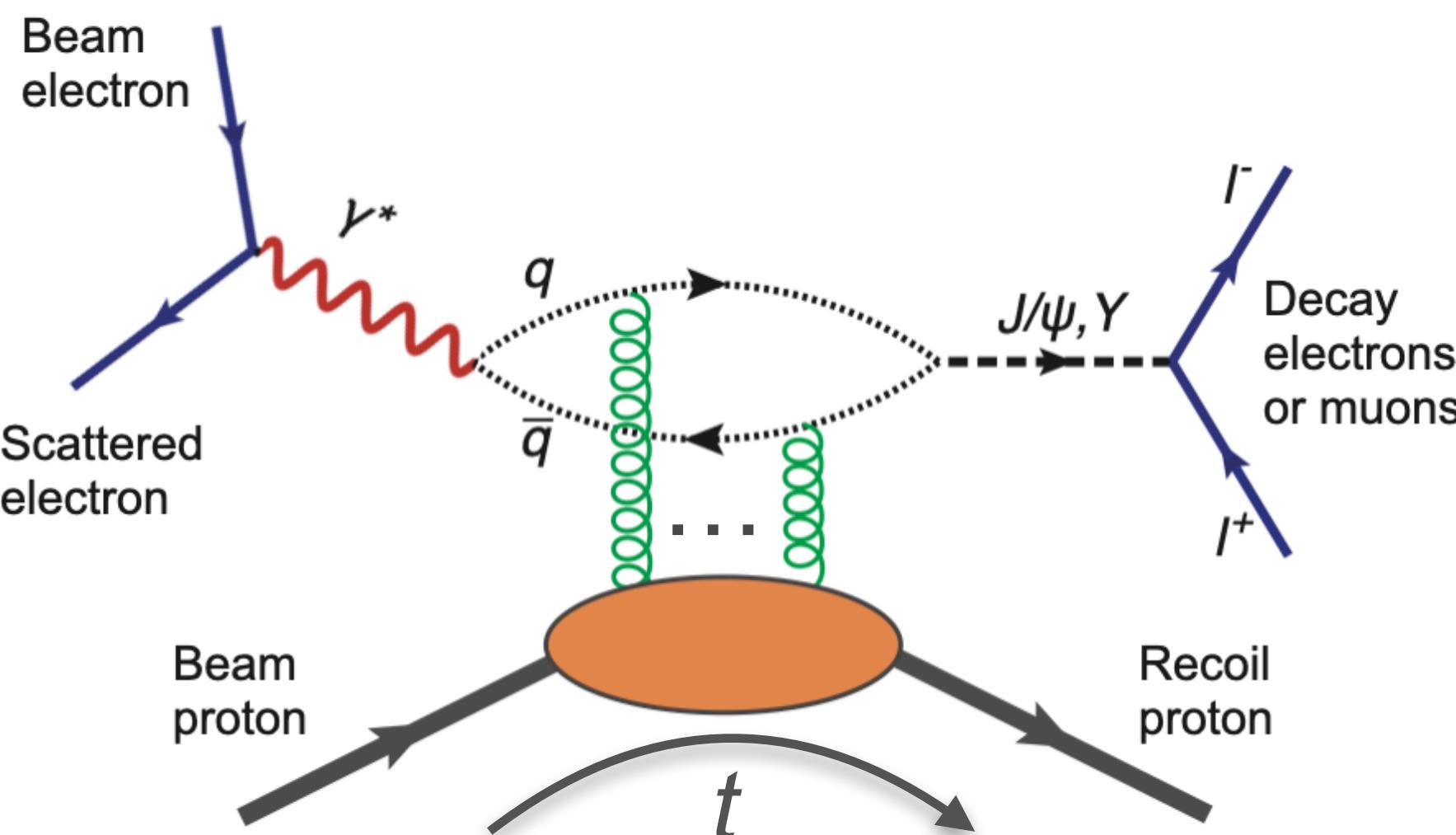
$$t \approx -8.1 \text{ GeV}^2$$

- Phase space limits defined by quarkonium direction



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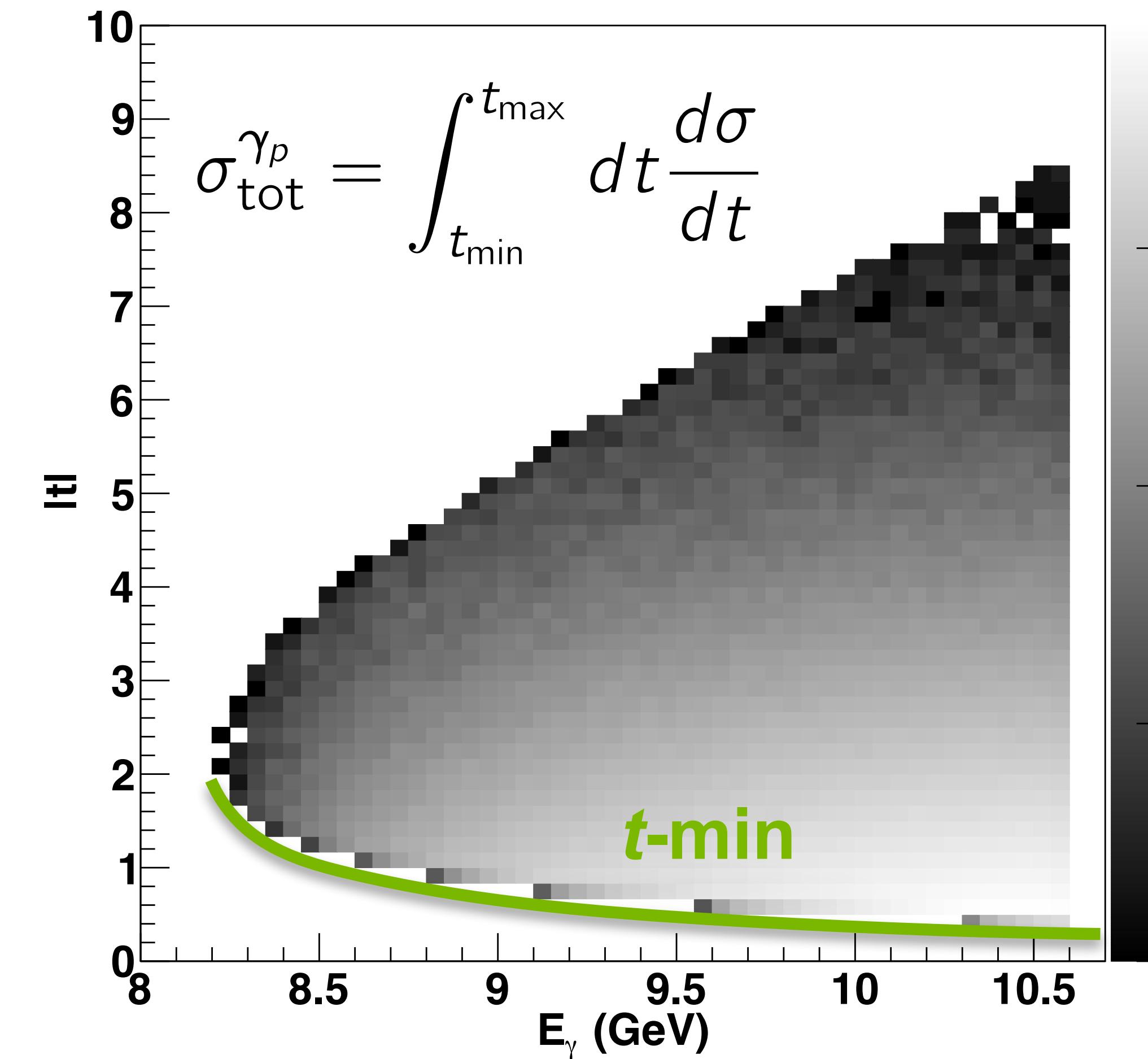
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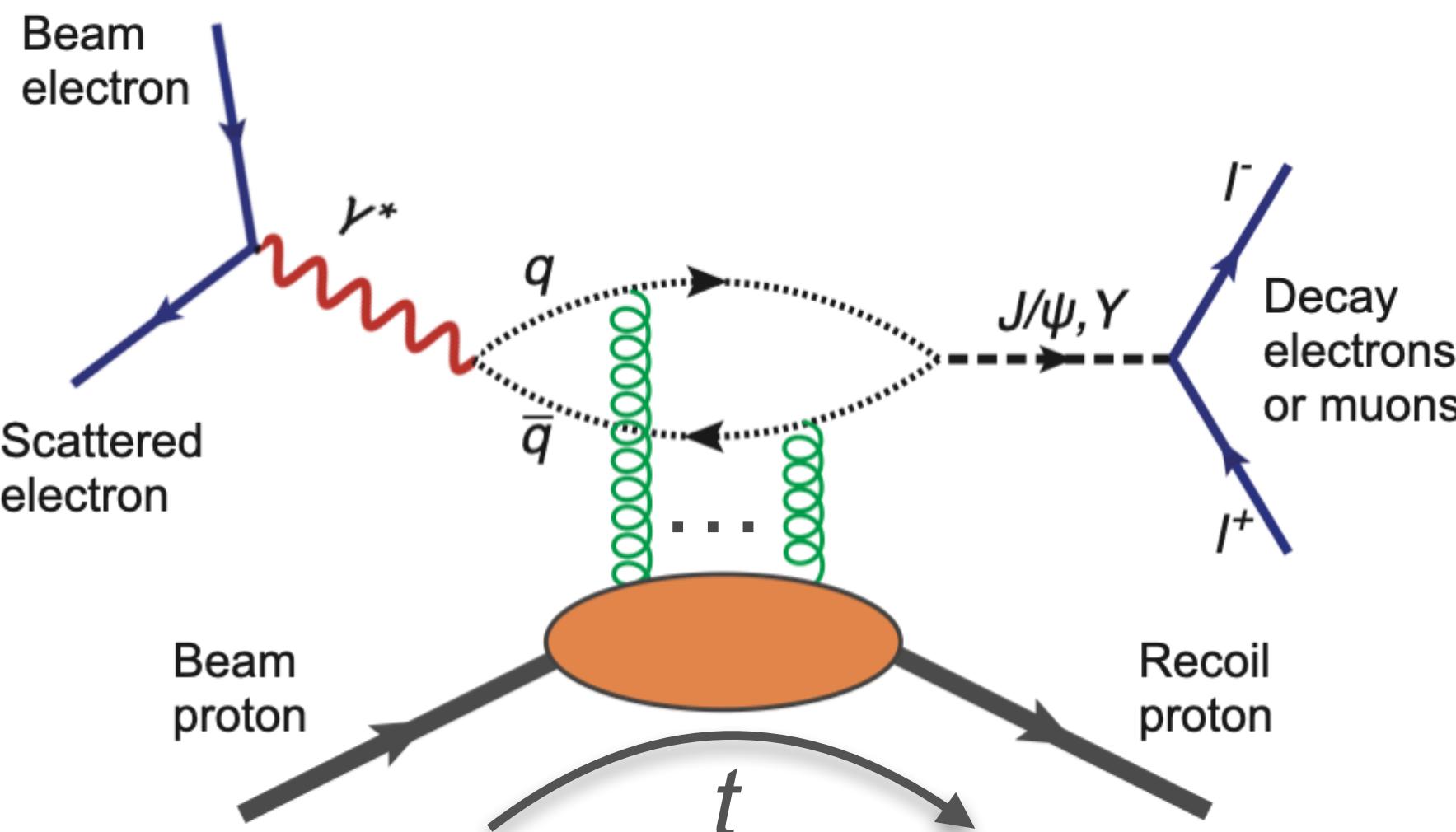
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- Phase space limits defined by quarkonium direction
- Forward (with photon):  $t = t_{\min}$



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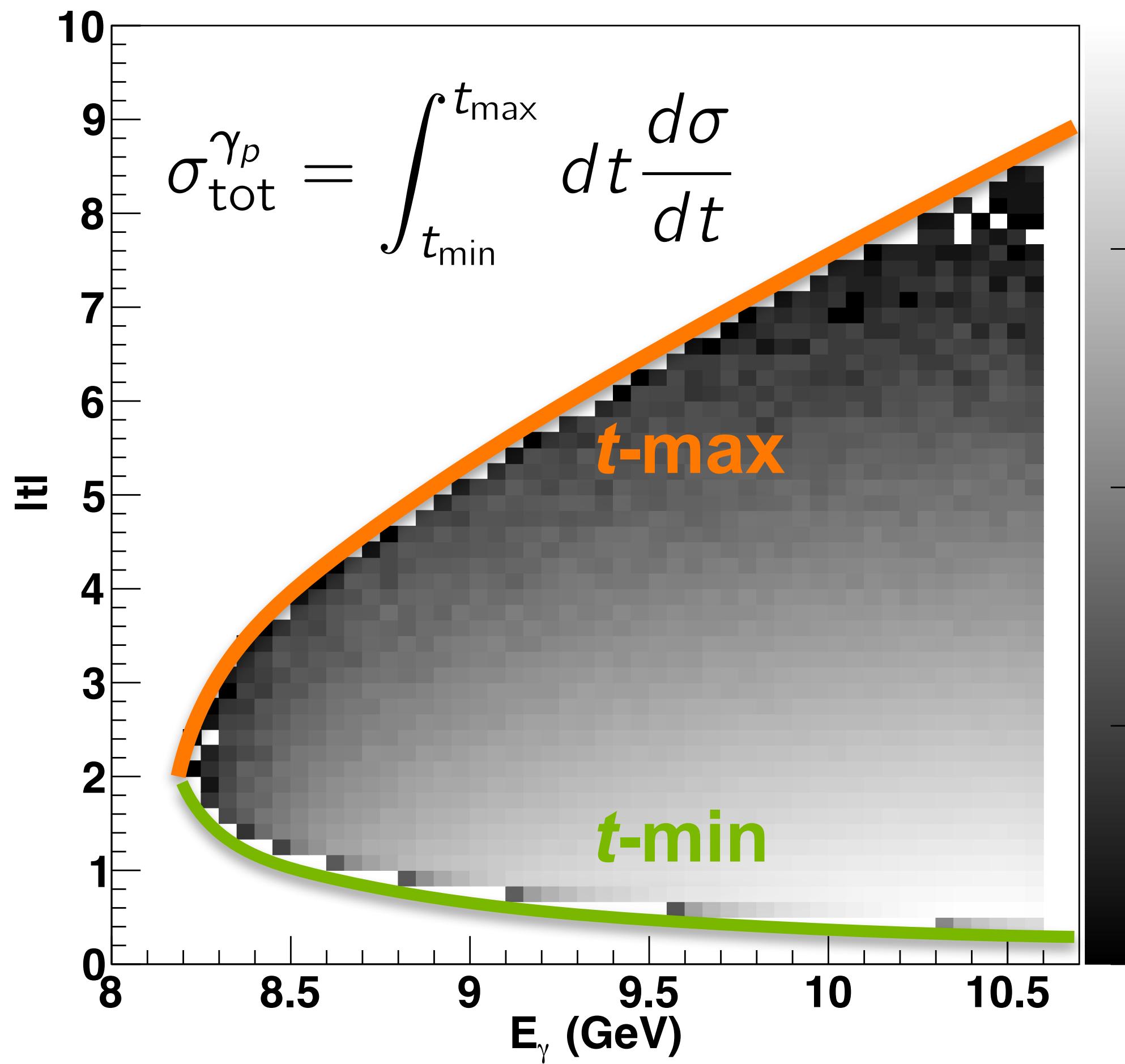
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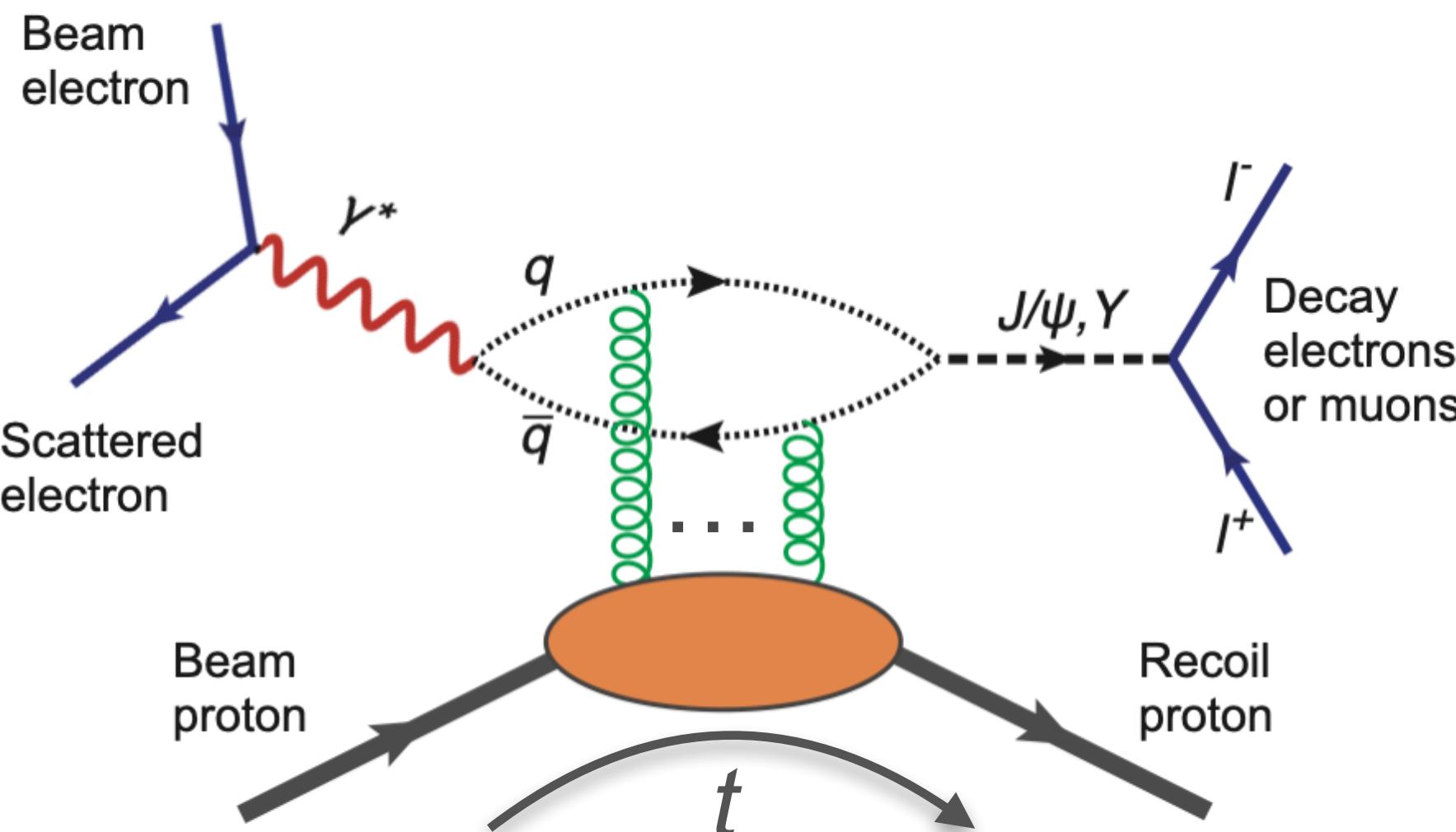
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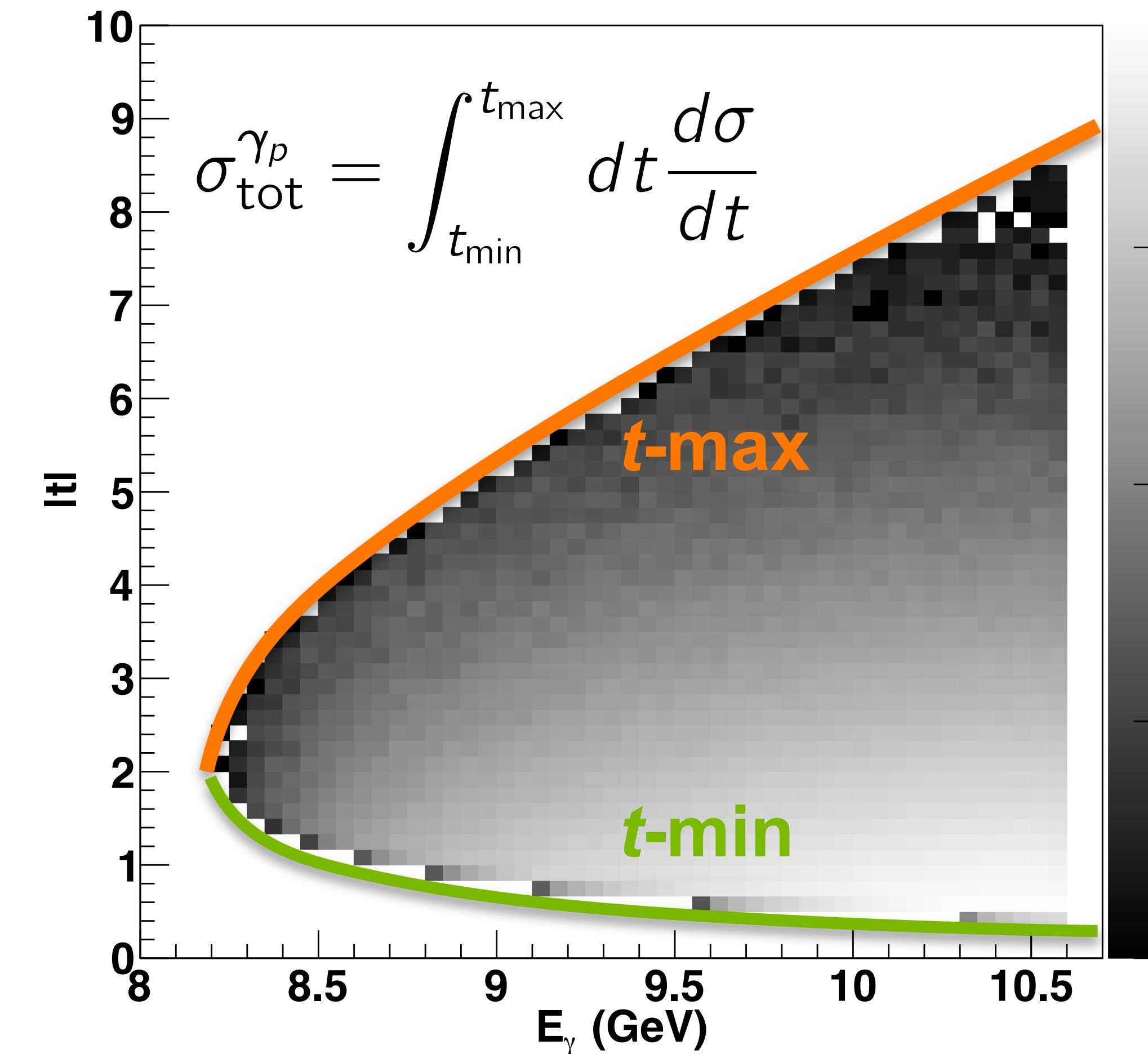
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- Forward (with photon):  $t = t_{\min}$
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- Forward direction preferred:  $t$ -dependence  $\sim$ exponential



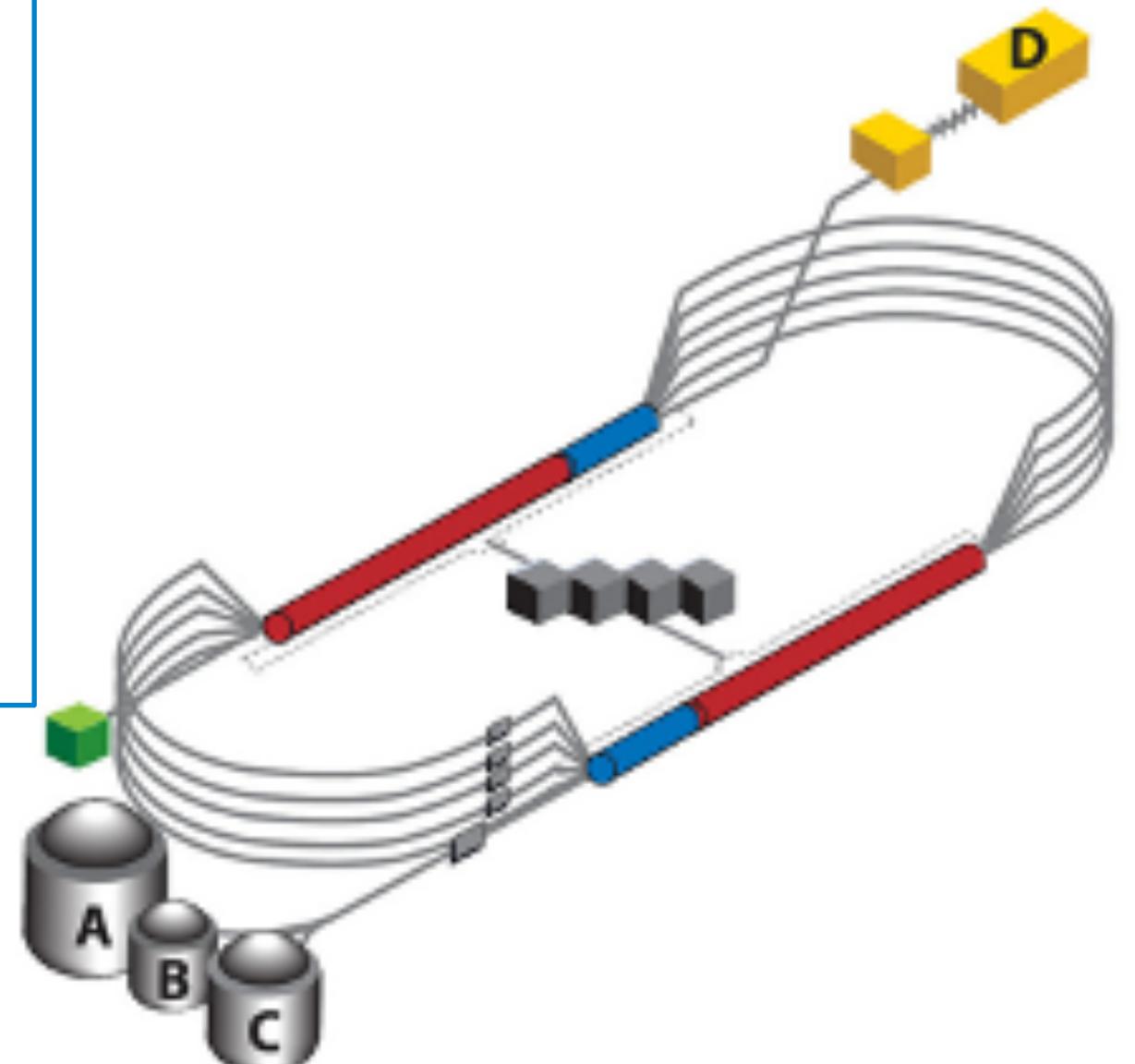
# QUARKONIUM AT JEFFERSON LAB AND EIC

## Jefferson Lab

CEBAF: very high luminosity ( $10^{35}$ - $10^{39} \text{ cm}^{-2}\text{s}^{-1}$ ) continuous electron beam on fixed target

- 4 experimental halls:
- 11GeV in Hall A, B & C
  - 12GeV in Hall D

Jefferson Lab is the ideal laboratory to measure  $\text{J}/\psi$  near threshold, due to luminosity, resolution and energy reach



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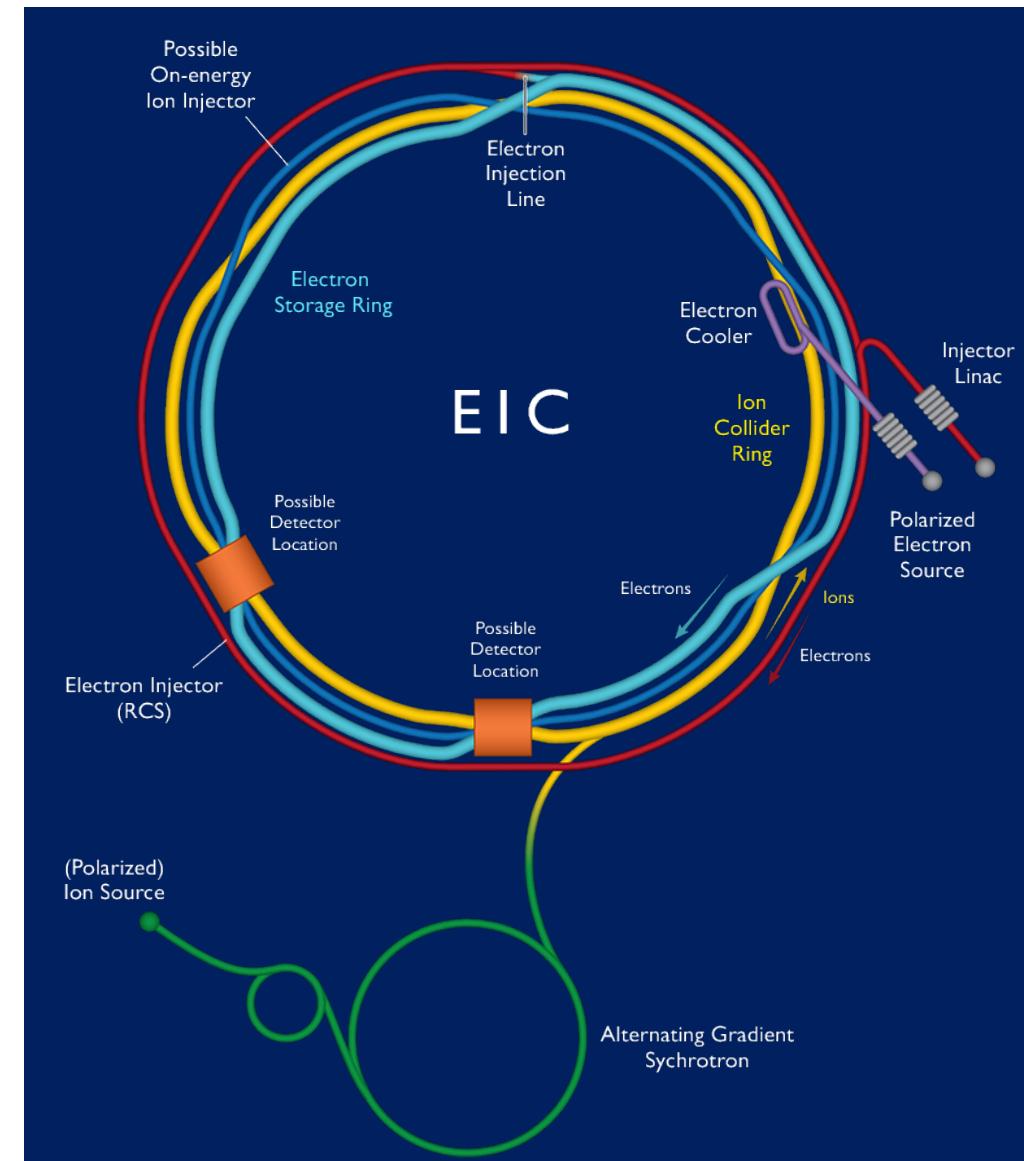
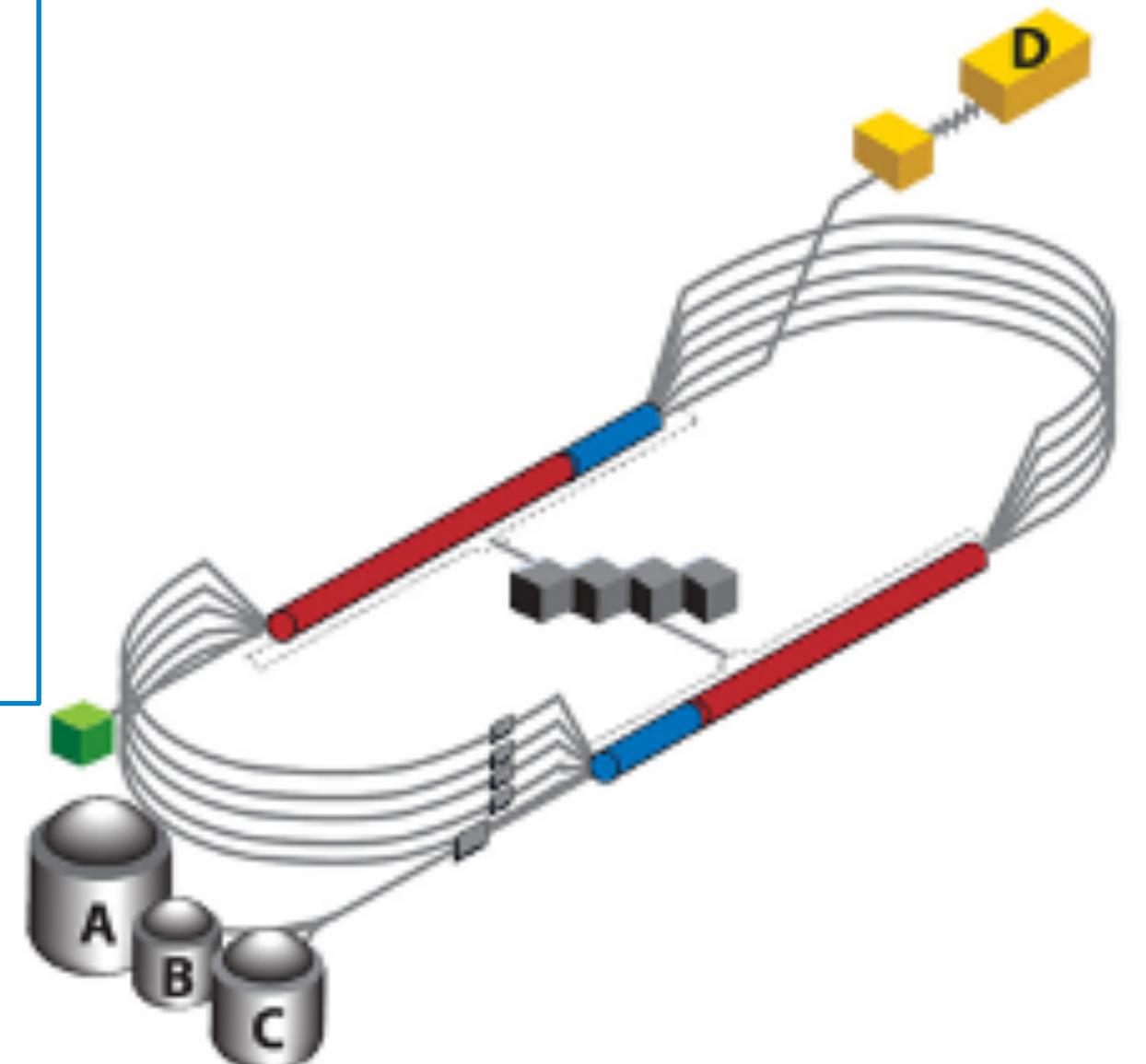
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## Electron-ion Collider

EIC: high luminosity ( $10^{33}$ - $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ ) polarized electron polarized ion collider

Variable CM energies: 29-140 GeV with 2 possible interactions regions

Ideally suited to study to  $\text{J}/\psi$  electroproduction at higher energies, sufficient energy and luminosity to produce  $\text{Y}(1S)$ .

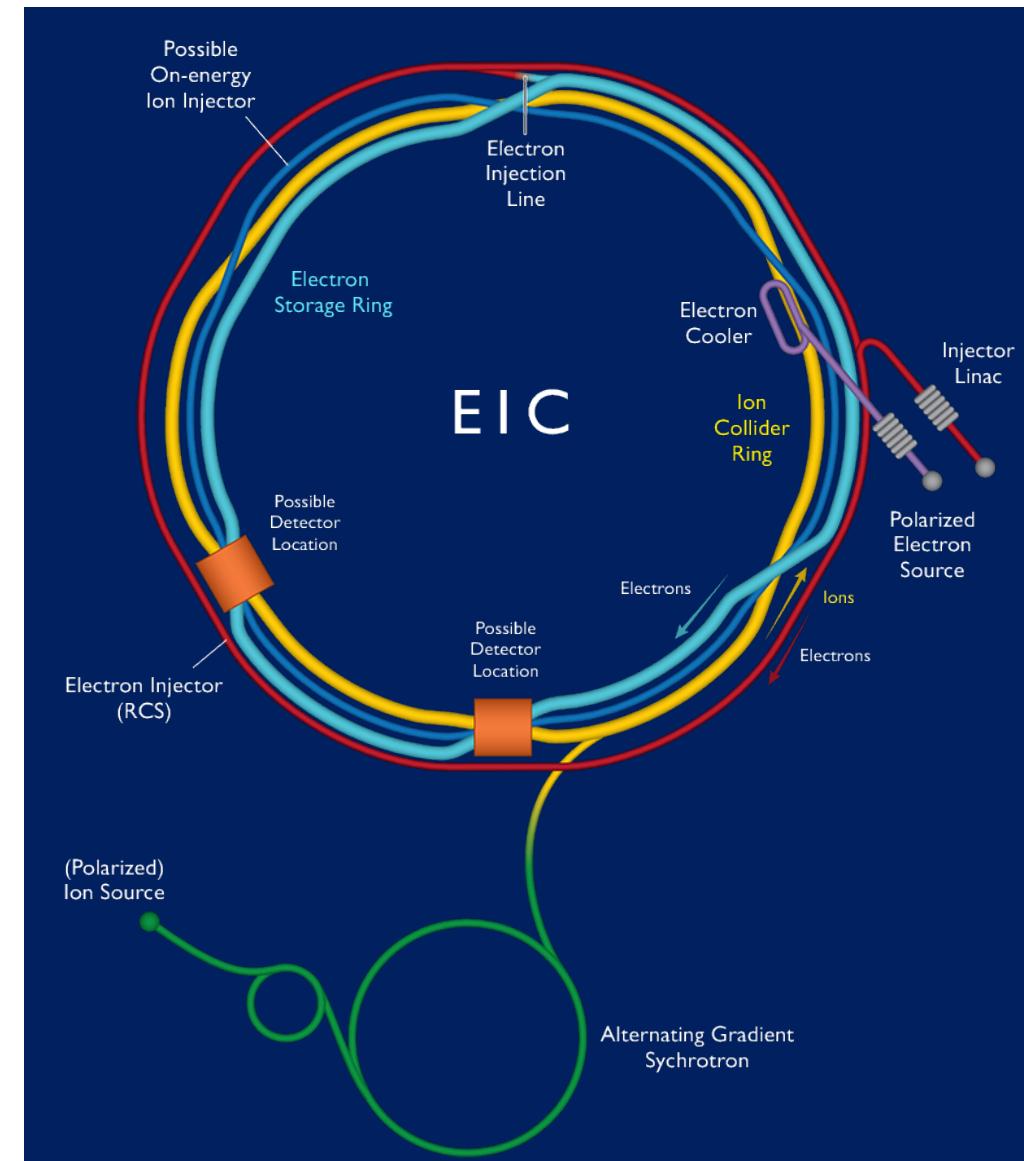
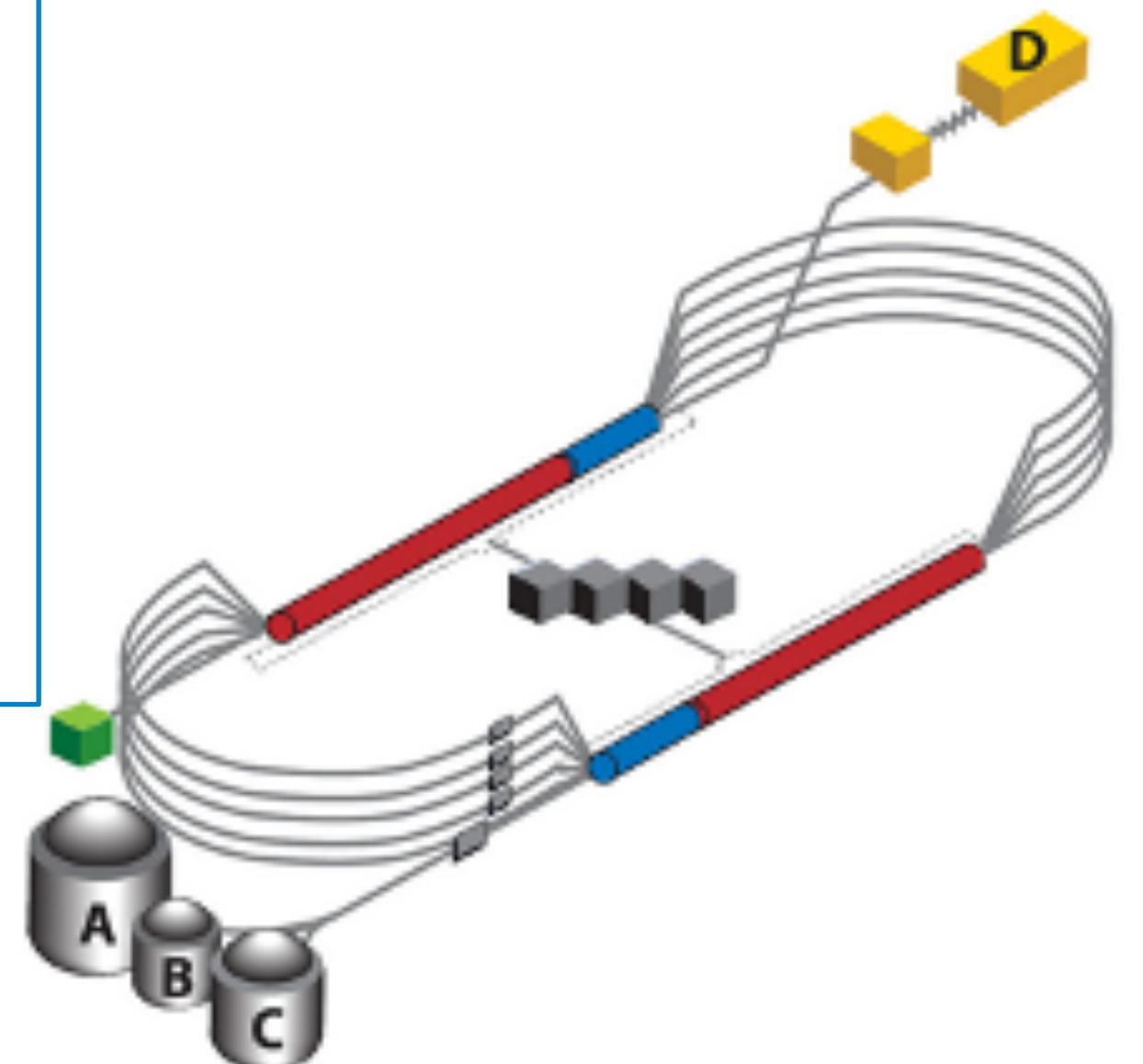
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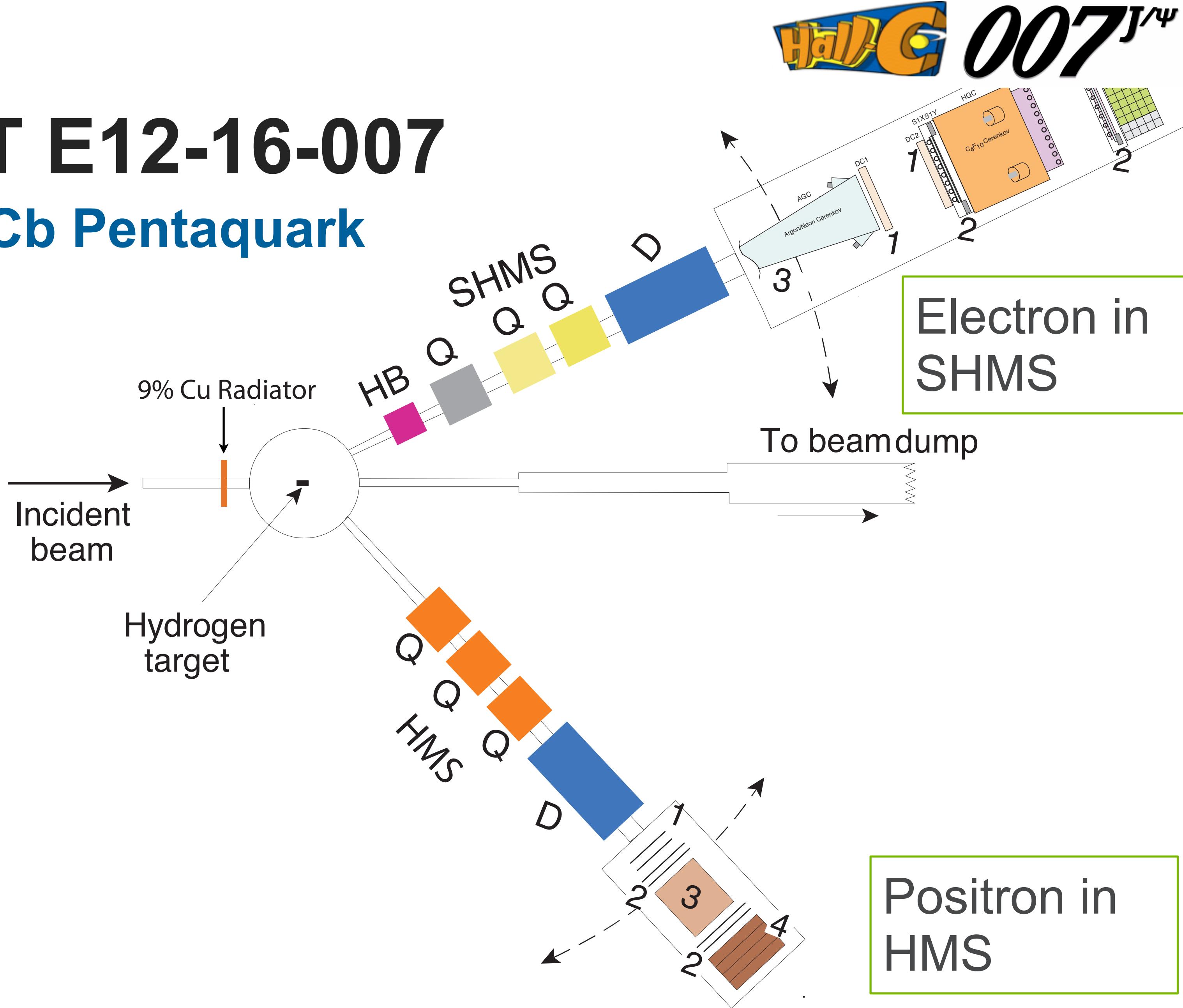
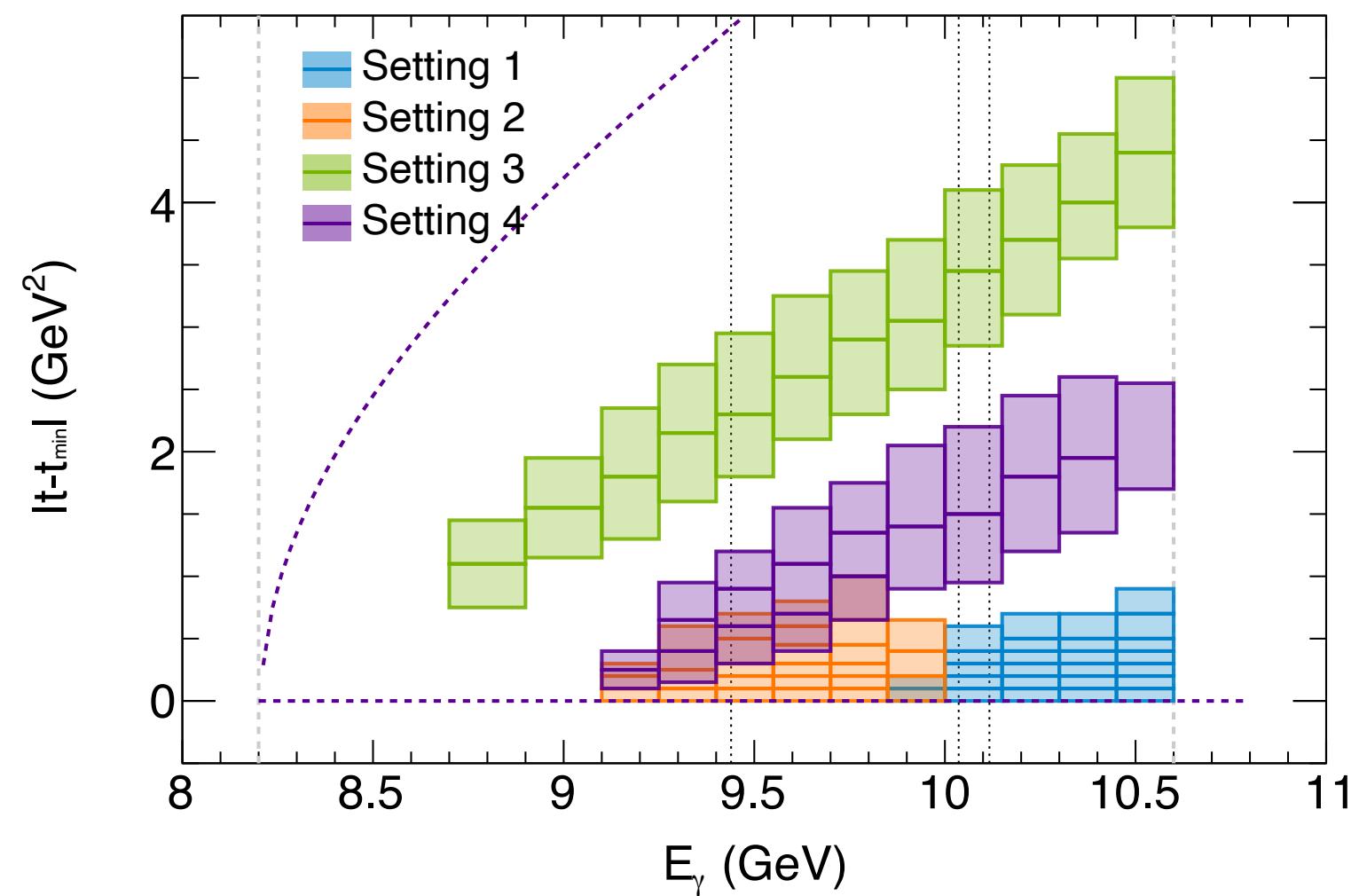
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**Complementary programs:** Jefferson Lab is the ideal laboratory to access GFFs with  $\text{J}/\psi$  production, and EIC has sufficient luminosity and energy for 3-D gluonic imaging.

# JLAB EXPERIMENT E12-16-007

## J/ $\psi$ -007: Search for the LHCb Pentaquark

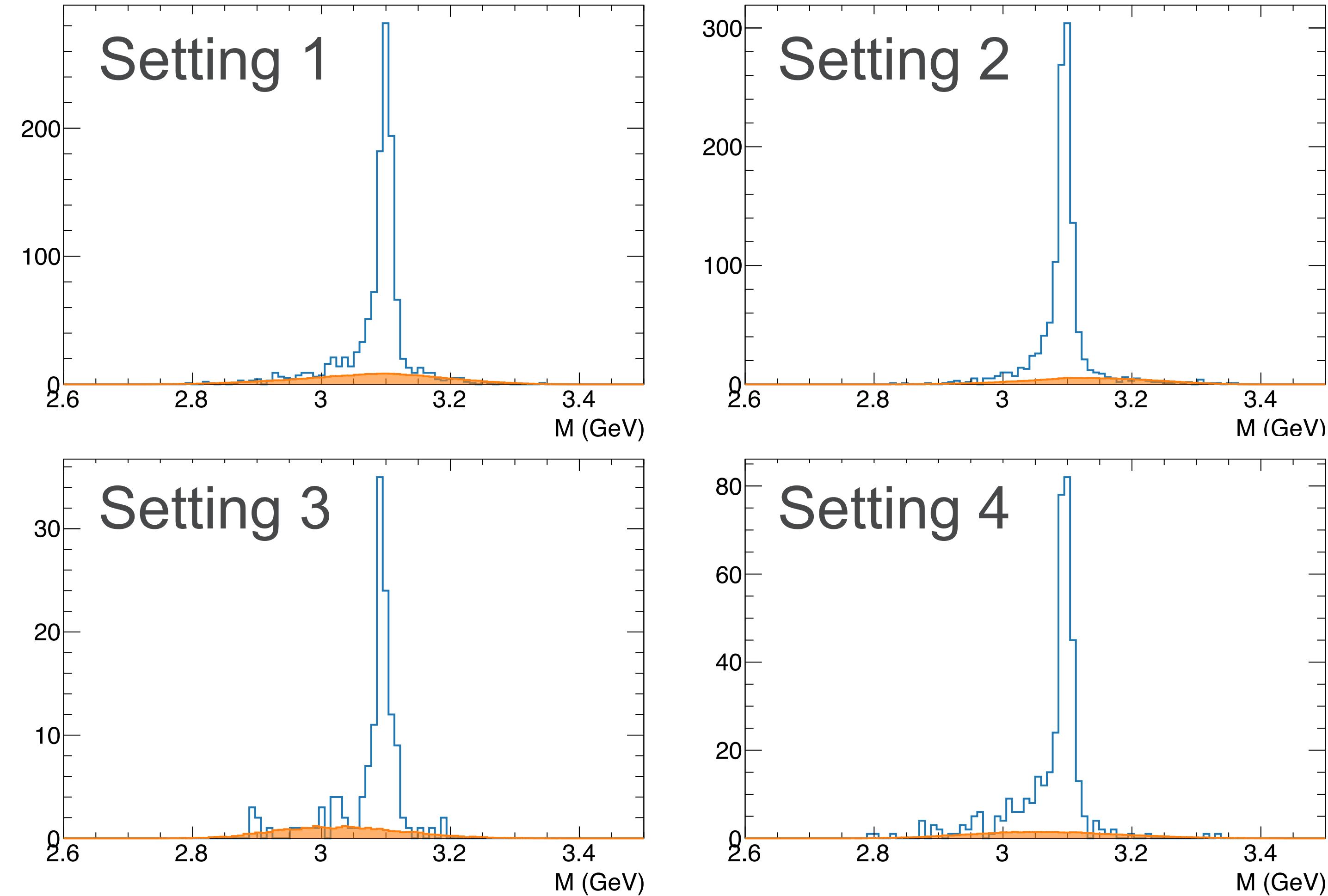
- Ran February 2019 for ~8 PAC days
- High intensity real photon beam (50 $\mu$ A electron beam on a 9% copper radiator)
- 10cm liquid hydrogen target
- Detect J/ $\psi$  decay leptons in coincidence
  - Bremsstrahlung photon energy fully constrained



# CLEAR J/ $\Psi$ SIGNAL WITH MINIMAL BACKGROUND

007<sup>J/ $\Psi$</sup>

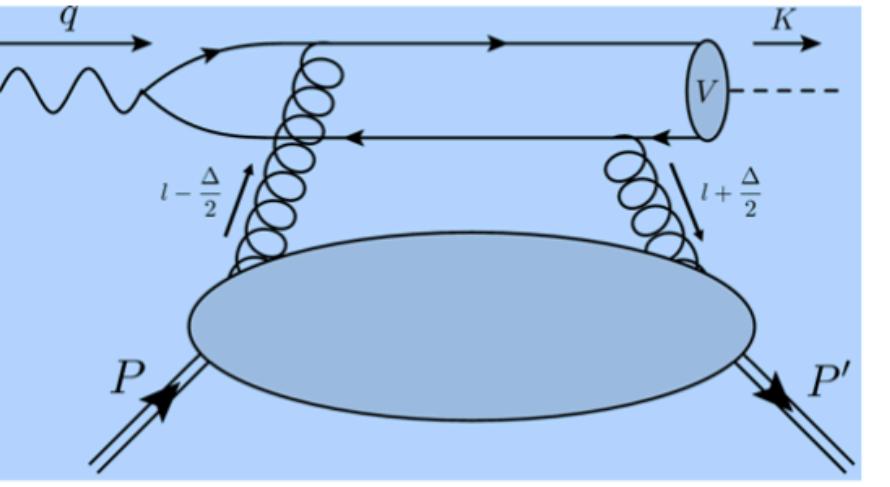
settings	HMS	SHMS	target	charge [C]	goal
setting 1	19.1° at +4.95GeV	17.0° at -4.835GeV	LH2 with radiator dummy with radiator LH2, no radiator	5.2 0.6 0.1	low- $t$ and high energy target wall electroproduction
setting 2	19.9° at +4.6GeV	20.1° at -4.3GeV	LH2 with radiator dummy with radiator	8.2 0.3	low- $t$ and low energy target wall
setting 3	16.4° at +4.08GeV	30.0° at -3.5GeV	LH2 with radiator	13.8	high- $t$
setting 4	16.5° at +4.4GeV	24.5° at -4.4GeV	LH2 with radiator dummy with radiator	6.9 0.2	medium- $t$ target wall



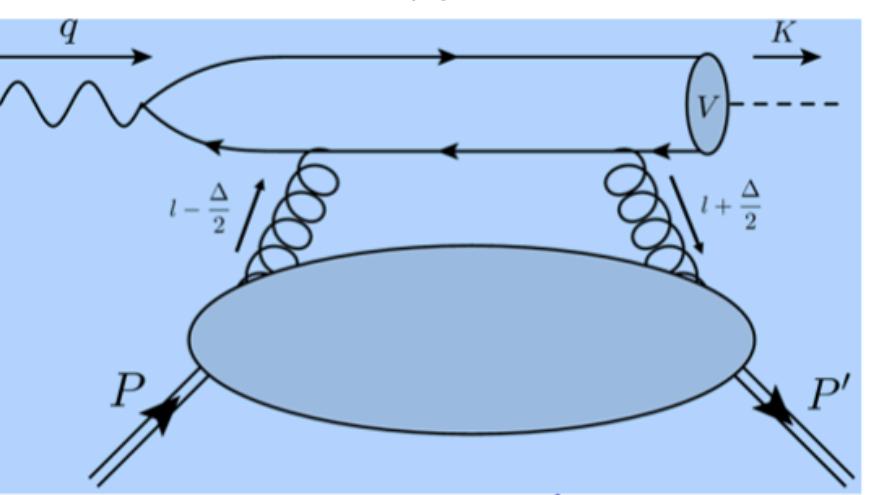
# THE GENERALIZED PARTON DISTRIBUTION MODEL

2D fit to extract  $A(t)$  &  $C(t)$  assuming  $B(t)$  negligible

Y. Guo, X. Ji and Y. Liu, Phys. Rev. D **103**, no.9, 096010 (2021) and Y. Guo, X. Ji and Y. Liu, J. Yang, Phys. Rev. D **108** (2023) no.3, 034003



$$\frac{d\sigma}{dt} = \frac{\alpha_e m e_Q^2}{4(W^2 - m_N^2)^2} \frac{(16\pi\alpha_s)^2}{3M_V^2} |\psi_{NR}|^2 |G(t, \xi)|^2$$



$$G(t, \xi) = \sum_0^\infty \frac{1}{\xi^{2n+2}} \int_{-1}^1 dx x^{2n} F_g(x, \xi, t)$$

$$|G(t, \xi)|^2 = \frac{4}{\xi^4} \left\{ \left(1 - \frac{t}{4m_N^2}\right) E_2^2 - 2E_2(H_2 + E_2) + (1 - \xi^2)(H_2 + E_2)^2 \right\}$$

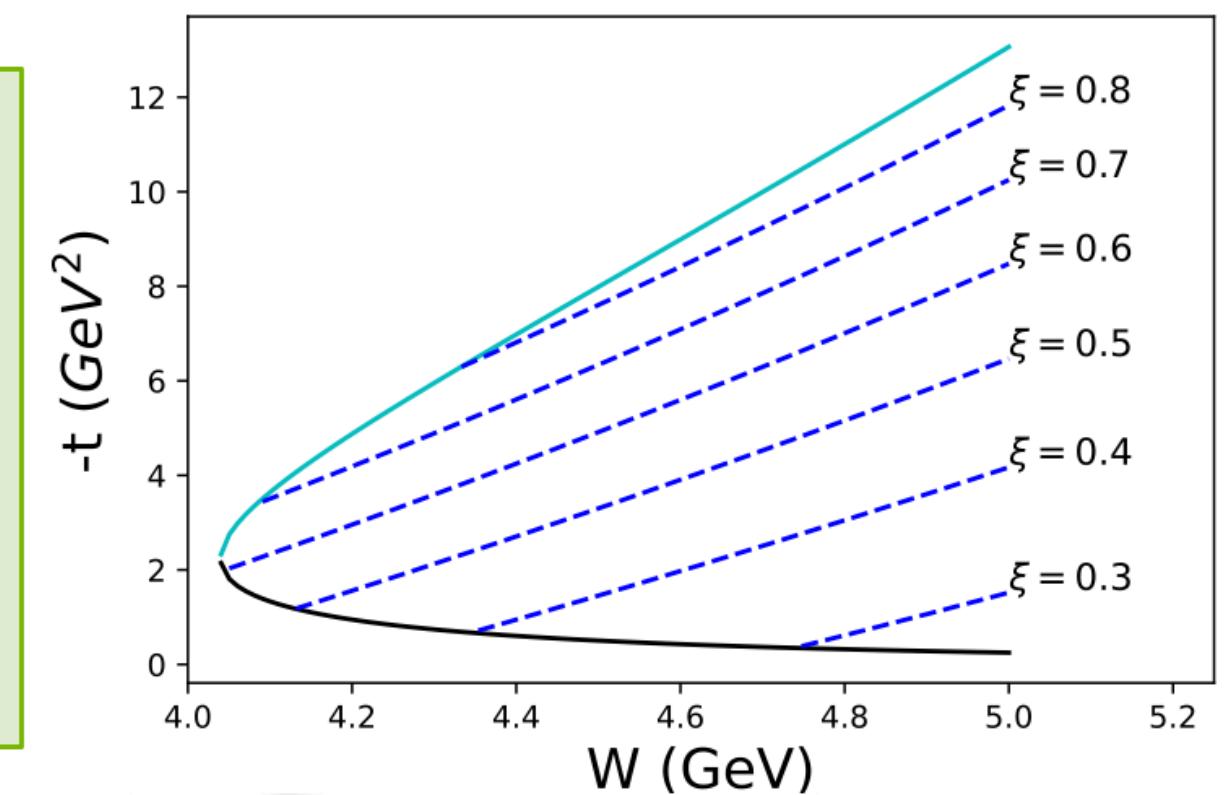
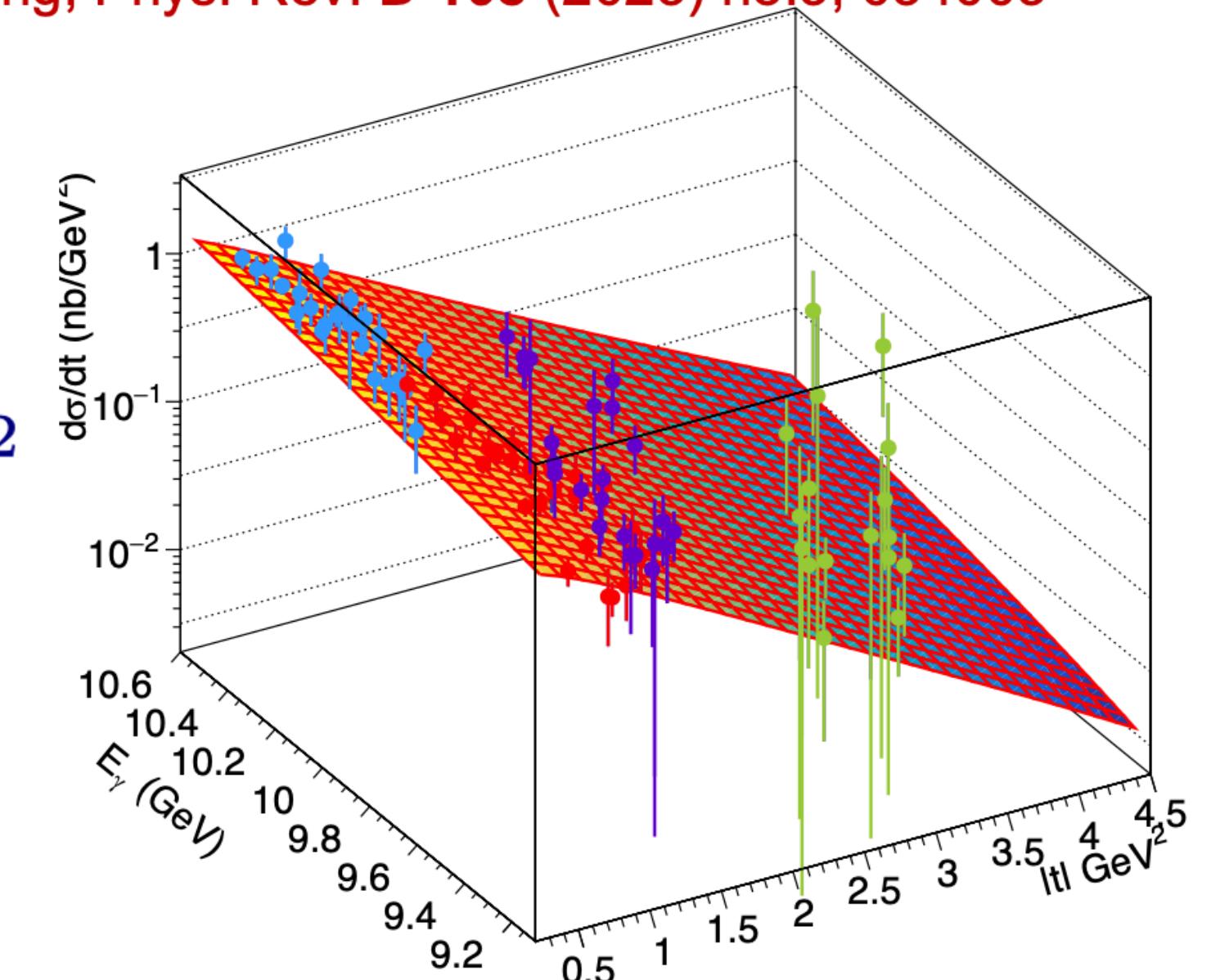
$$\int_0^1 dx H_g(x, \xi, t) = A_{2,0}^g(t) + (2\xi)^2 C_2^g \equiv H_2(t, \xi)$$

$$\int_0^1 dx E_g(x, \xi, t) = B_{2,0}^g(t) - (2\xi)^2 C_2^g \equiv E_2(t, \xi)$$

20

$$A_g(t) = \frac{A_g(0)}{\left(1 - \frac{t}{m_A^2}\right)^3}$$

$$C_g(t) = \frac{C_g(0)}{\left(1 - \frac{t}{m_C^2}\right)^3}$$



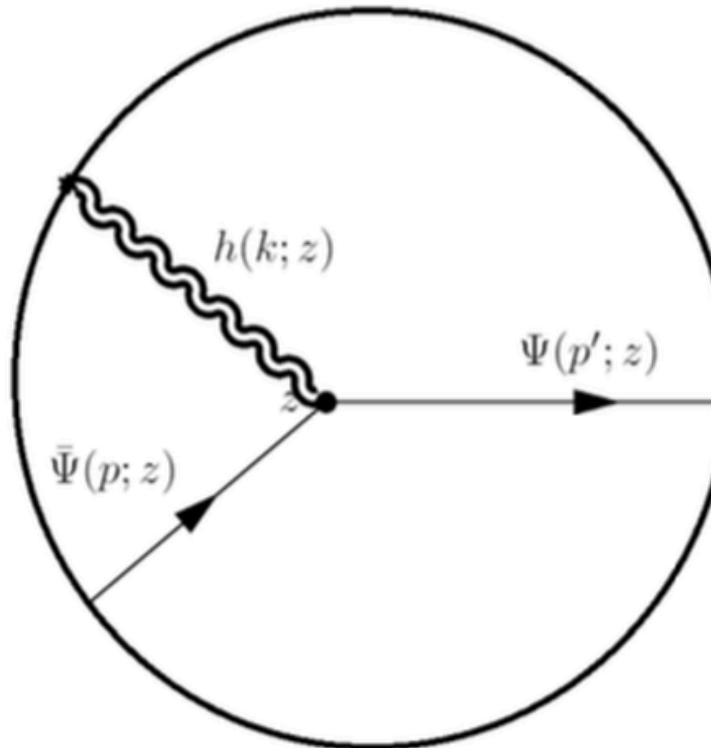
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# THE HOLOGRAPHIC QCD MODEL

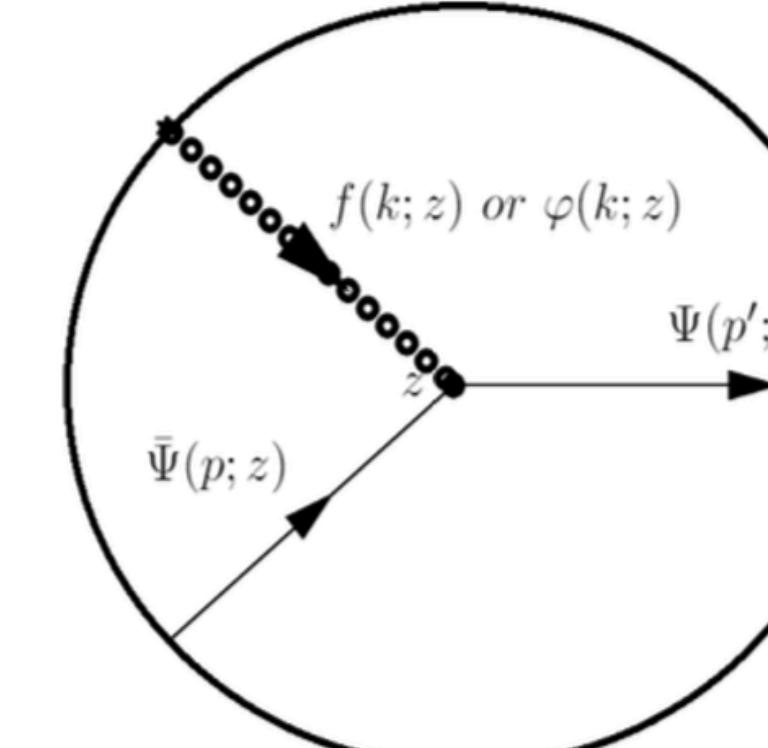
2D fit to extract the  $A(t)$  &  $C(t)$  assuming  $B(t)$  to be small

M-Z: K. Mamo & I. Zahed, PRD 103, 094010 (2021) and 2204.08857 (2022)

A tensor component and a scalar component



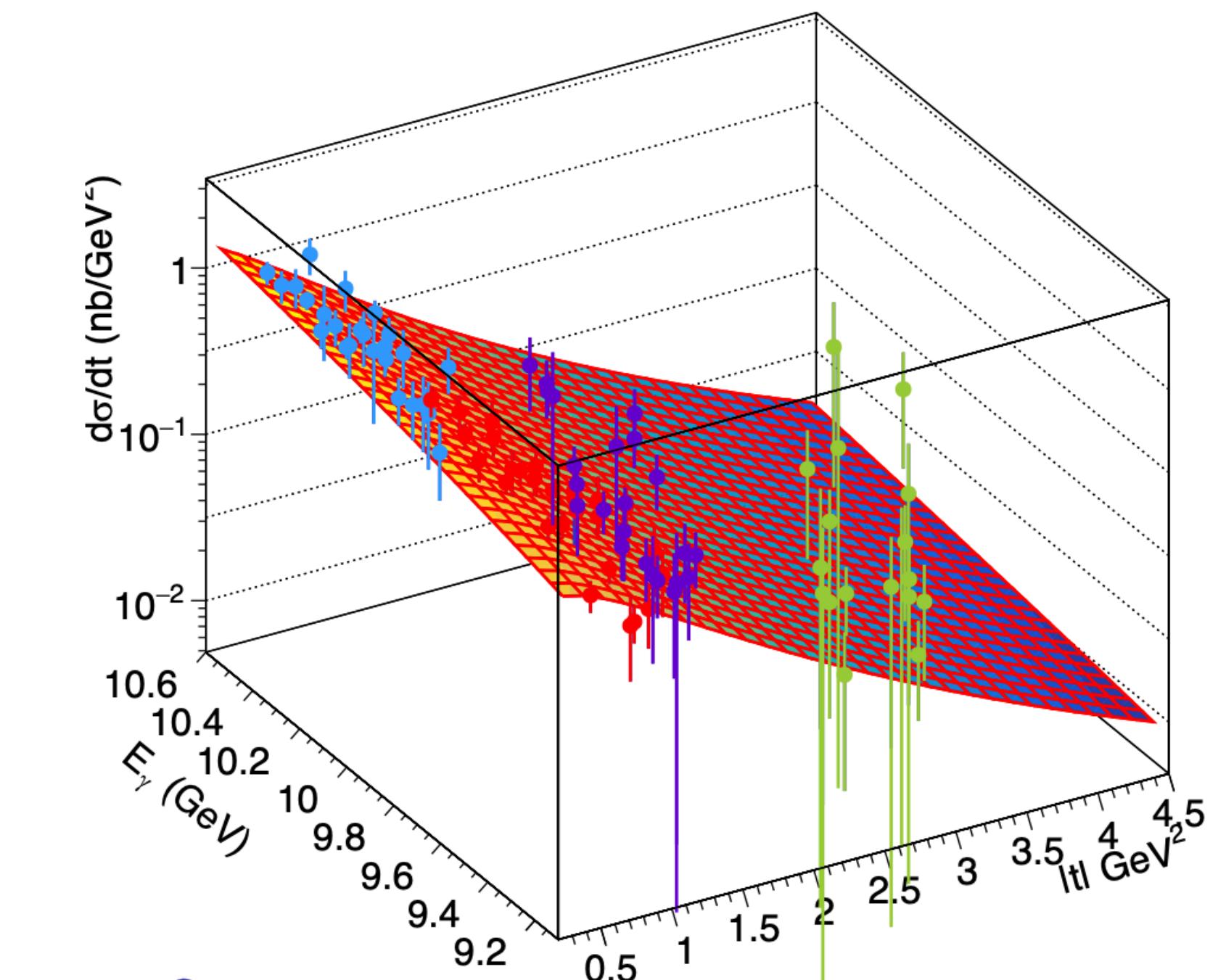
Spin-2 :  $\langle p_2 | T^{xy}(0) | p_1 \rangle$



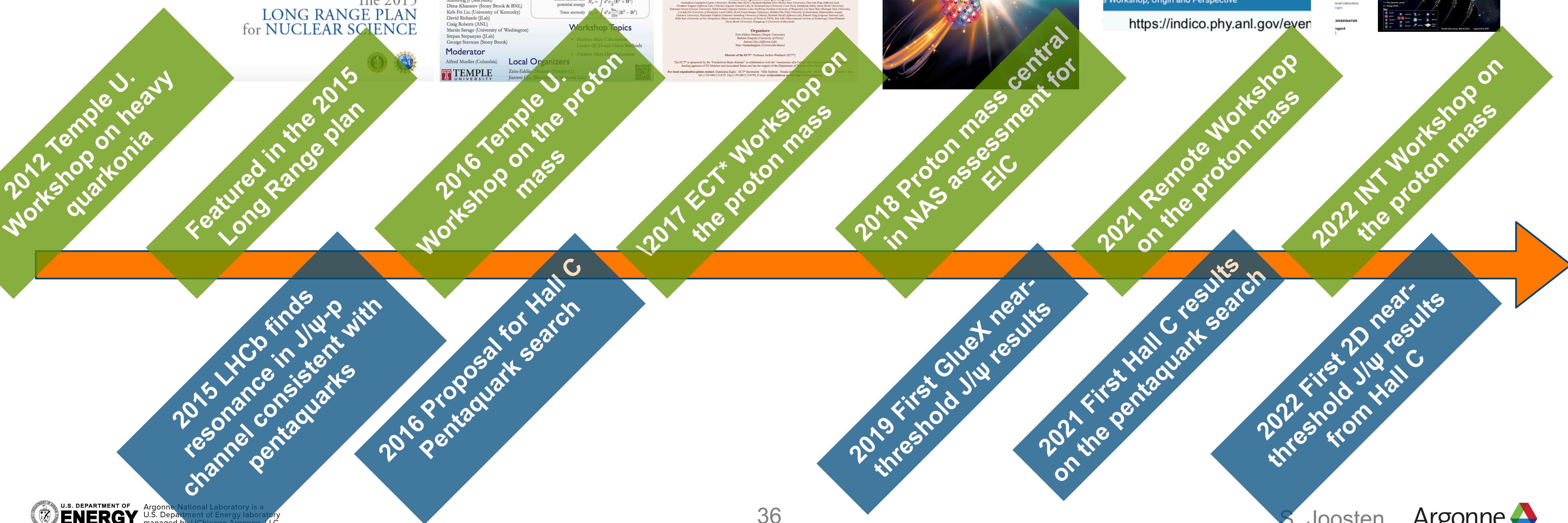
Spin-0 :  $\langle p_2 | T_\mu^\mu(0) | p_1 \rangle$

$$\frac{d\sigma}{dt} = \mathcal{N} \times \frac{e^2}{64\pi(s - m_N^2)^2} \times \frac{A(-t, \kappa_T) + \eta^2 D(-t, \kappa_T, \kappa_S)]^2}{A^2(0)} \times F(\tilde{s}) \times 8$$

- $A(t)$  and  $D(t)$  shapes are fully calculated; However, dipole forms are assumed as very good approximations and are used in the fits to the data.  $A_g(0) = \langle x_q \rangle$  is fixed to the DIS value from global fit CT18.
- $B(t)$  is neglected and  $\mathcal{N}$  is normalized to the cross section.



# The proton mass: An important topic in contemporary hadronic physics! RAPIDLY EVOLVING



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