

TMD: from CLAS12 to EIC

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Research Supported by



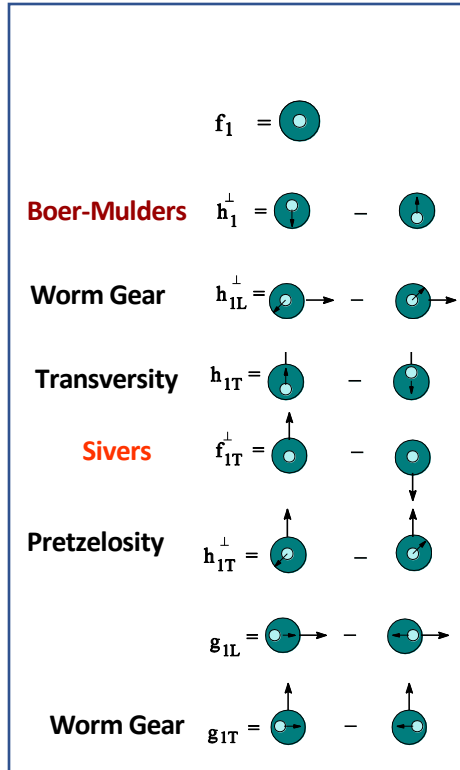
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SIDIS X-section in the Parton Model

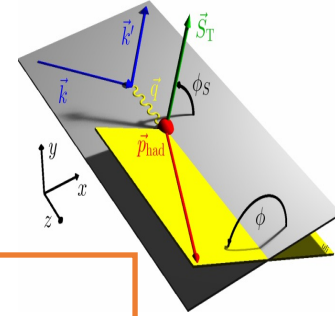


$$d^6\sigma = \frac{4\pi\alpha^2 sx}{Q^4} \times$$

$$\{ [1 + (1-y)^2] \sum e_q^2 f_1^q(x) D_1^q(z, P_{h\perp}^2) + (1-y) \frac{P_{h\perp}^2}{4z^2 M_N M_h} \cos(2\phi_h^l) \sum_{q,\bar{q}} e_q^2 h_1^{\perp(1)q}(x) H_1^{\perp q}(z, P_{h\perp}^2) \}$$

$$\begin{aligned} & - |S_L| (1-y) \frac{P_{h\perp}^2}{4z^2 M_N M_h} \sin(2\phi_h^l) \sum_{q,\bar{q}} e_q^2 h_{1L}^{\perp(1)q}(x) H_1^{\perp q}(z, P_{h\perp}^2) \\ & + |S_T| (1-y) \frac{P_{h\perp}}{zM_h} \sin(\phi_h^l + \phi_S^l) \sum_{q,\bar{q}} e_q^2 h_1^q(x) H_1^{\perp q}(z, P_{h\perp}^2) \\ & + |S_T| (1-y + \frac{1}{2}y^2) \frac{P_{h\perp}}{zM_N} \sin(\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 f_{1T}^{\perp(1)q}(x) D_1^q(z, P_{h\perp}^2) \\ & + |S_T| (1-y) \frac{P_{h\perp}^3}{6z^3 M_N^2 M_h} \sin(3\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 h_{1T}^{\perp(2)q}(x) H_1^{\perp q}(z, P_{h\perp}^2) \\ & + \lambda_e |S_L| y(1 - \frac{1}{2}y) \sum_{q,\bar{q}} e_q^2 g_1^q(x) D_1^q(z, P_{h\perp}^2) \end{aligned}$$

$$+ \lambda_e |S_T| y(1 - \frac{1}{2}y) \frac{P_{h\perp}}{zM_N} \cos(\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 g_{1T}^{\perp(1)q}(x) D_1^q(z, P_{h\perp}^2) \}$$



Unpolarized

Polarized target

Polarized beam and target

S_L and S_T : Target Polarizations; λ_e : Beam Polarization
 x : momentum fraction carried by struck quark, z : fractional energy of hadron

SIDIS cross-section

$$\begin{aligned}
 & \frac{d\sigma}{dx dy d\phi_S dz d\phi_h dP_{h\perp}^2} \\
 &= \frac{\alpha^2}{x y Q^2} \frac{y^2}{2(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 &+ \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} + S_L \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 &+ S_L \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 &+ S_T \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \right. \\
 &+ \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} \\
 &+ \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] + S_T \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 &+ \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}
 \end{aligned}$$

- Disentangling the different contributions is not trivial

- Ratio of T to L flux

- At fixed x e.g. change Q

$$\varepsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2},$$

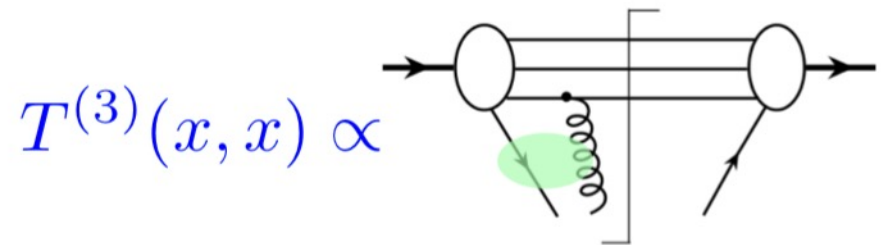
$$\gamma = \frac{2Mx}{Q}.$$

Beyond the parton picture

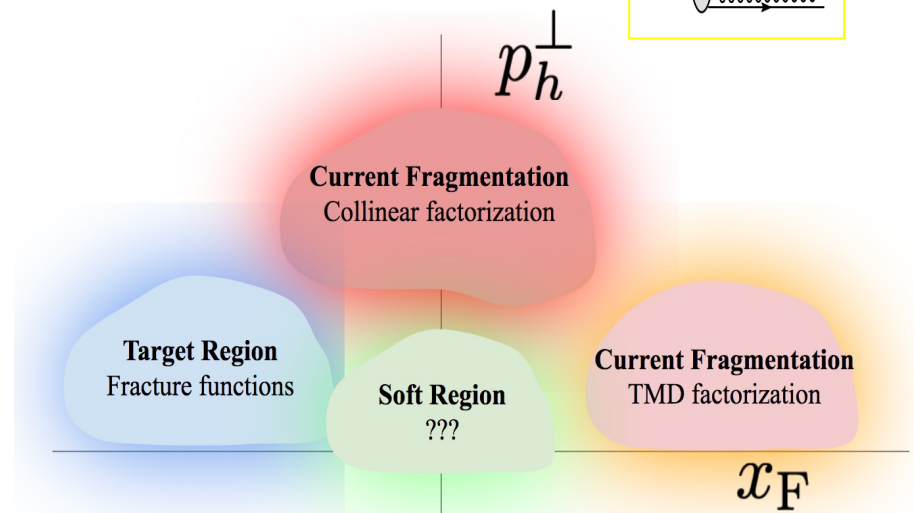
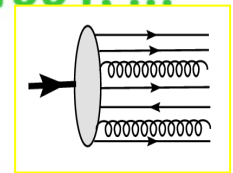
- Higher Twist Contributions
- Overlap of regions that are not captured by factorized TMD picture
- VM Meson decays
- Radiative effects
- Assumption of suppressed long. photon contributions

One persons 'complication' is another person's signal...

→ Need high lumi, leverarm in kinematics to disentangle various contributions



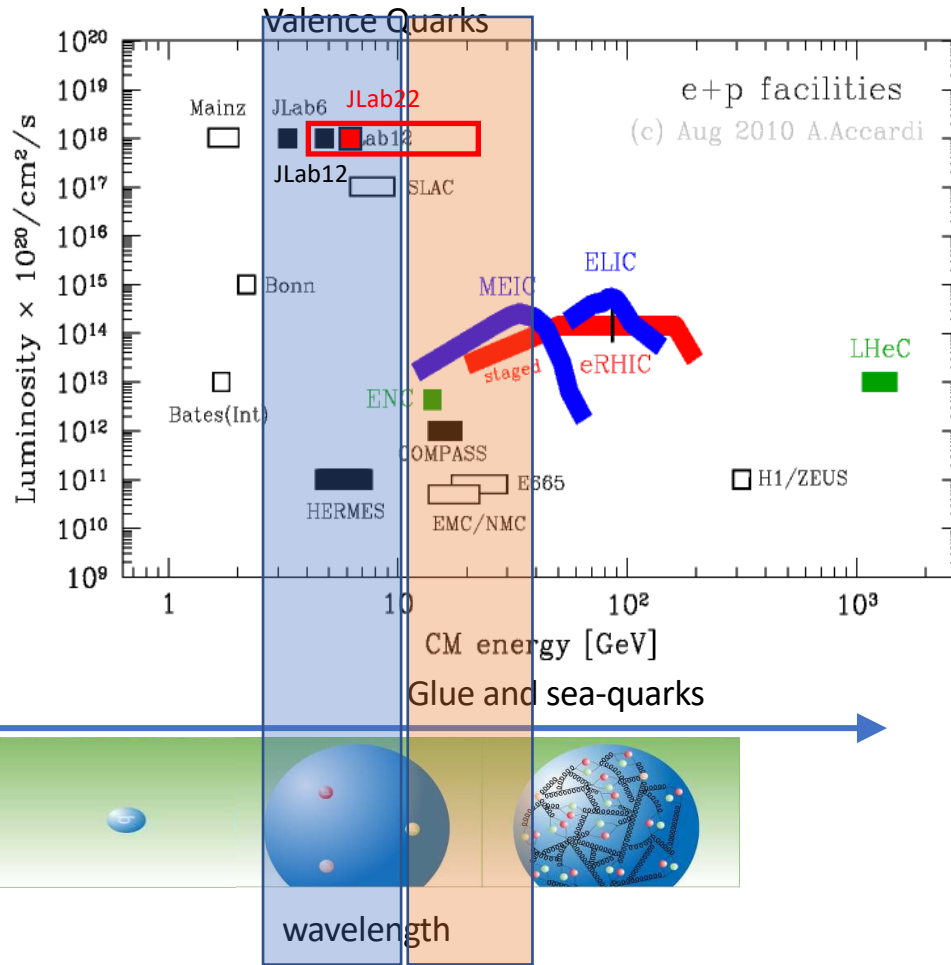
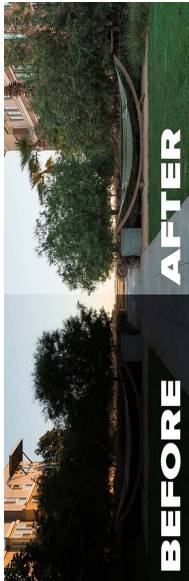
Qiu, Sterman, 1991....



Boglione, Produkin et al

Landscape

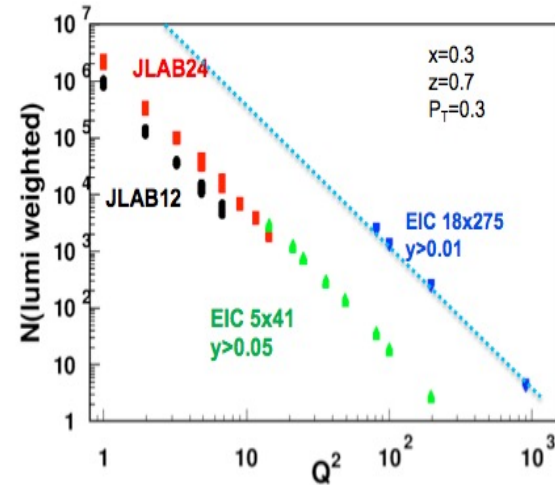
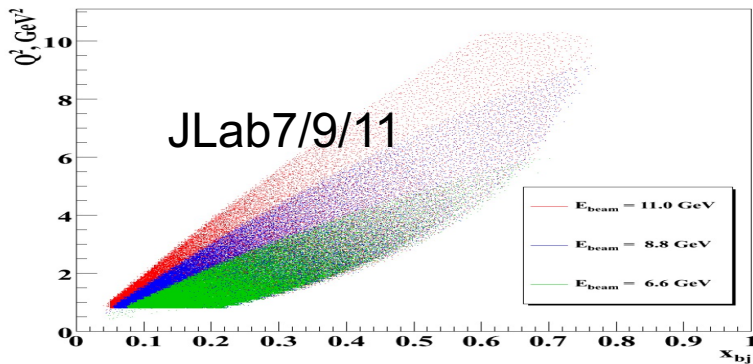
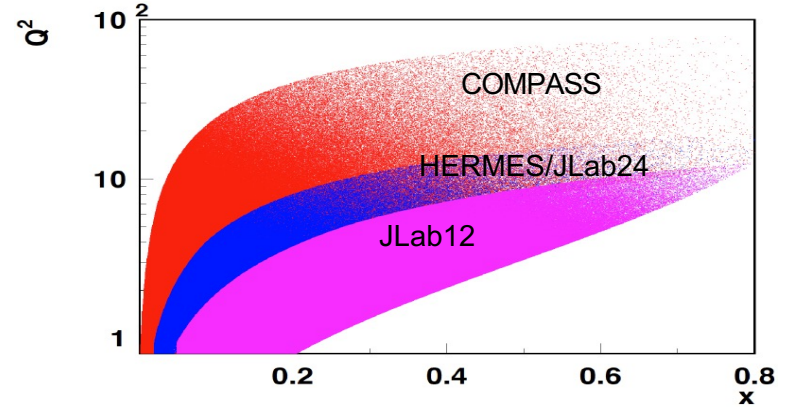
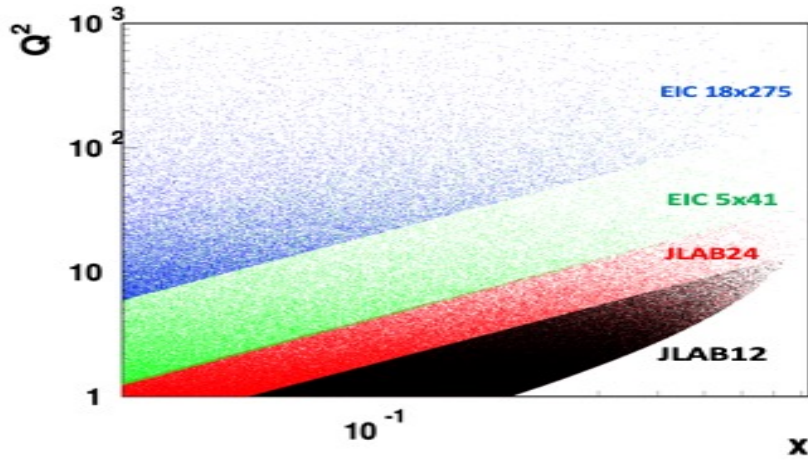
Exposure ↑



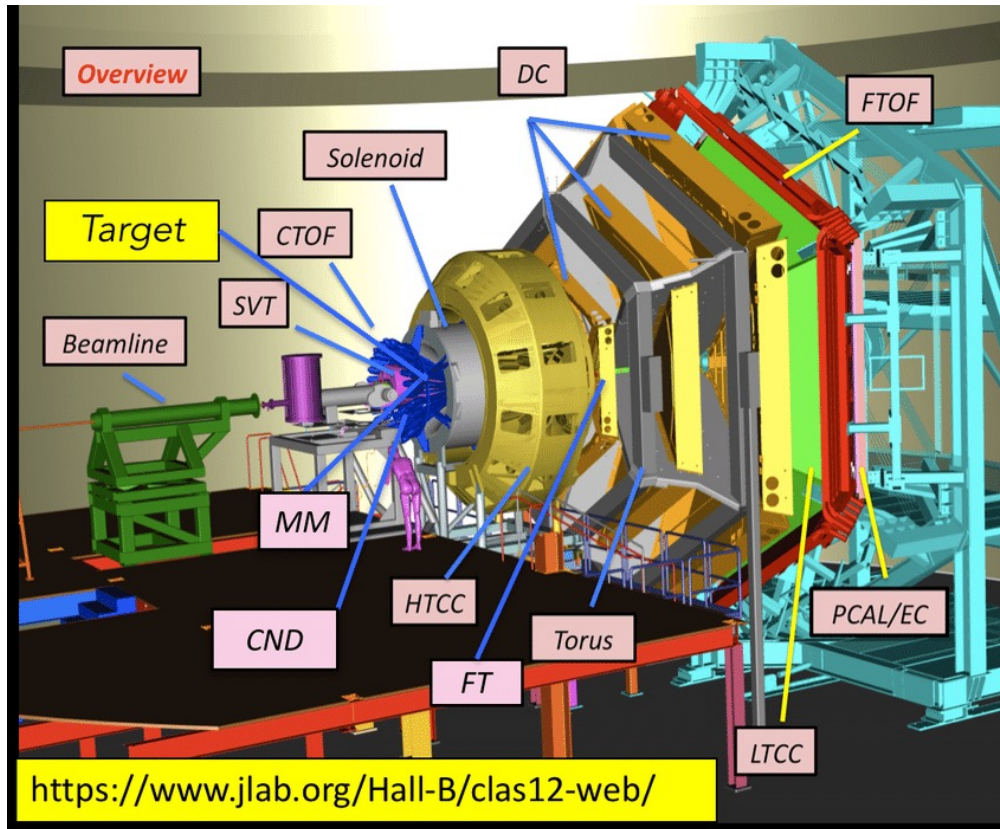
To extract 3D structure
Need luminosity

- Jlab12: 2018+
- JLab22:

Kinematic comparisons



NB: Kinematic slice heavily biased towards Jlab



Longitudinally Polarized Electron Beam

- $E = 10.6 \text{ GeV}$
- $P = 86\text{--}89\%$

Unpolarized Liquid H_2 Fixed Target

- Torus magnet \rightarrow electrons inbending

Run Group A (Unpolarized LH_2 target)

- ★ unpolarized SIDIS cross section off proton
- ★ A_{LU} in Beam Spin Asymmetries

Run Group B (Unpolarized LD_2 target)

- ★ Complementary to RG-A
 \rightarrow allow for u/d quark flavor separation

Run Group C (longitudinally polarized NH_3 and ND_3)

- ★ F_{UL} and F_{LL}

Run Group K (Unpolarized LH_2 target)

- ★ - 6.5, 7.5, 8.4 GeV e- beam
- ★ $F_{UU,L}, F_{UU,T}$ Separation

Run Group H (transversely polarized NH_3)

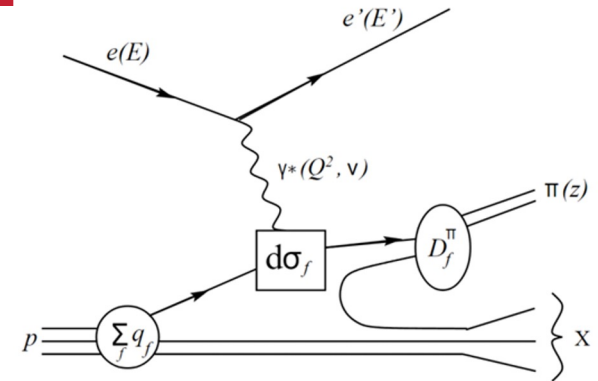
- ★ F_{UT} structure function

\Rightarrow should be completed before 1st Shutdown \approx 2032

Unpolarized Multiplicities of $ep \rightarrow e\pi^0(X)$

- ★ Measurements of neutral pion multiplicities
 - π^0 yields normalized by number of DIS electrons

$$\sigma^{\pi^0} \approx \sigma^{DIS} \otimes f^p(x, Q^2) \otimes D^{p \rightarrow \pi^0}(z, Q^2)$$

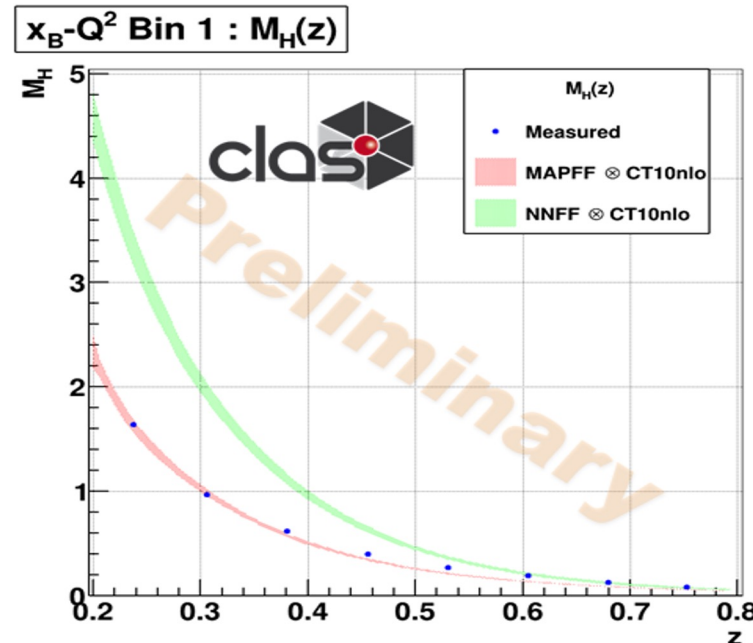


- ★ Study integrates over the azimuthal ϕ_h angle

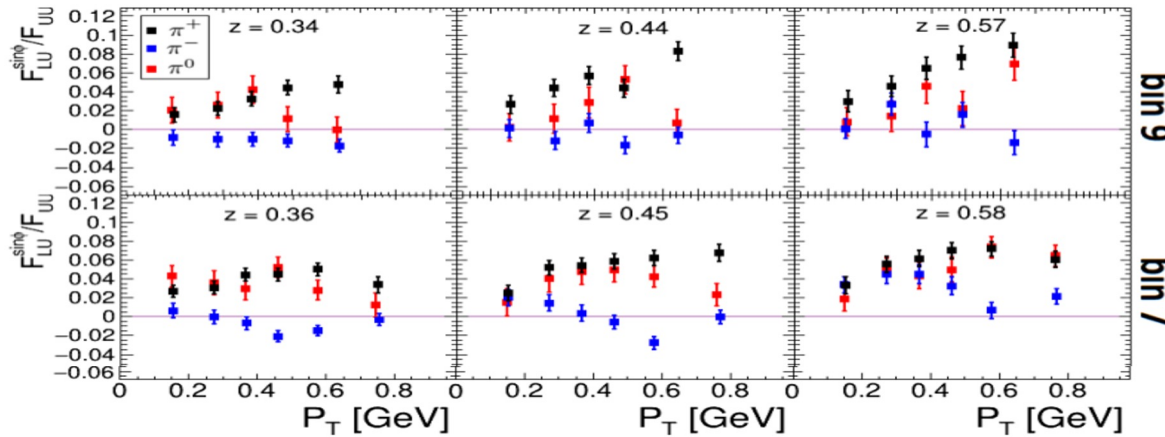
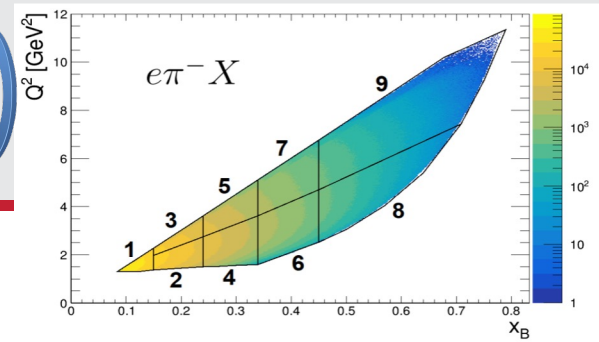
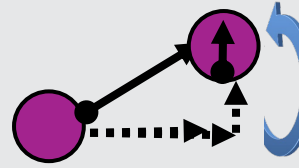
$$F_{UU,T} = C[f_1 D_1] \quad D_1^{\pi^0/q} = \frac{1}{2} \left(D_1^{\pi^+/q} + D_1^{\pi^-/q} \right)$$

- ★ Invariant mass fits over the diphoton spectrum are performed to calculate $N(\pi^0)$

- ★ **Ongoing Work:** Bayesian unfolding, ϕ_h modulation fits



CLAS12 pion BSAs



Higher Twist PDFs

N/q	U	L	T
U	f^\perp	g^\perp	h, e
L	f_L^\perp	g_L^\perp	h_L, e_L
T	f_T, f_T^\perp	g_T, g_T^\perp	$h_T, e_T, h_T^\perp, e_T^\perp$

S. Diehl

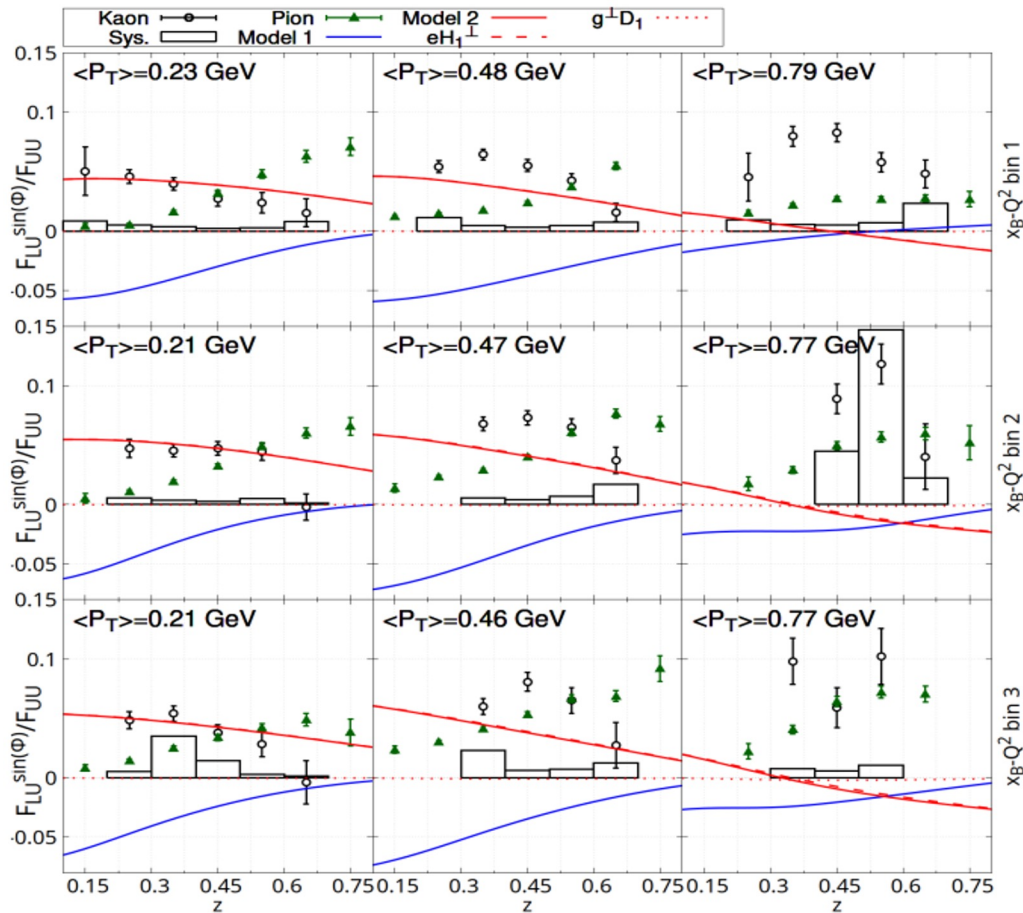
$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q} \mathcal{C} \left[-\frac{\hat{h} \cdot k_T}{M_h} \left(x e \boxed{H_1^\perp} + \frac{M_h}{M} \boxed{f_1} \frac{\boxed{\tilde{G}^\perp}}{z} \right) + \frac{\hat{h} \cdot p_T}{M} \left(x \boxed{g^\perp} \boxed{D_1} + \frac{M_h}{M} \boxed{h_1^\perp} \frac{\boxed{\tilde{E}}}{z} \right) \right]$$

twist-3 pdf unpolarized PDF twist-3 t-odd PDF Boer-Mulders
Collins FF twist-3 FF unpolarized FF

★ If Collins term only (H_1^\perp) \rightarrow hierarchy of the A_{LU} 's
 $A_{LU}(\pi^-) < A_{LU}(\pi^0) = 0 < A_{LU}(\pi^+)$

★ Observed is more **Sivers-like** (g^\perp), asymmetry comes from struck u-quark
 $A_{LU}(\pi^-) < A_{LU}(\pi^0) = A_{LU}(\pi^+)$

Kaon Asymmetries are larger



A. Kripko

- Reasonable Assumption

- u quark dominance

- Difference due to $D_1^{K^+}/u$, $H_1^{\perp K^+}/u$

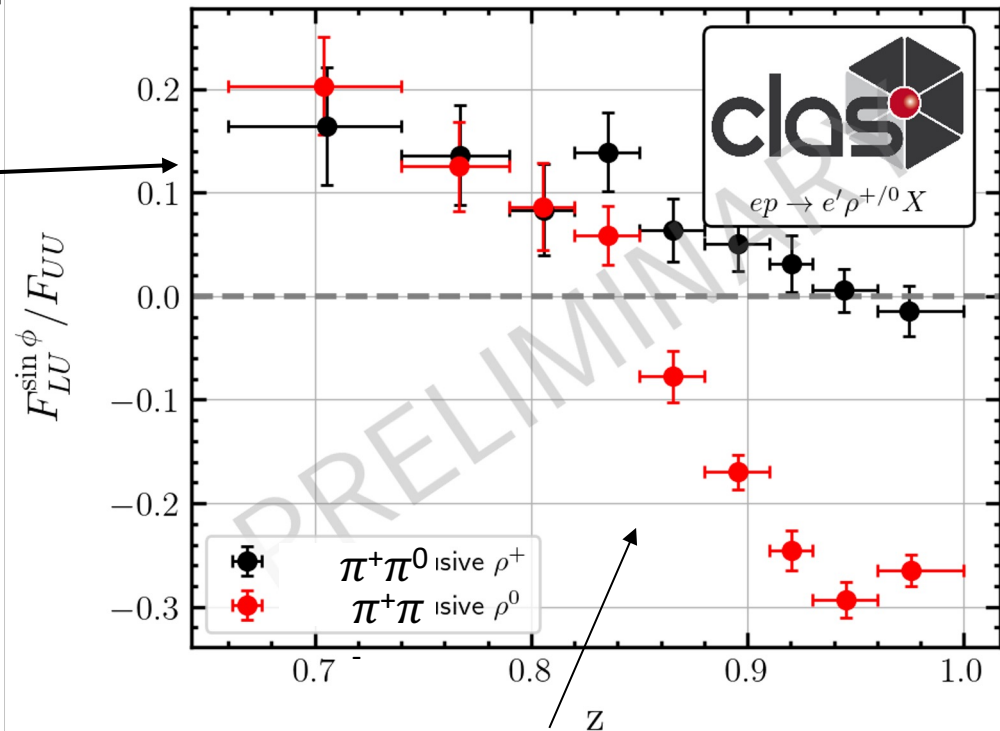
Models by Mao/Lu
 using different models
 of $e(x)$, g^\perp
 EPJC 73, 2557 (2013) and 74, 2910 (2014)

Near-exclusive $\pi^+\pi^-$, $\pi^+\pi^0$ production

★ We can constrain/better understand the contribution of ρ^0 , ρ^+ decays on our single hadron asymmetries by looking at near exclusive ($M_X < 1.1$ GeV) channels

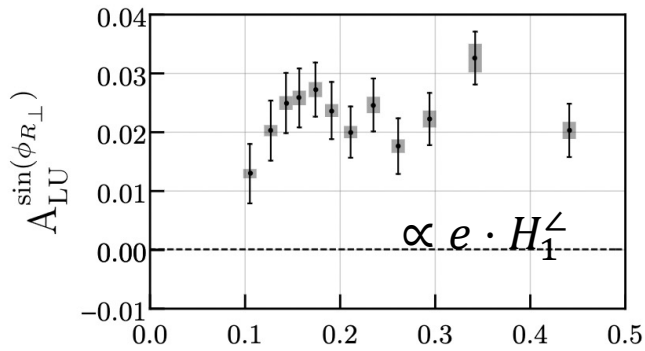
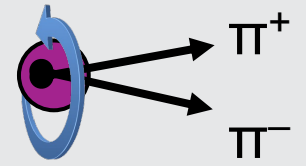
★ Strong yet similar asymmetries observed (**both productions came from struck u quark**)

→ See talk by K. Joo

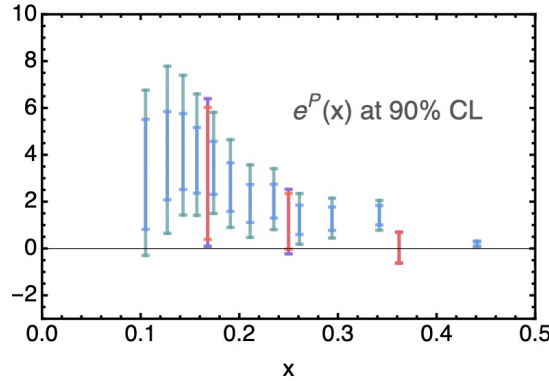


★ Different mechanism for neutral ρ^0 at high z (low $|t|$) → GPDs, gluon contributions

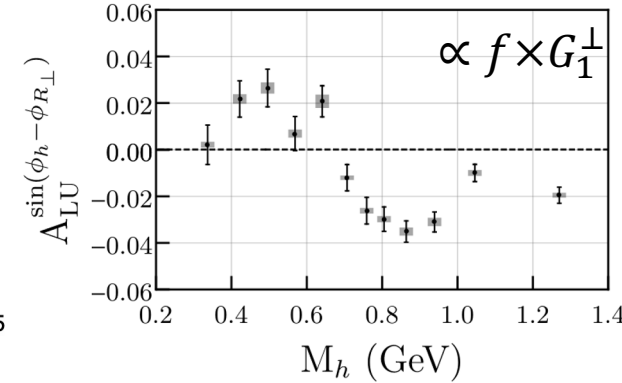
Better: di-hadrons



Phys.Rev.Lett. 126 (2021) 152501



Phys.Rev.D 106 (2022) 1, 014027

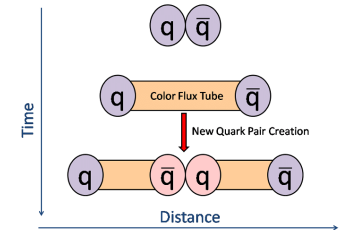


$$F_{LU}^{\sin \phi_R} = -x \frac{|\mathbf{R}| \sin \theta}{Q} \left[\frac{M}{m_{hh}} x e^q(x) H_1^{\triangleleft q}(z, \cos \theta, m_{hh}) + \frac{1}{z} f_1^q(x) \tilde{G}^{\triangleleft q}(z, \cos \theta, m_{hh}) \right],$$

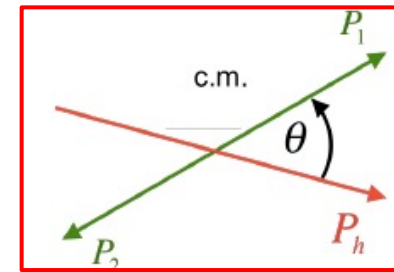
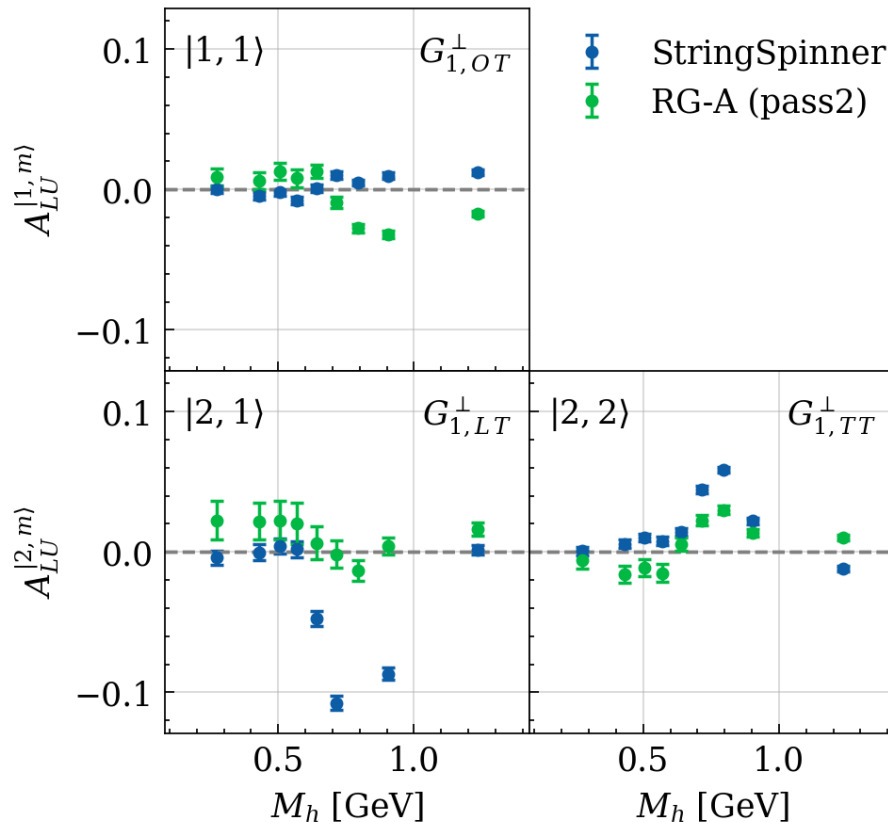
- First extraction of $e(x)$
- Further constrains from F_{UL} and F_{LL}

Compare Partial Wave Decomposition in MC and Data

- Comparing to Polarized Lund model here (StringSpinner $3P_0$ model, A. Kerbizi et al, *Comput.Phys.Commun.* 272 (2022))



Twist-2 A_{LU} Amplitudes

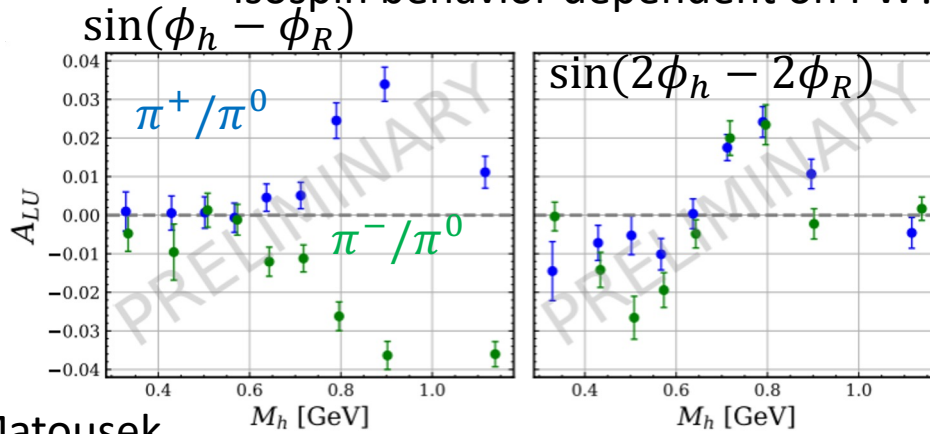


- See more MC tuning studies in QCD whitepaper
- E.g. charge, flavor correlations (Phys.Rev.D 105 (2022) 5, L051502)

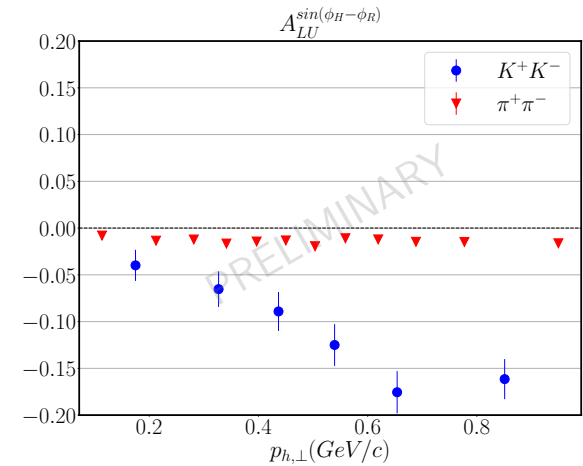
G Matousek

π^0 and Kaon combinations (SIDIS@CLAS12)

Isospin behavior dependent on PW?



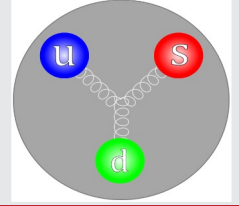
G Matousek



C Pecar

- $A_{LU} \propto \frac{f(x,k_t)G_1^\perp(z,p_t)}{f(x,k_t)D(z,p_t)} \approx \frac{G_1^\perp}{D_1}$
- Kaon \gg Pions for sp interference (not all PW terms)
 - FF effect?
 - π^\pm/π^0 ordering dependent on PW

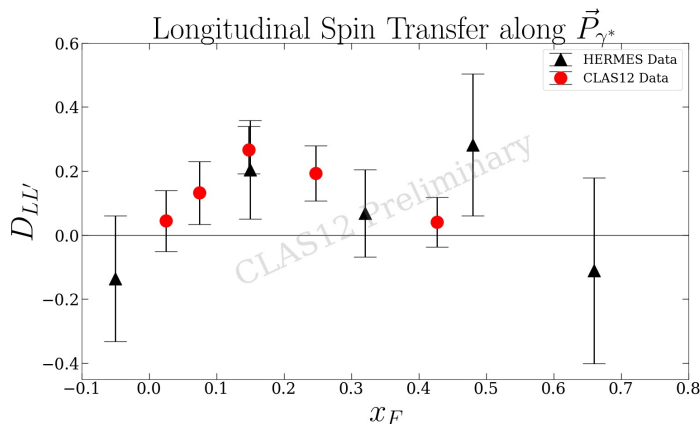
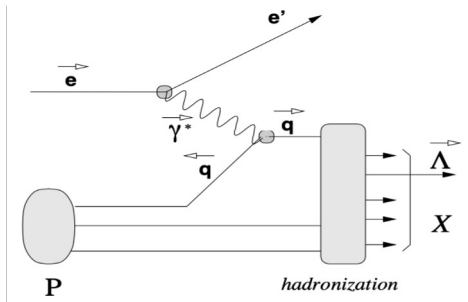
Lambda Program at CLAS12



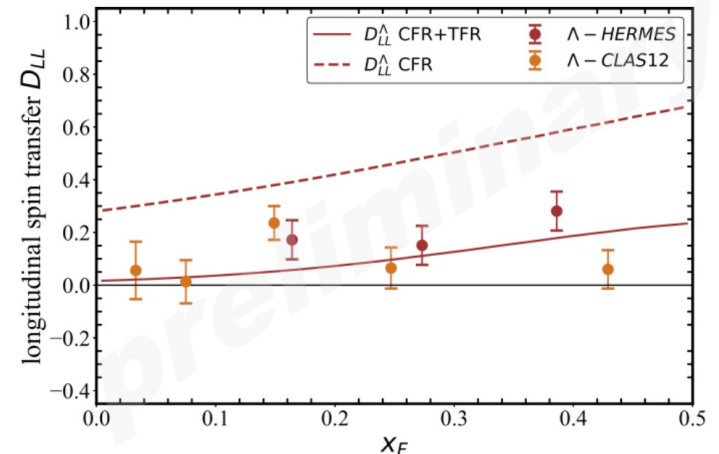
- **Constituent Quark Model (CQM)**
 - Predicts s quark carries 100% of the Λ hyperon spin
- “Do polarized u -quarks from current fragmentation transfer their longitudinal spin to the lambda?” → Test spin structure

$$P_{\Lambda} = P_b D(y) D_{LL'}^{\Lambda},$$

longitudinal spin-transfer



Xiaoyan Zhao at

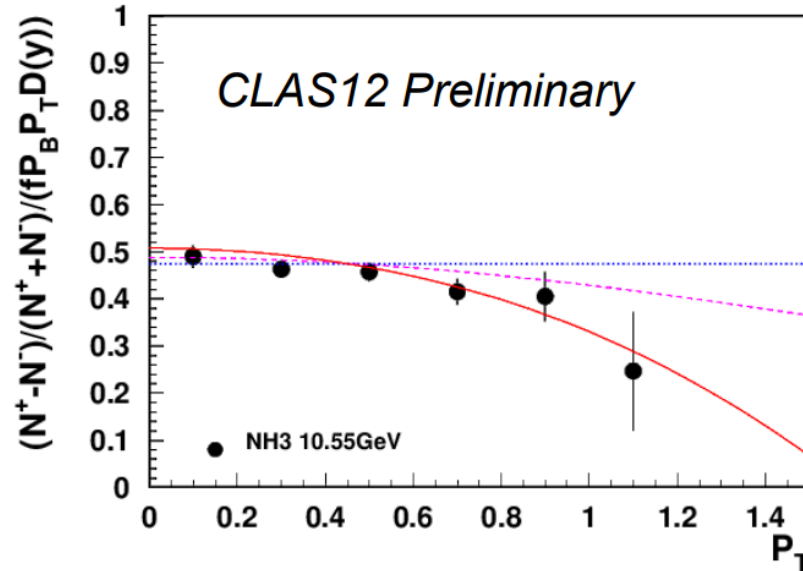


M. McEneaney

Part of planned extensive Lambda program with larger statistics: Transverse, polarizing...

Longitudinal target results with RGC

- Results represent 5% proton target (Ammonia, NH_3)



Sensitive to the k_T width of f_1 and g_1

- Dilution factor $\approx \frac{3}{17}$
- Polarization $\approx 85\%$

$$F_{LL} \propto g_1(x, k_T) \otimes D_1(z, p_T)$$

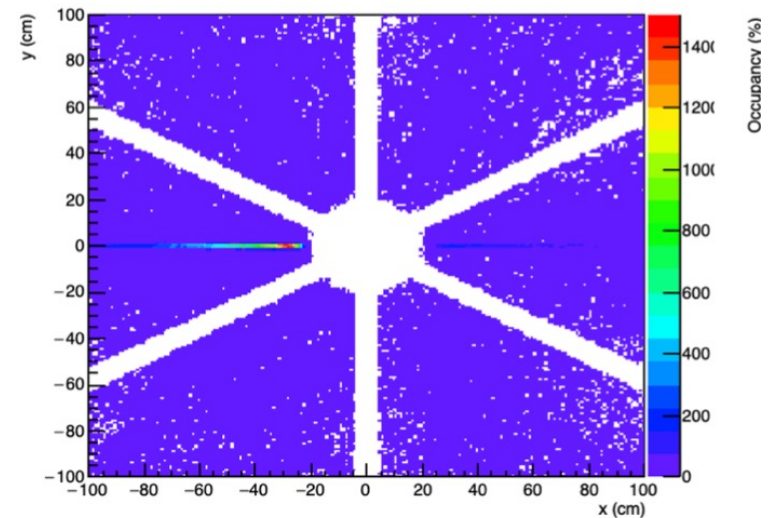
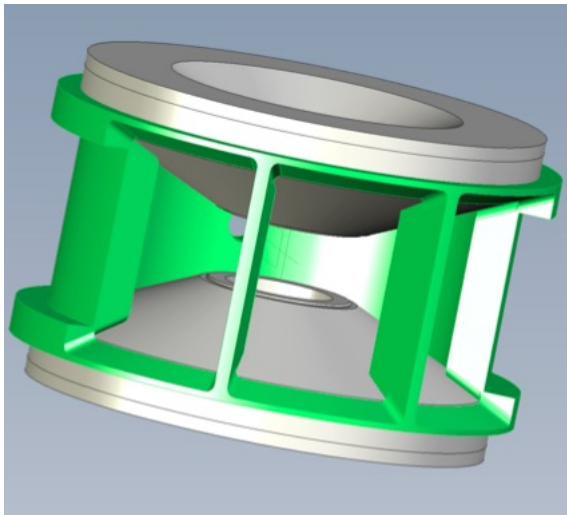
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Convolution over transverse momentum space

- Rich program underway

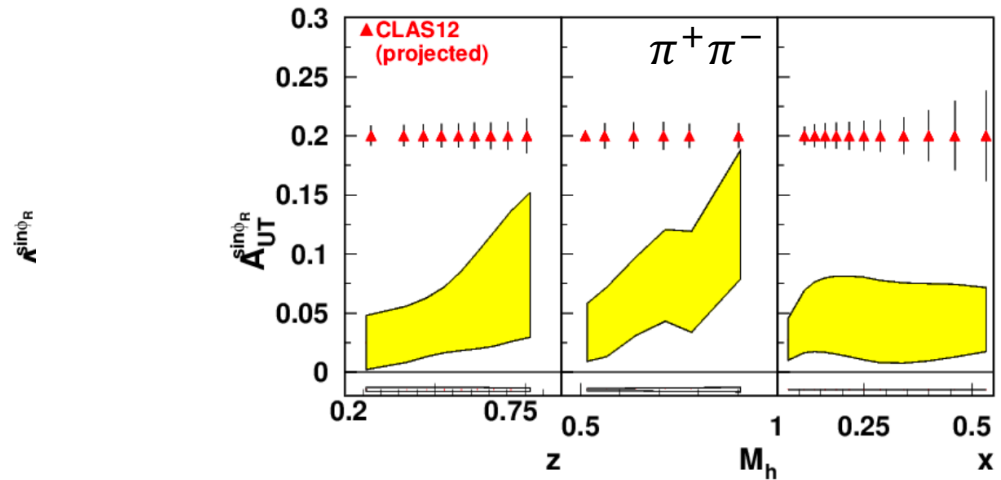
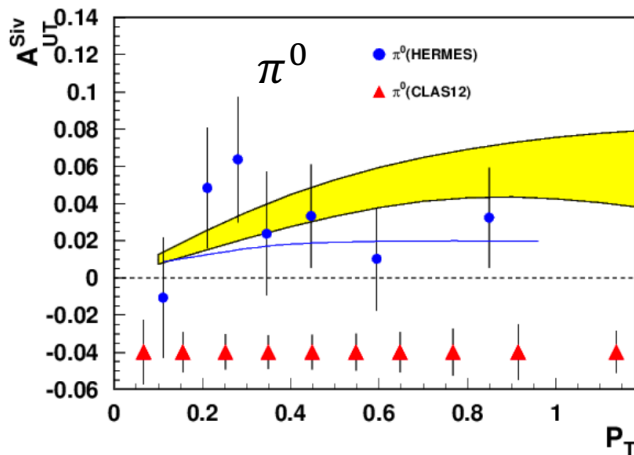
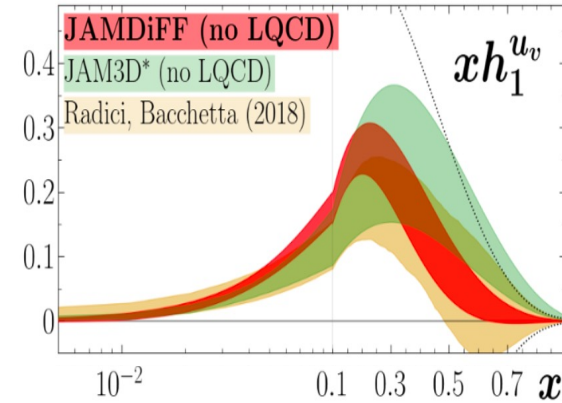
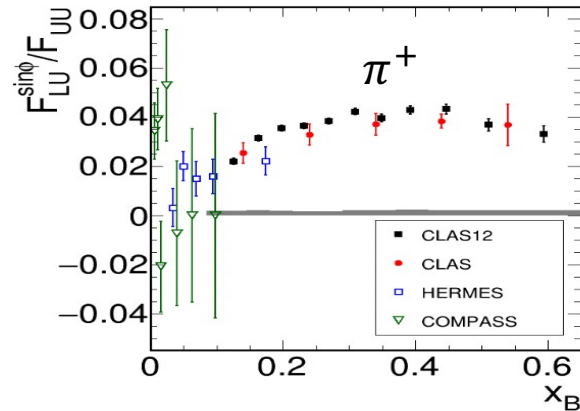
CLAS future: Transverse Spin NH_3 target (RGH)

- Add "conventional" NH_3 target to CLAS
 - Similar to longitudinal target, dilution 3/17, Polarization $\approx 85\%$
 - Transverse Holding field: Moeller Scattering limits luminosity
- ➔ For proposal use **very** conservative $5 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
(about $\frac{1}{50}$ of 'regular' CLAS12)

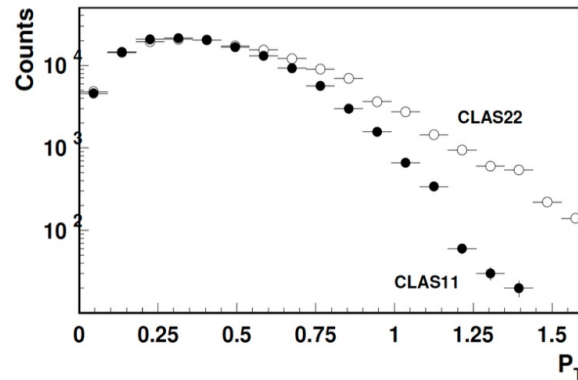
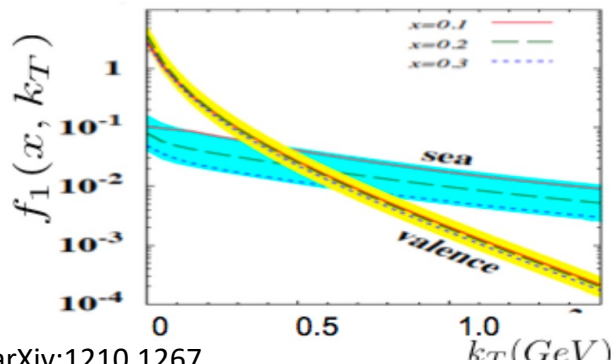
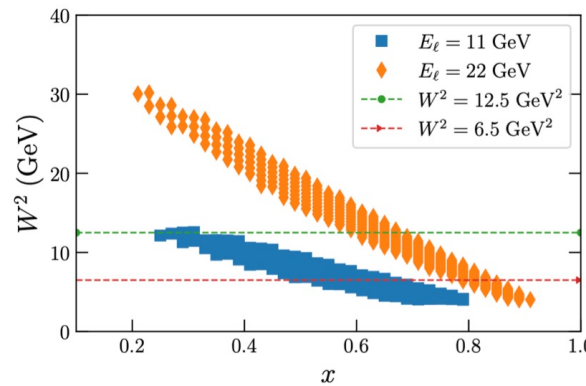
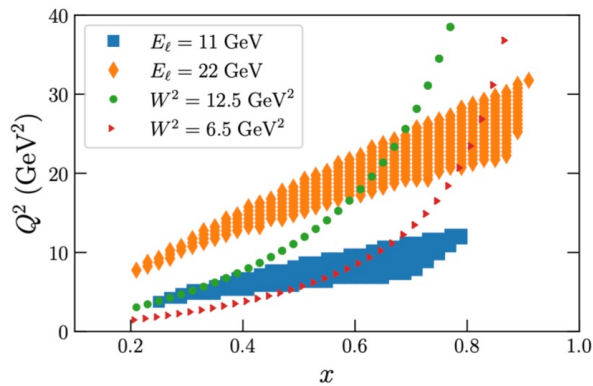


Physics with a transverse Target: Example Transversity

- Data $x > 0.2$ very sparse
→ Models diverge



High x at Jlab 22



P.Schweitzer et al. arXiv:1210.1267

Strong Interaction Physics at the Luminosity Frontier
with 22 GeV Electrons at Jefferson Lab

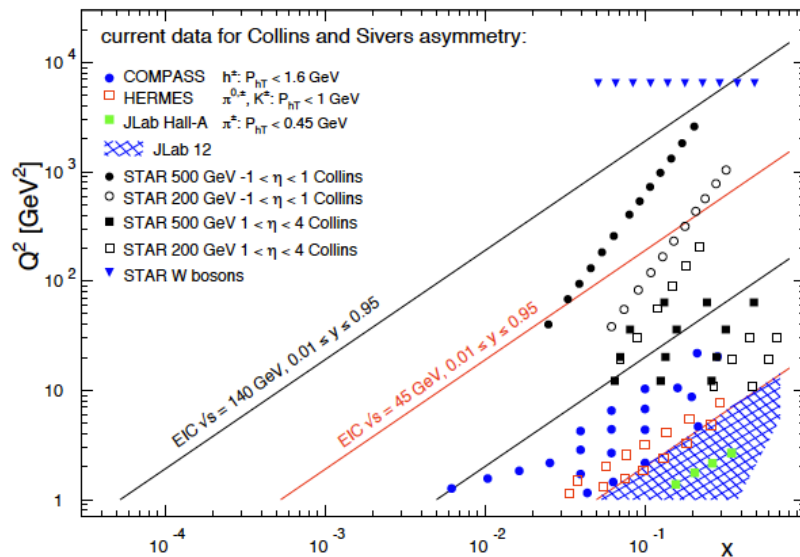
t: [2306.09360](https://www.jlab.org/programs/accelerator/2306.09360) [nucl-ex]

- \approx doubling beam energy significantly increases phase space
- Pin down valence structure of the proton
- Integration in global analyses (e.g., strange distributions, CS Kernel)

SIDIS physics at an EIC: Coverage

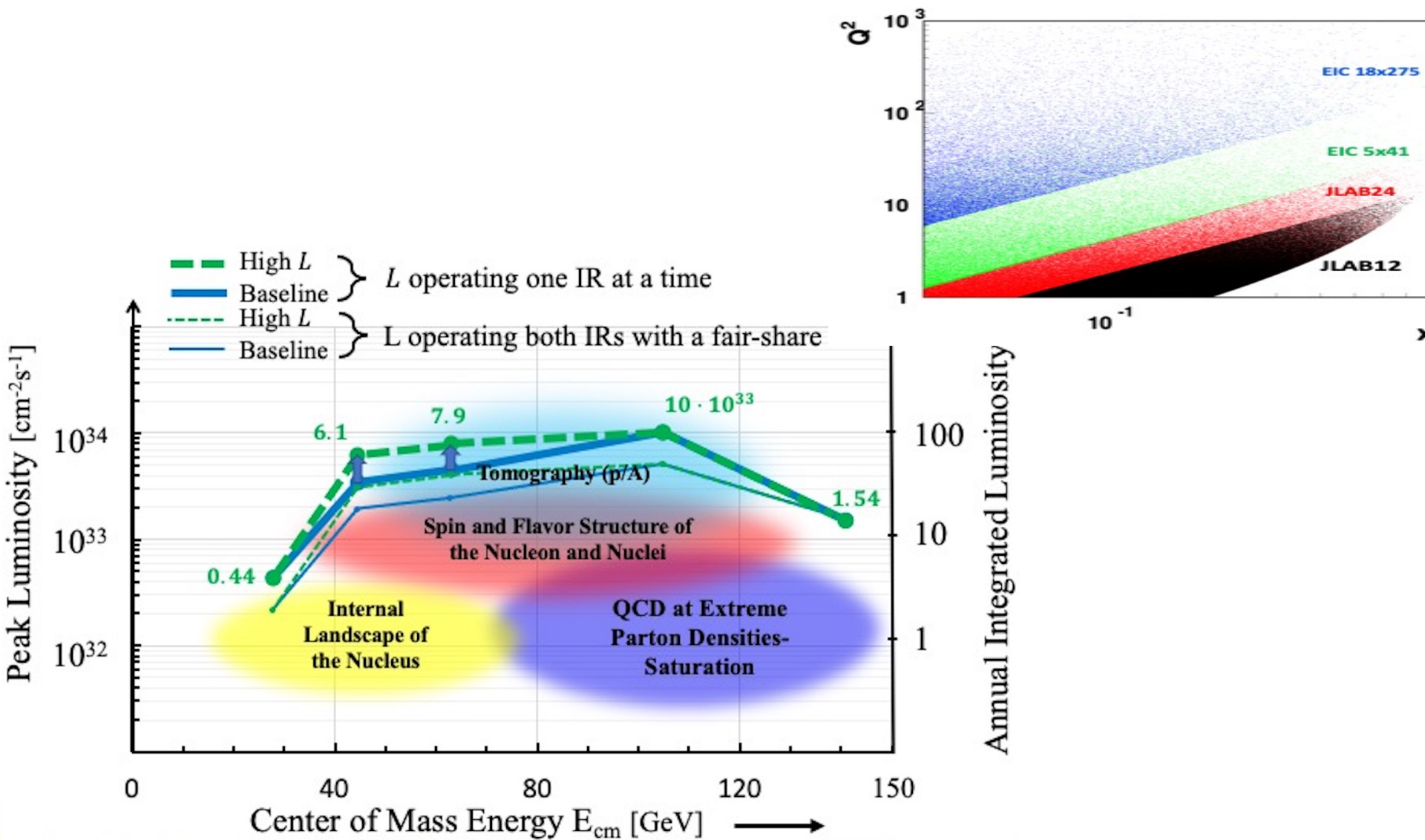
- Common theme on EIC impact
 - Extended **kinematic coverage** and **precision**, along with polarization and possible beam charge degrees of freedom allow multi-pronged approach → needed to extract multidimensional objects
 - TMD factorization is valid

Large Q^2 lever arm: probe evolution, disentangle contributions to σ



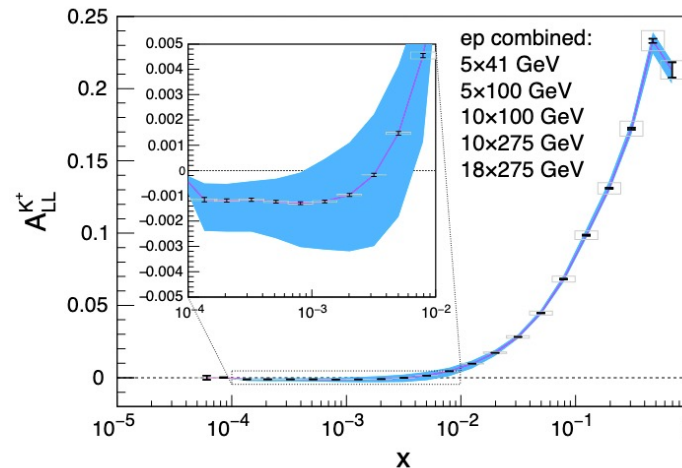
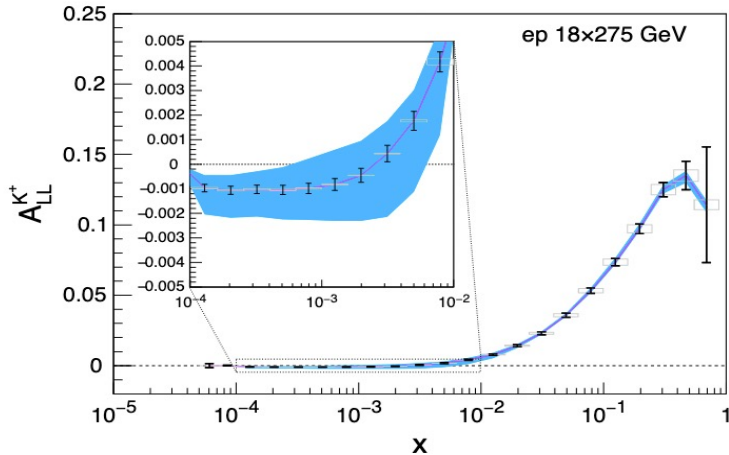
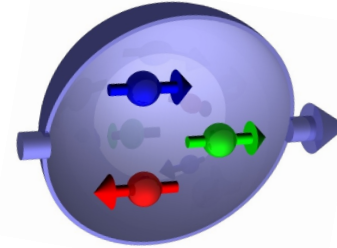
Coverage to low x : access sea and gluon distributions

Order of magnitude in luminosity depending on \sqrt{s} (beware of projections with fixed $\int L$)

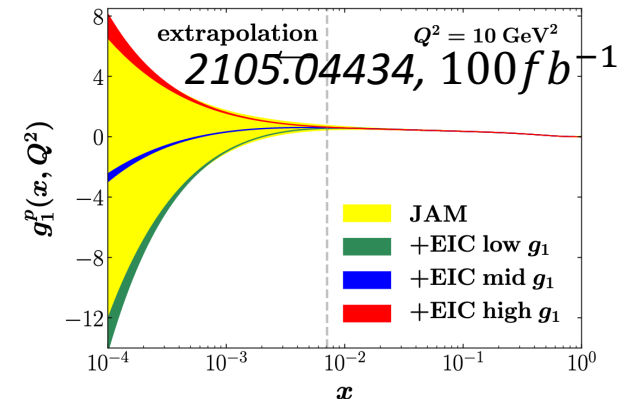


Longitudinal double spin asymmetries

$$\bullet A_{LL} = \frac{\sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\uparrow\downarrow}} \propto g_1$$



- Projections for Athena (2022 *JINST* 17 P10019)
- 3% point-to-point, 2% scale uncertainties (from Hera experience)
- $z > 0.2$
- 15.5 fb^{-1} at $18x275$, other datasets scaled accordingly

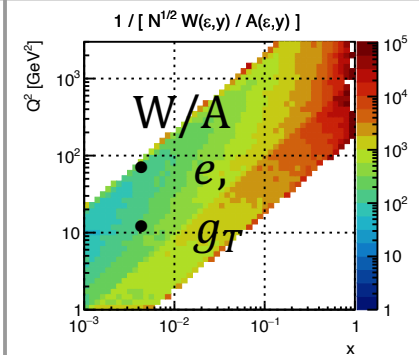
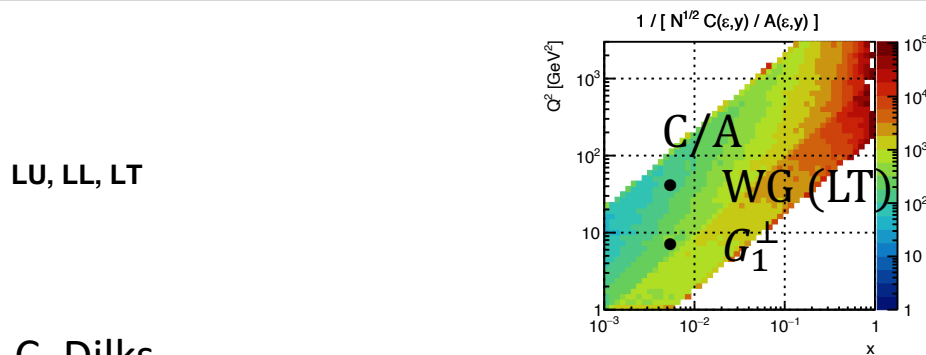
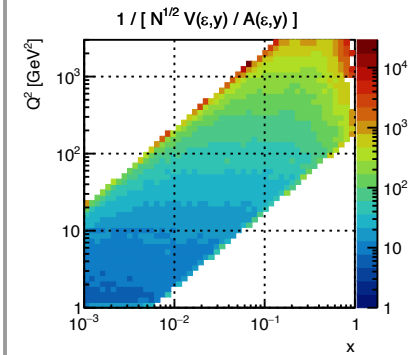
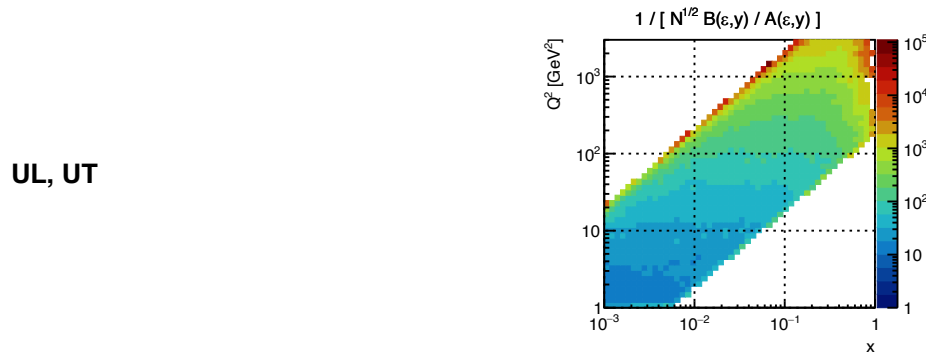
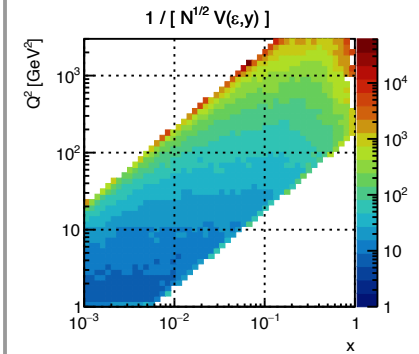
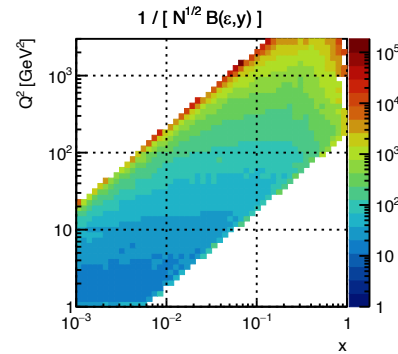
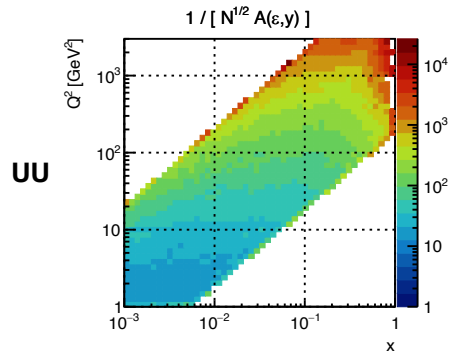


Access to TMDs: Kinematic factors

Twist 2

	Polarization	Depolarization
Boer-Mulders	UU	B
Sivers	UT	1
Transversity	UT	B/A
Kotzinian-Mulders	UL	B/A
Wormgear (LT)	LT	C/A
Helicity DiFF G_1^\perp	LU	C/A
	UL	1
<u>Twist 3</u>		
$e(x)$	LU	W/A
$h_L(x)$	UL	V/A
$g_T(x)$	LT	W/A

Statistical uncertainty scaling factor for $18x275$



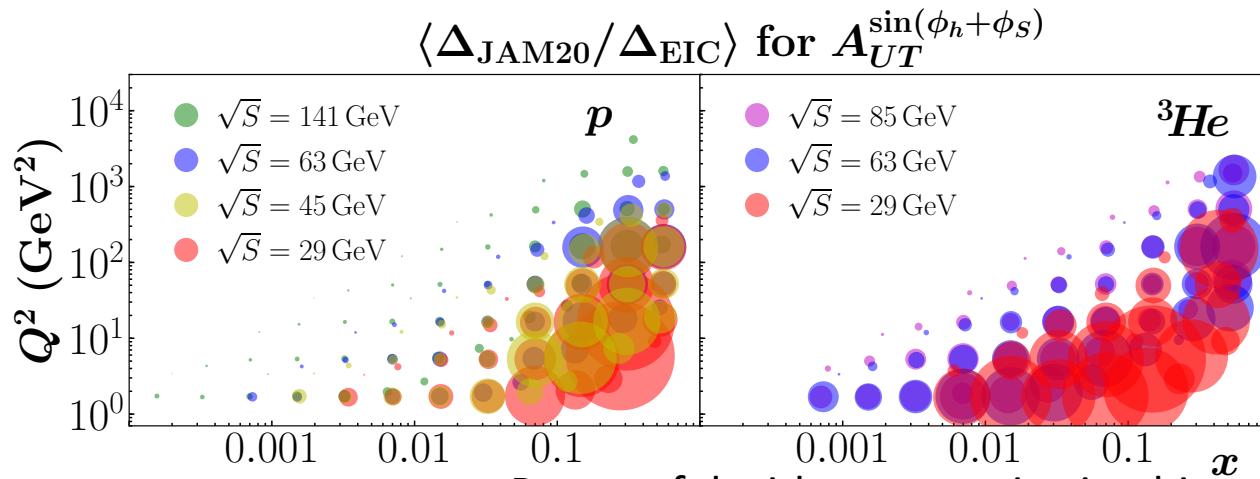
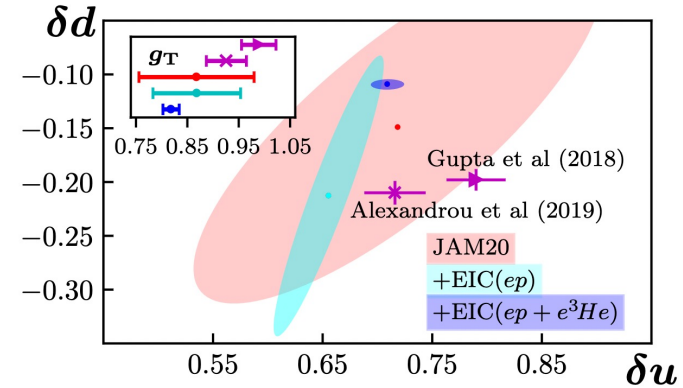
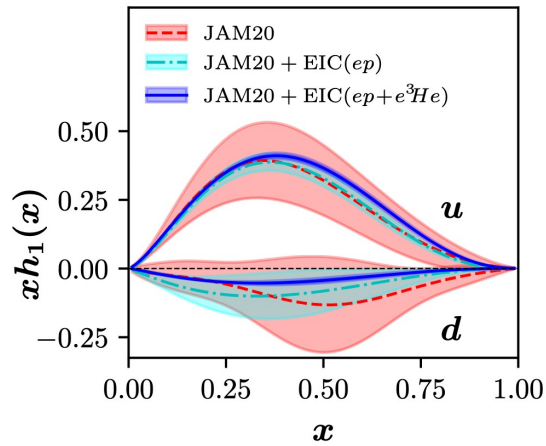
Depolarization Factors

Twist 2

	Polarization	Depolarization
Boer-Mulders	UU	B
Sivers	UT	1
Transversity	UT	B/A
Kotzinian-Mulders	UL	B/A
Wormgear (LT)	LT	C/A
Helicity DiFF G_1^\perp	LU	C/A
	UL	1
<u>Twist 3</u>		
$e(x)$	LU	W/A
$h_L(x)$	UL	V/A
$g_T(x)$	LT	W/A

Suppressed at EIC

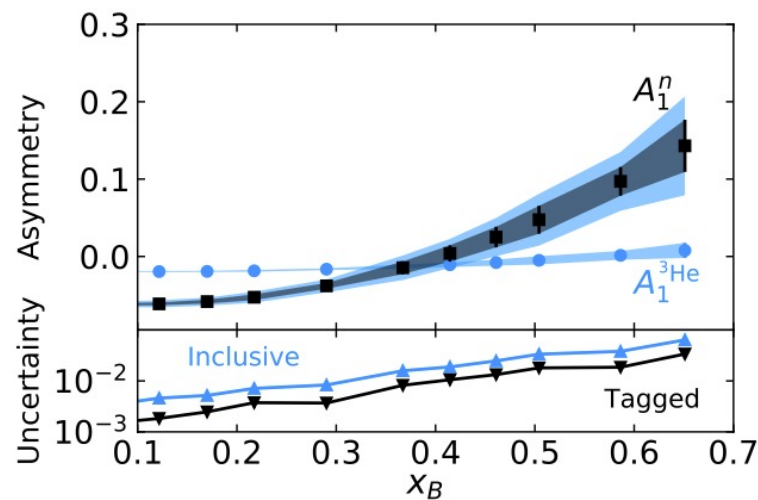
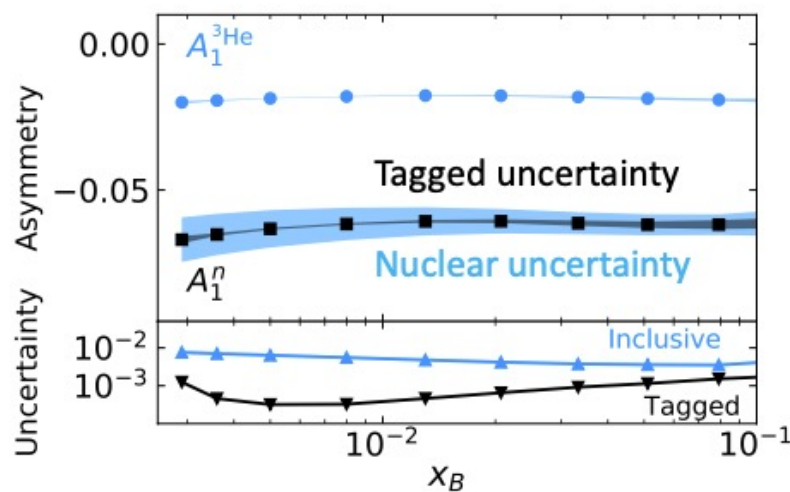
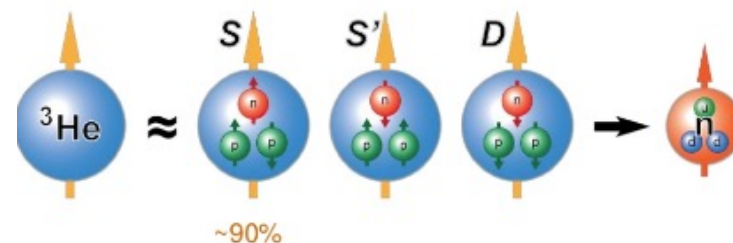
Example: transversity extraction from Jlab and the EIC



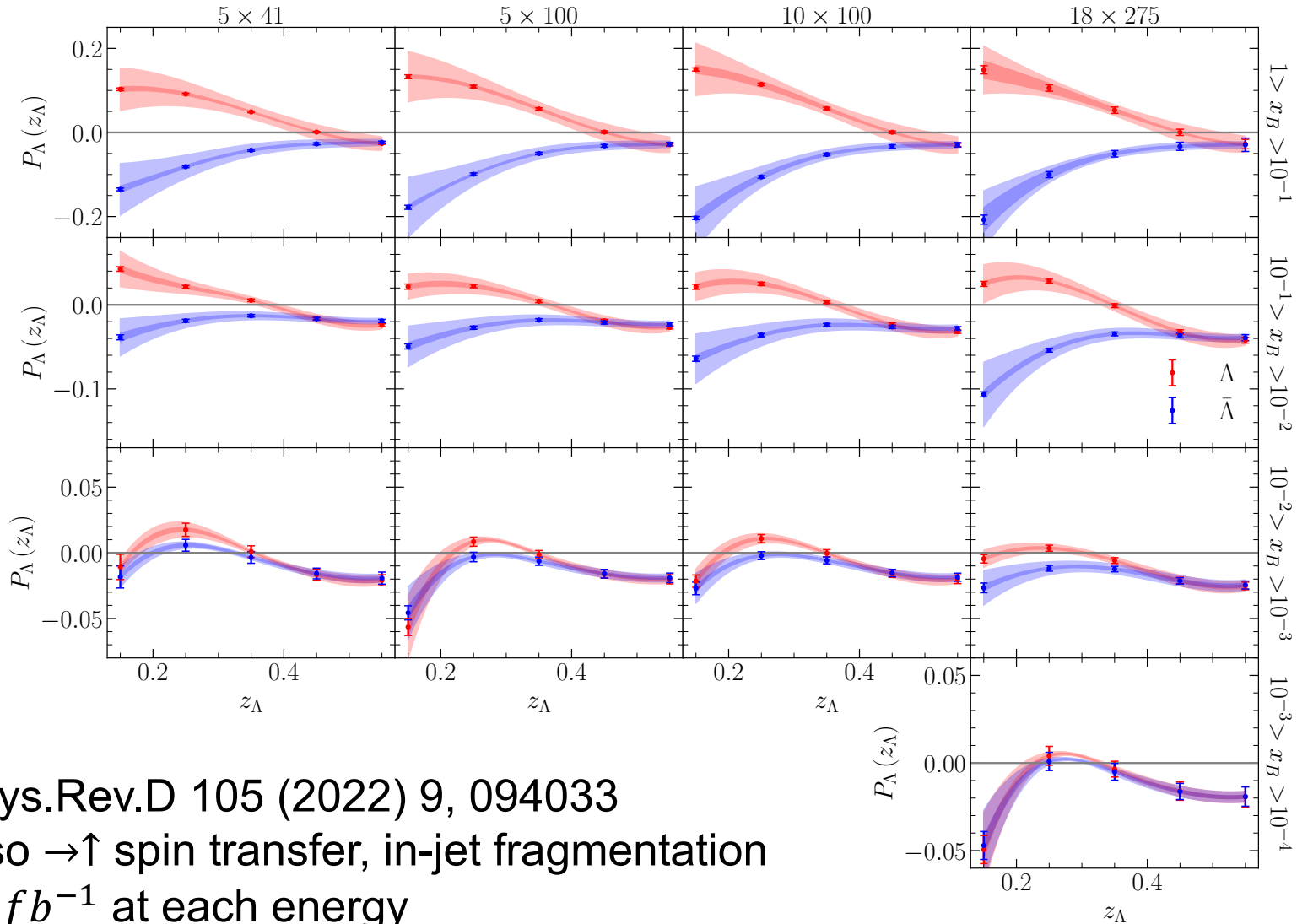
But careful with parametrization bias...

He^3 Double Tagging at the EIC allows clean neutron measurement

- Neutron is to 87% polarized
- Double tagged events thus provide access to polarized neutron beam
- Reconstruction of initial neutron momentum from tagged protons allows reduction of uncertainties from nuclear corrections



Precision Λ physics at the EIC



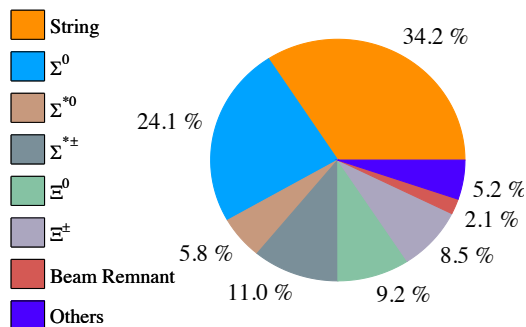
- Phys.Rev.D 105 (2022) 9, 094033
- Also $\rightarrow \uparrow$ spin transfer, in-jet fragmentation
- 40 fb^{-1} at each energy
- Significant impact of low \sqrt{s} data

Lambda feed-down composition vs JLab20

EIC

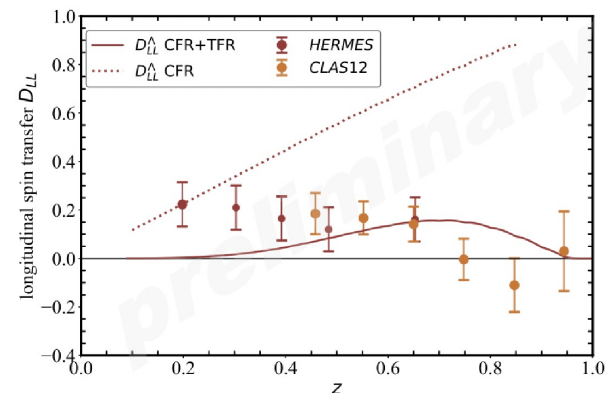
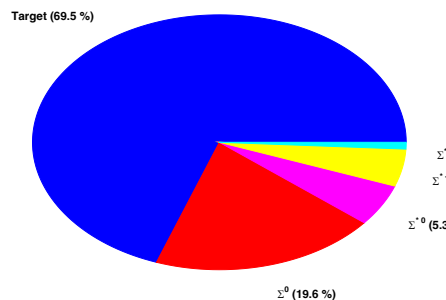
Λ ($0.1 < z < 0.3$)

18×275 GeV



JLAB12

Λ Parent 10.6GeV $x_F > 0$



Xiaoyan Zhao at S



Study by M. McEneaney(Duke)
JLab22 similar

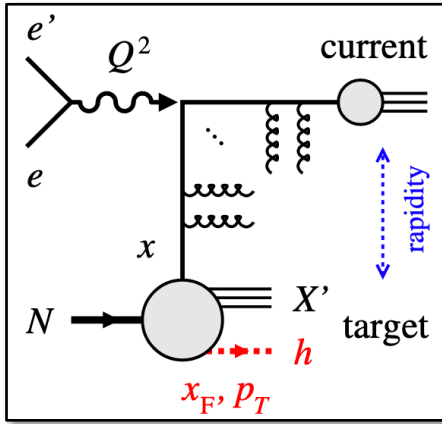
- Possible to unfold at the EIC (not so much at Jlab)
- ML methods might help

Summary/Conclusion/Outlook

- **Broad and diverse SIDIS program**
 - Present: JLab12, COMPASS
 - Future: JLab22, EIC
- **Wide kinematic reach enabling us to understand the full QCD picture**
- Jlab: unprecedented precision in valence quark regime
- Already exiting results, e.g. kaon final states
- EIC: new frontiers

- What is missing in this talk:
 - Target Fragmentation/Fracture Functions
 - Exclusive limit
 - TMDs in medium
 - Charged current at the EIC
 - Details on JLab12, 22 and EIC programs can be found in review papers and detector proposals
 - Yellow Report, Athena, ECCE, CORE proposals, PSQ report...
 - ...(and a lot more)

Fracture Functions to describe Target Region



- probability for the target (p/n) remnant to form a hadron *given* ejected quark q
No hard/soft energy scale separation
- Direct relationship to traditional PDFs by integrating over fractional longitudinal nucleon momentum ζ

$$\frac{d\sigma^{\text{TFR}}}{dx_B dy dz} = \sum_a e_a^2 (1 - x_B) M_a(x_B, (1 - x_B)z) \frac{d\hat{\sigma}}{dy}$$

$$\sum_h \int_0^{1-x} d\zeta \zeta \hat{u}_1(x, \zeta) = (1-x) f_1(x)$$

$$\sum_h \int_0^{1-x} d\zeta \zeta \hat{l}_{1L}(x, \zeta) = (1-x) g_{1L}(x)$$

$$\sum_h \int_0^{1-x} d\zeta \zeta M_a(x, \zeta) = (1-x) f_a(x)$$

M. Anselmino et al., Phys. Lett. B. 699 (2011), 108, [hep-ph] 1102.4214

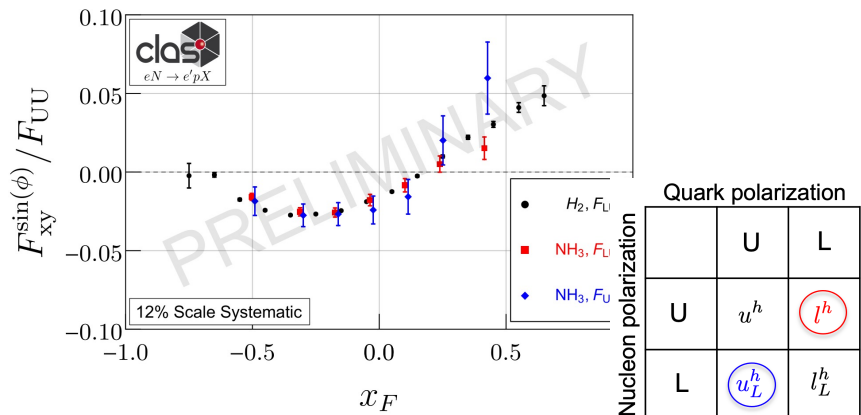
		Quark polarization					Quark polarization			
		U	L	T			U	L	T	
Nucleon polarization	U	f_1		h_1^\perp	Unpolarized PDF analog		\hat{u}_1	$\hat{l}_1^{\perp h}$	$\hat{t}_1^h, \hat{t}_1^\perp$	
	L		g_{1L}	h_{1L}^\perp		helicity PDF analog		$\hat{u}_{1L}^{\perp h}$	\hat{l}_{1L}^h	$\hat{t}_{1L}^h, \hat{t}_{1L}^\perp$
	T	f_{1T}^\perp	g_{1T}	h_{1T}, h_{1T}^\perp		etc. etc.		$\hat{u}_{1T}^h, \hat{u}_{1T}^\perp$	$\hat{l}_{1T}^h, \hat{l}_{1T}^\perp$	$\hat{t}_{1T}^h, \hat{t}_{1T}^{\perp h}$ $\hat{t}_{1T}^\perp, \hat{t}_{1T}^{\perp h}$

M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132

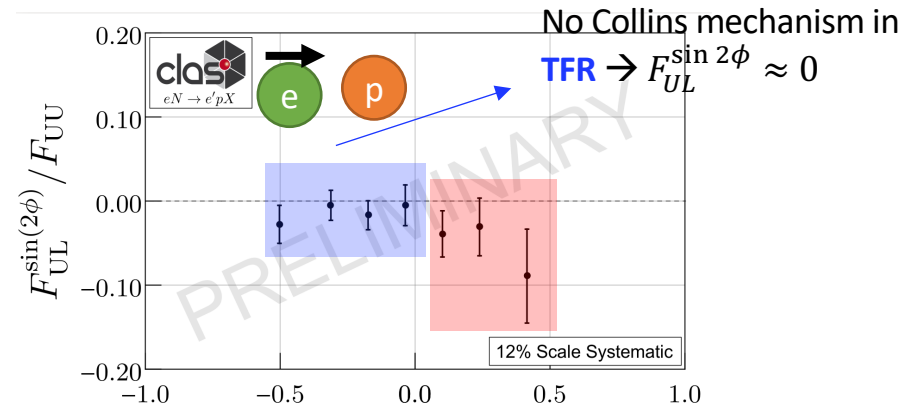


Preliminary Analysis: Fracture Functions

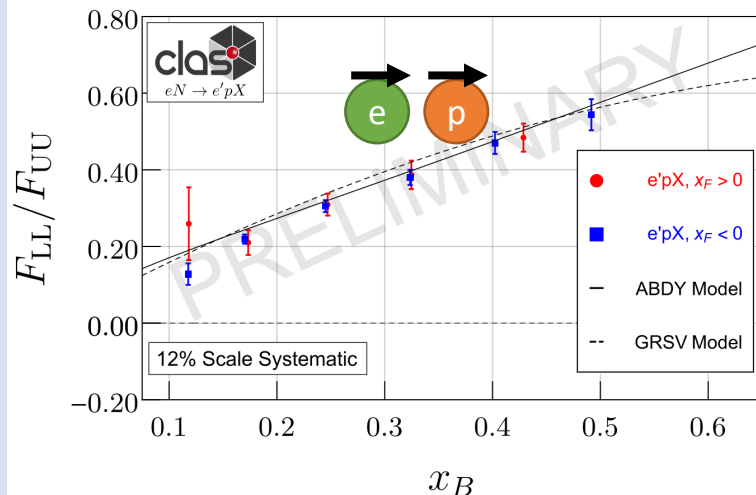
- First observation of correlations between Current and target region
- Visible separation between TFR ($x_F < 0$) and CFR ($x_F > 0$)



Twist-3 Collinear terms:
Chen, K. B., Ma, J. P. and Tong, X. B., [hep-ph] 2308.11251



$$F_{UL}^{\sin 2\phi_h} = C \left[-\frac{2(\hat{h} \cdot \mathbf{k}_T)(\hat{h} \cdot \mathbf{p}_T) - \mathbf{k}_T \cdot \mathbf{p}_T}{MM_h} h_{1L}^\perp H_1^\perp \right]$$



TFR Access to helicity distribution g_{1L}

$$A_{LL} = \lambda_\ell S_L \frac{\sqrt{1 - \epsilon^2} F_{LL}}{F_{UU,T}}$$

Integral relation holds!

Quark polarization

	U	L
U	\hat{u}_1	\hat{l}_1^h
L	\hat{u}_{1L}^\perp	\hat{l}_{1L}^\perp

Nucleon polarization

M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132

SIDIS Datasets, present and future

CLAS12 in Hall B

Run Group A (Unpolarized LH₂ target)

- ★ unpolarized SIDIS cross section off proton
- ★ A_{LU} in Beam Spin Asymmetries

Run Group B (Unpolarized LD₂ target)

- ★ Complementary to RG-A
→ allow for u/d quark flavor separation

Run Group C (longitudinally polarized NH₃ and ND₃)

- ★ F_{UL} and F_{LL}

Run Group K (Unpolarized LH₂ target)

- ★ - 6.5, 7.5, 8.4 GeV e- beam
- ★ $F_{UU,L}$, $F_{UU,T}$ Separation

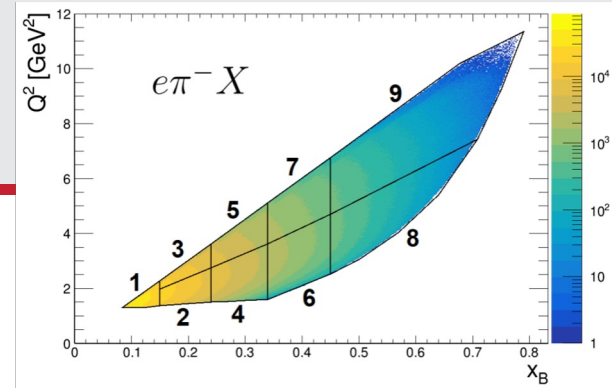
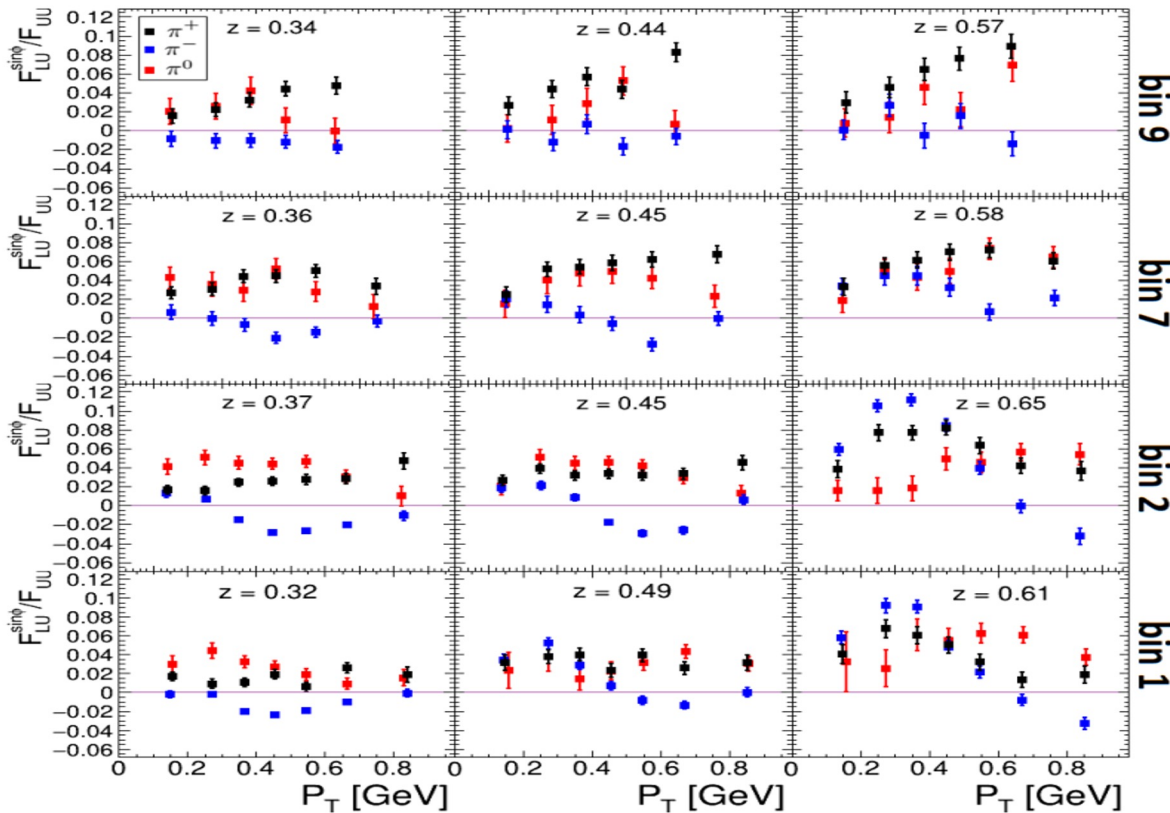
Run Group H (transversely polarized NH₃)

- ★ F_{UT} structure function

Hall A

- ≈ 2028 SoLID with He3/proton target (long/transverse)

⇒ should be
completed before 1st Shutdown ≈ 2032



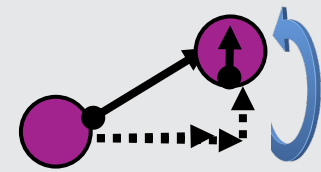
S. Diehl

- ★ If Collins term only (H_1^\perp) \rightarrow hierarchy of the A_{LU} 's
 $A_{LU}(\pi^-) < A_{LU}(\pi^0) = 0 < A_{LU}(\pi^+)$

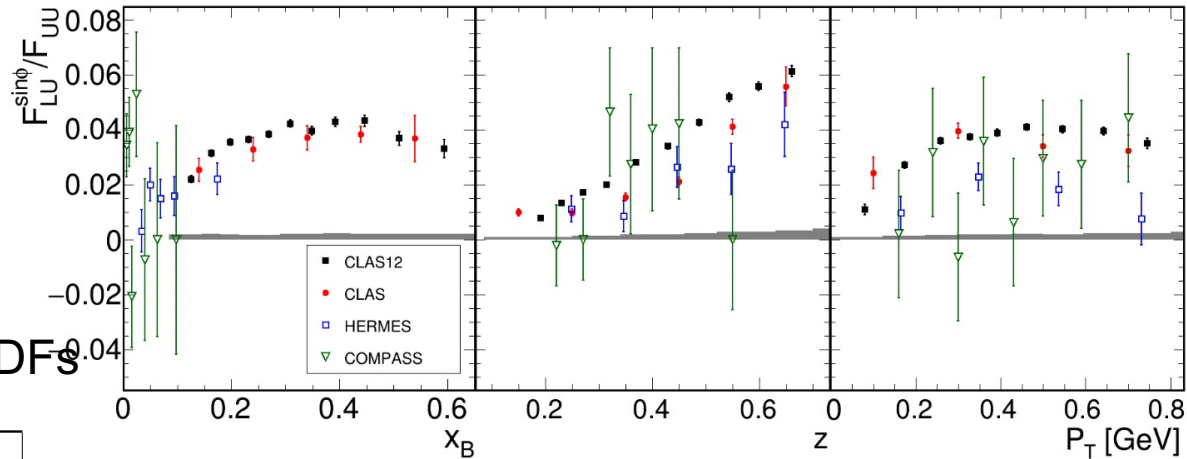
- ★ Observed is more **Sivers-like (g^\perp)**, asymmetry comes from struck u-quark

$$A_{LU}(\pi^-) < A_{LU}(\pi^0) = A_{LU}(\pi^+)$$

CLAS12 pion BSAs



Phys.Rev.Lett. 128 (2022) 6, 062005



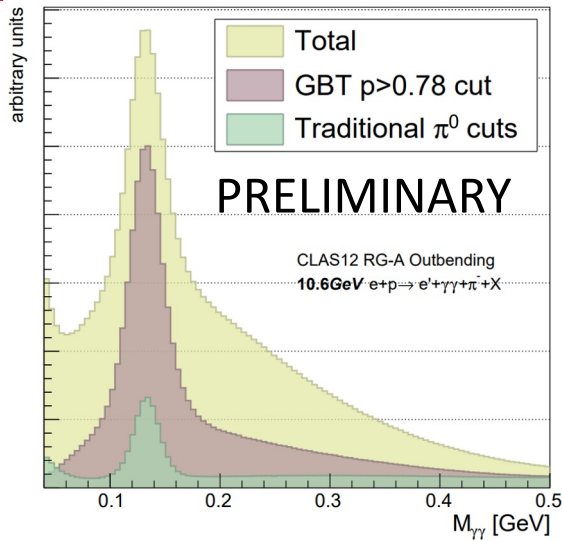
Higher Twist PDFs

N/q	U	L	T
U	f^\perp	g^\perp	h, e
L	f_L^\perp	g_L^\perp	h_L, e_L
T	f_T, f_T^\perp	g_T, g_T^\perp	$h_T, e_T, h_T^\perp, e_T^\perp$

S. Diehl

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q} \mathcal{C} \left[-\frac{\hat{h} \cdot k_T}{M_h} \left(x \overset{\text{twist-3 pdf}}{\underbrace{e}_{\text{Collins FF}}} \overset{\text{unpolarized PDF}}{\underbrace{H_1^\perp}} + \frac{M_h}{M} \overset{\text{twist-3 t-odd PDF}}{\underbrace{f_1}_{\text{twist-3 FF}}} \frac{\overset{\text{Boer-Mulders}}{\underbrace{\tilde{G}^\perp}}}{z} \right) + \frac{\hat{h} \cdot p_T}{M} \left(x \overset{\text{twist-3 t-odd PDF}}{\underbrace{g^\perp}_{\text{unpolarized FF}}} \overset{\text{Boer-Mulders}}{\underbrace{D_1}} + \frac{M_h}{M} \overset{\text{twist-3 t-odd PDF}}{\underbrace{h_1^\perp}_{\text{Boer-Mulders}}} \frac{\overset{\text{Boer-Mulders}}{\underbrace{\tilde{E}}}}{z} \right) \right]$$

Dihadron Production $ep \rightarrow e\pi^\pm\pi^0(X)$



★ Nearest-neighbor GBDT model to reduce γ background

★ Negative $\sin(\phi_R)$ asymmetry for $\pi-\pi^0 \rightarrow e(x)$ extraction

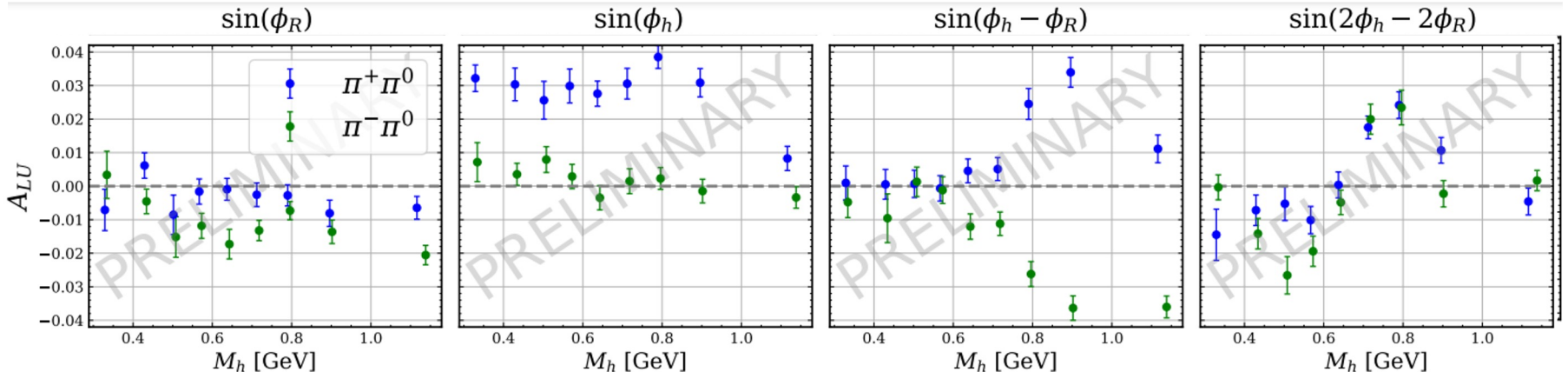
★ Strong positive $\sin(\phi_h)$ asymmetry for $\pi+\pi^0 \rightarrow u$ quark dominated channels (seen in 1h SIDIS frequently)

★ Isospin symmetries of G_1 DiFF observed in $\sin(\phi_h - \phi_R)$

★ Strong enhancement near resonant region

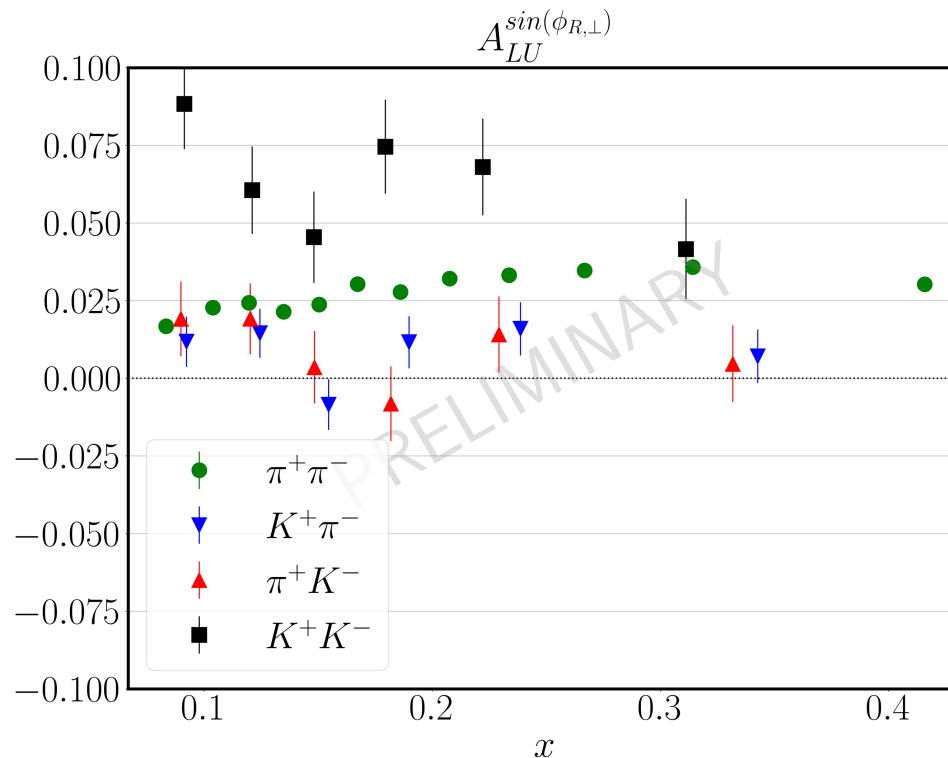
$$e \otimes H_1^\perp |l, m\rangle$$

$$f_1 \otimes G_1^\perp |l, m\rangle$$



Slide by G. Matousek

Kaons



- Kaon \gg Pions

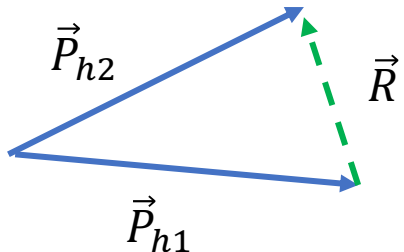
- Assuming u –quark dominance \rightarrow FF effect?
- Twist3 FF relevant?
- Or $e(x)$ for strange quarks

Better: Dihadron Fragmentation Functions

Additional Observable:

$$\vec{R} = \vec{P}_1 - \vec{P}_2 :$$

The relative momentum of the hadron pair is an additional degree of freedom:

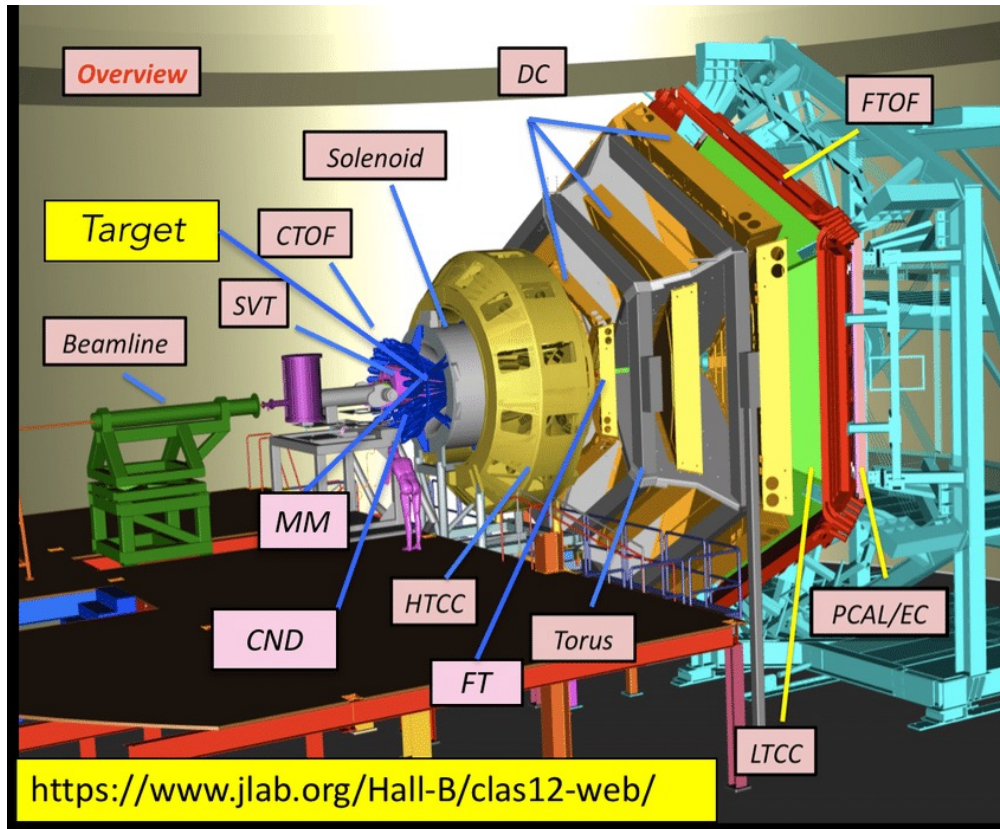


More degrees of freedom → More information about correlations in final state

Additional FFs that do not exist in single-hadron case $G_1^\perp \rightarrow$ related to jet handedness

Parton polarization → Hadron Polarization ↓	Spin averaged	longitudinal	transverse
spin averaged	$D_1^{h/q}(z, M)$ 		$H_1^{\perp h/q}(z, p_T, M, (\mathbf{P}_h), \theta)$ 'Di-hadron Collins'
longitudinal			
Transverse		$G_1^\perp(z, M, P_h, \theta) =$ T-odd, chiral-even → jet handedness QCD vacuum structure 	$H_1^{\perp}(z, M, (\mathbf{P}_h), \theta) =$ T-odd, chiral-odd Collinear

Jefferson Lab



Forward Detector

- Torus Magnet
- Drift Chamber (DC)
- Forward Time of Flight (FTOF)
- High-threshold Cherenkov Counter (HTCC)
- Low-threshold Cherenkov Counter (LTCC)
- Ring Imaging Cherenkov Detector (RICH)
- Preshower + Electromagnetic Calorimeter (PCAL/EC)
- Forward Tagger (FT)

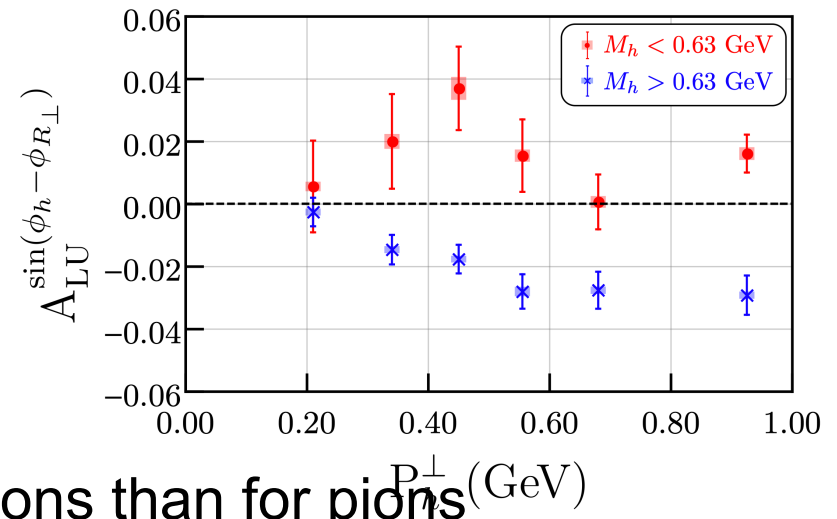
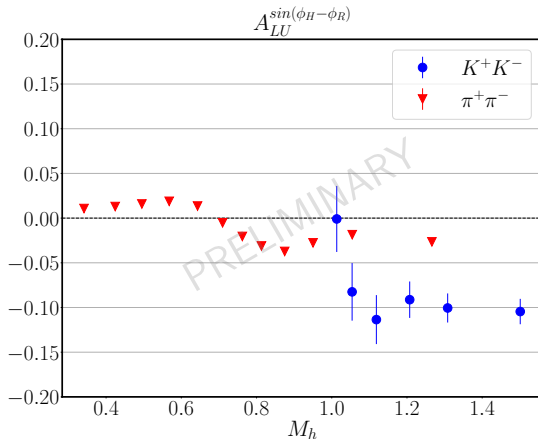
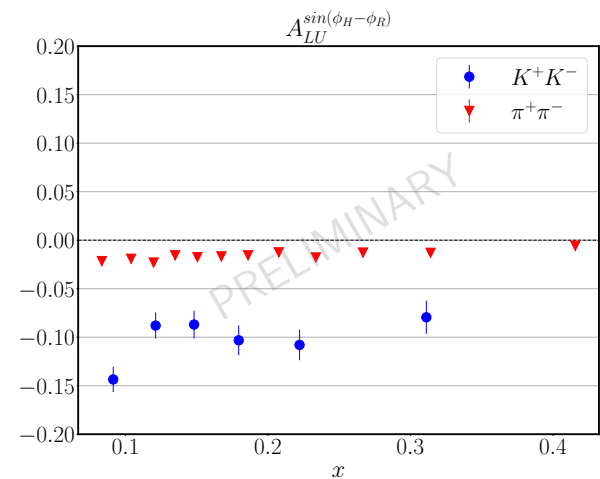
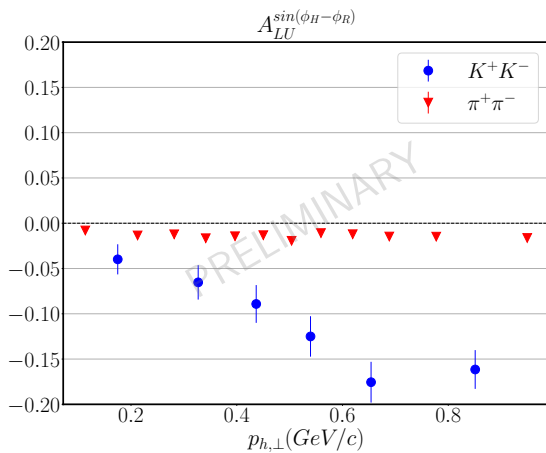
■ Longitudinally Polarized Electron Beam

- $E = 10.6 \text{ GeV}$
- $P = 86\text{--}89\%$

■ Unpolarized Liquid H_2 Fixed Target

- Torus magnet \rightarrow electrons inbending

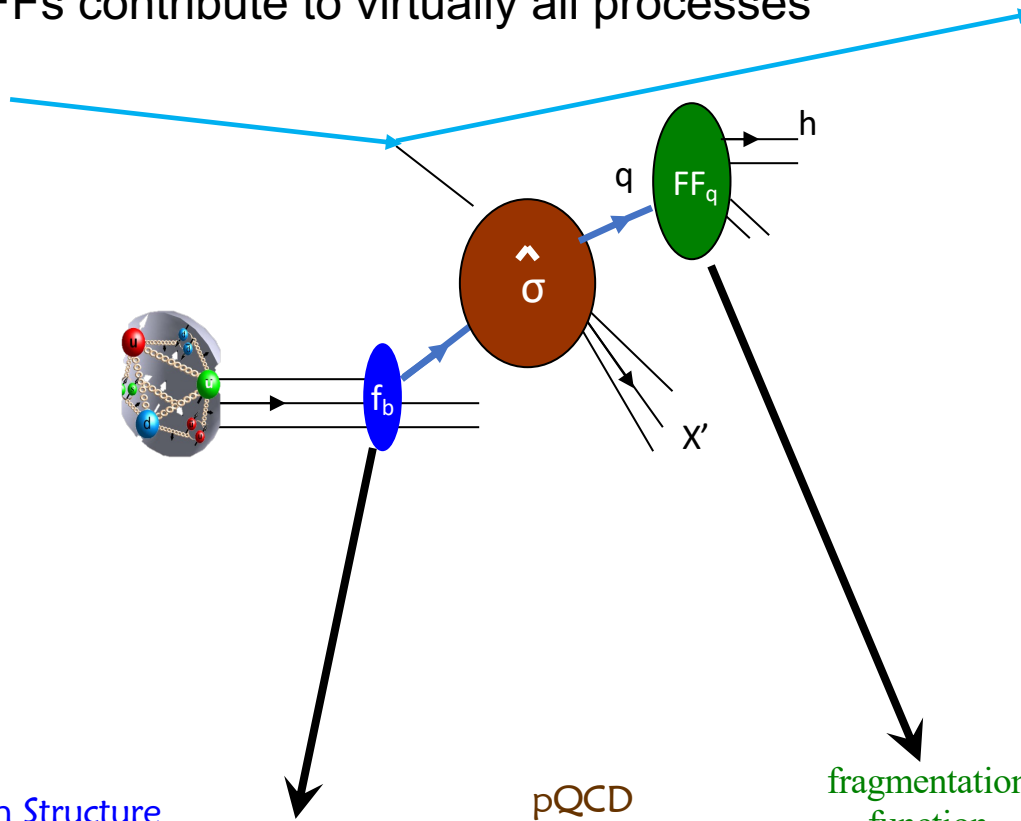
Asymmetries sensitive to G_1^\perp



- sp –interference term larger for kaons than for pions
- Not true for all interference terms (not shown)
- $M_{KK} > m_\phi$ can account for p_\perp dependence

Hard Scattering is a premier tool to probe the quark and gluon degrees of freedom

- Proton Structure extracted using QCD factorization theorem
- FFs contribute to virtually all processes



Wigner distributions $\rho(x, \vec{b}_\perp, \vec{k}_\perp)$
 5 dimensional!!

Transverse momentum \vec{k}_\perp
 Transverse position \vec{b}_\perp
 partons
 Transverse plane
 Longitudinal momentum $k^+ = xP^+$

- Q^2 “scale” of the probe, “resolution”
- x : bjorken **scaling variable**, in partonic picture momentum fraction of the struck parton, “shutter speed”
- z fractional quark momentum carried by hadron
- p_T








Proton Structure

pQCD

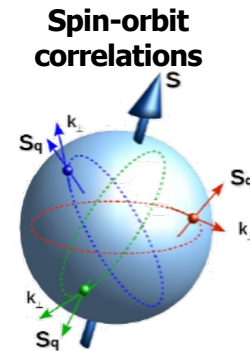
fragmentation function

$$\frac{d^2 \sigma(ep \rightarrow \pi X)}{dx dz p_{hT}} \propto q(x, k_T) \times \frac{d\sigma^2(eq \rightarrow e'q')}{dx} \times FF(z, p_T) \delta(p_{hT} - (k_T + p_T))$$

Momentum structure in the parton model parametrized by TMDs (spin $\frac{1}{2}$)

N \ q	U	L	T
U			
L			
T			

- In addition to the spin-spin correlations can have spin momentum correlations!



Run Group SIDIS programs at a glance

Run Group A (Unpolarized LH₂ target - 10.6 GeV e⁻ beam)

- ★ Measurements of **unpolarized SIDIS cross section** off proton (*ex: π multiplicities*)
- ★ Access to **higher-twist PDFs** through A_{LU} beam-spin asymmetries (BSAs)
- ★ Study impact of struck quark's spin/flipavor/momentum on **hadronization**
 - Separate contributions from vector meson decays (*ex: direct π vs. decay π*)
- ★ Observe correlations between struck quark and target breakup

Run Group B (Unpolarized LD₂ target - 10.6 GeV e⁻ beam)

- ★ Complementary to RG-A → allow for ***u/d* quark flavor separation** of observables

Run Group C (Dynamic **longitudinally** polarized NH₃ and ND₃ - 10.6 GeV e⁻ beam)

- ★ Access to **F_{UL}** and **F_{LL}** structure functions → Sensitive to different PDFs and FFs
 - Dihadron SIDIS will give first measurements of **higher-twist** fragmentation functions

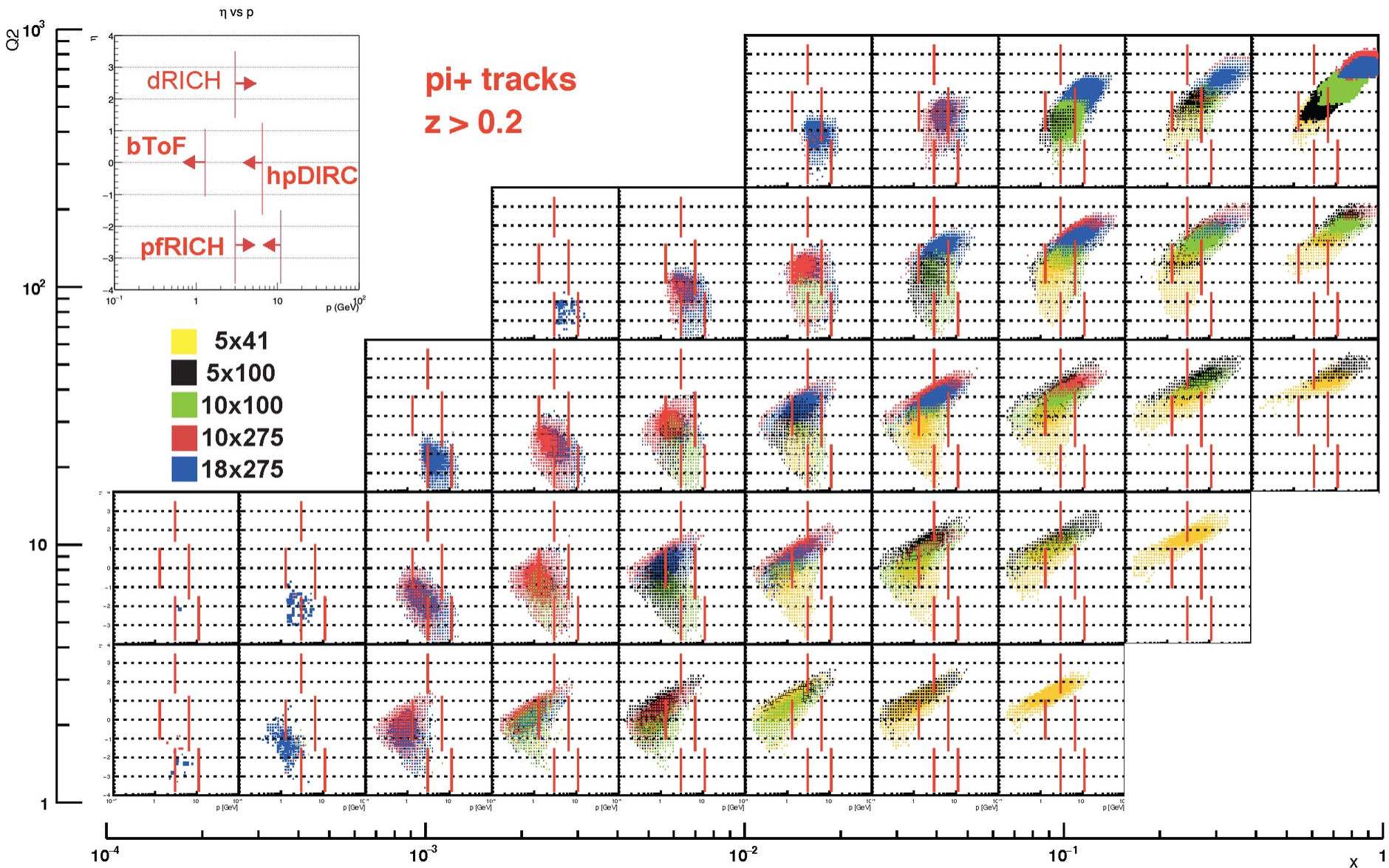
Run Group K (Unpolarized LH₂ target - 6.5, 7.5, 8.4 GeV e⁻ beam)

- ★ Separation of longitudinal (F_{UU,L}) and transverse (F_{UU,T}) photons from SIDIS cross section

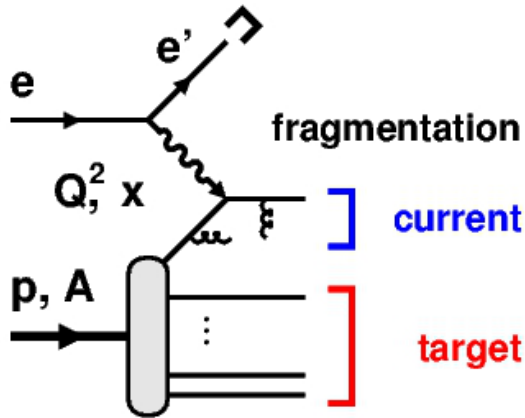
Run Group H (Dynamic **transversely** polarized NH₃ - 10.6 GeV e⁻ beam)

- ★ Access to **F_{UT}** structure function → Transverse spin structure at high (essentially unmeasured) x

Wide Coverage



Target Fragmentation → Fracture Functions



- Correlations between CFR-TFR allows to extract fracture functions coupled to FFs

Quark polarization

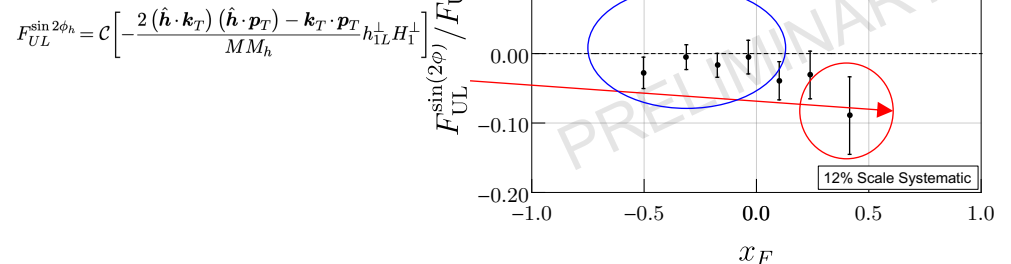
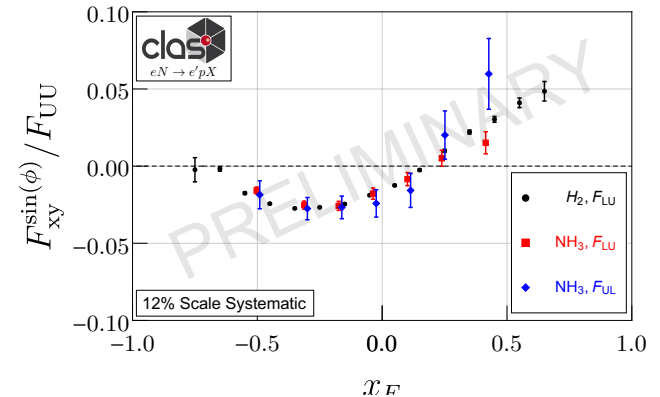
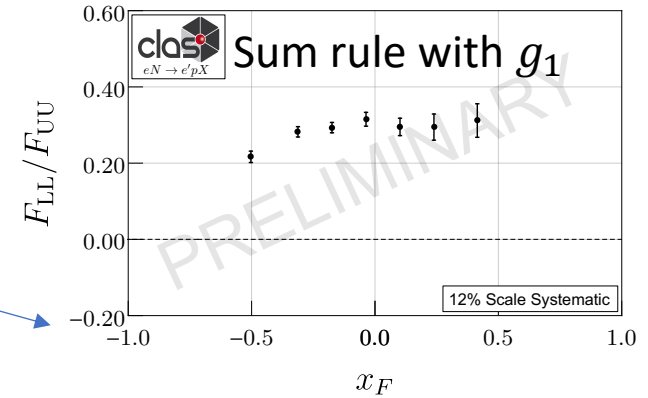
	U	L	T
U	\hat{u}_1	$\hat{l}_1^{\perp h}$	$\hat{t}_1^h, \hat{t}_1^\perp$
L	$\hat{u}_{1L}^{\perp h}$	\hat{l}_{1L}	$\hat{t}_{1L}^h, \hat{t}_{1L}^\perp$
T	$\hat{u}_{1T}^h, \hat{u}_{1T}^\perp$	$\hat{l}_{1T}^h, \hat{l}_{1T}^\perp$	$\hat{t}_{1T}^h, \hat{t}_{1T}^{\perp h}$ $\hat{t}_{1T}^\perp, \hat{t}_{1T}^{\perp h}$

M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph/1109.1132]

Quark polarization

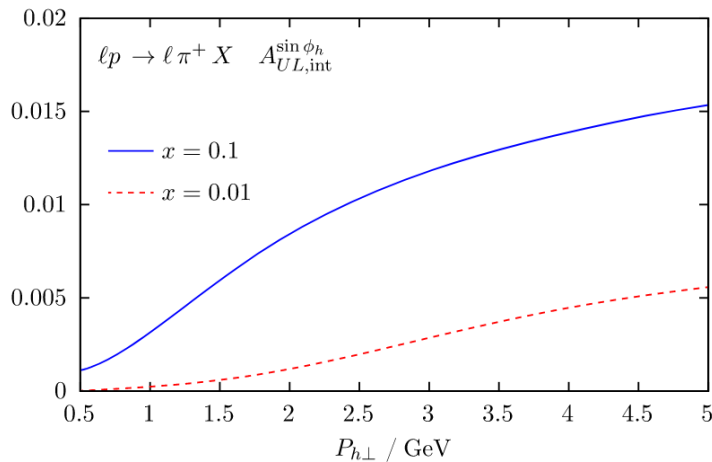
	U	L
U	u^h	l^h
L	u_L^h	l_L^h

Twist-3 Collinear terms:
Chen, K. B., Ma, J. P. and Tong, X. B., [hep-ph/2308.11251]

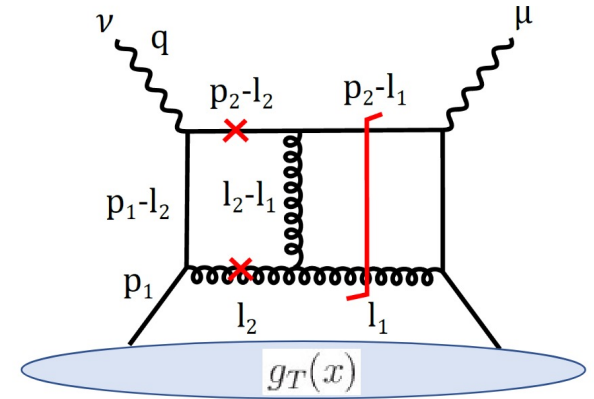


Appearance of perturbative asymmetries

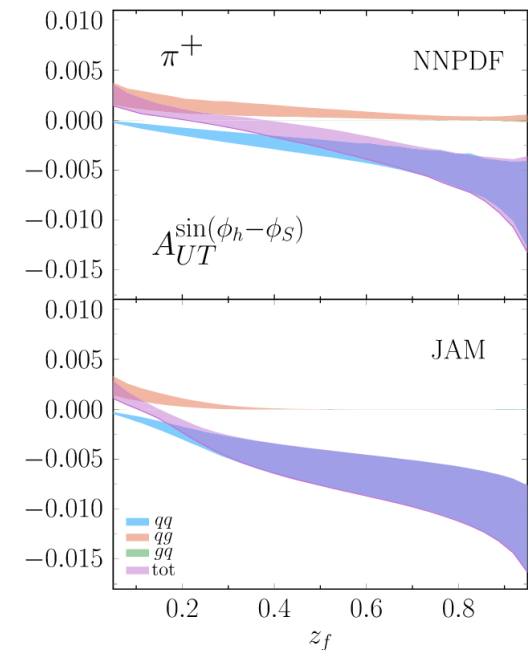
- Additionally: perturbative generation of asymmetries from g_T
- What about di-hadron correlations?



Abele, Aicher, Piacenza, Schafer, Vogelsang (2022)



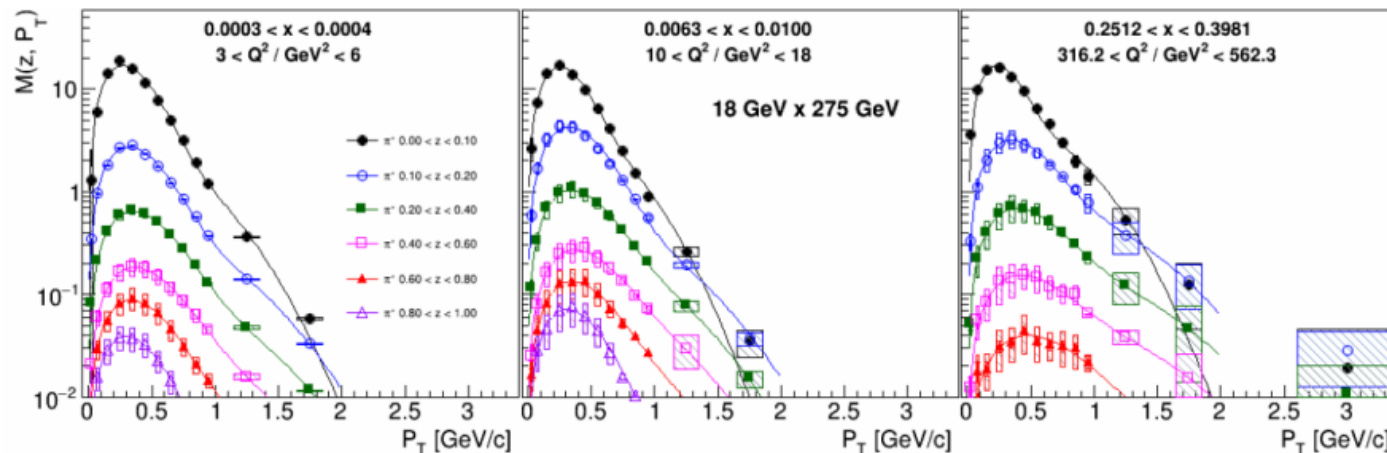
EIC, $\sqrt{S_{ep}} = 45 \text{ GeV}$



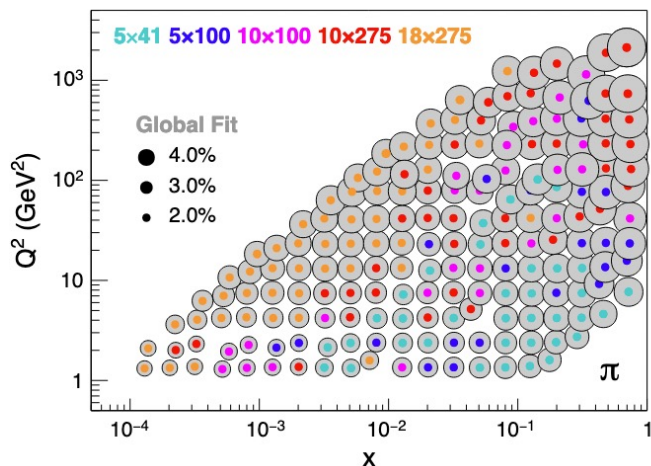
Benic et al. Phys. Rev. D 104, 094027

Unpolarized TMDs

- Top: Explicit z dependence of select pion multiplicities in 3 x - Q^2 bins, including the double-Gaussian fits

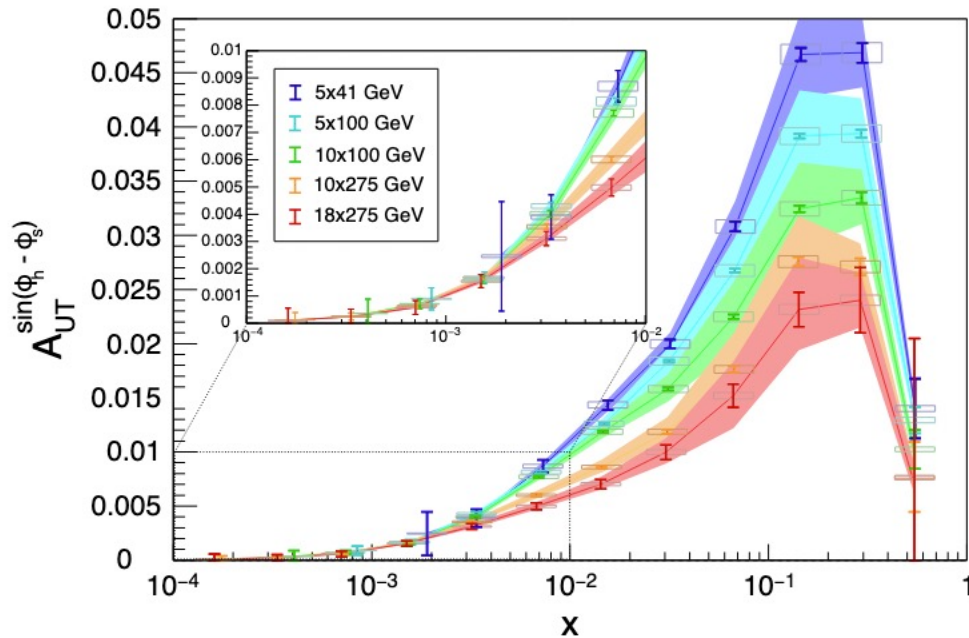


[hep-ex:2207.10893](https://arxiv.org/abs/hep-ex/2207.10893)



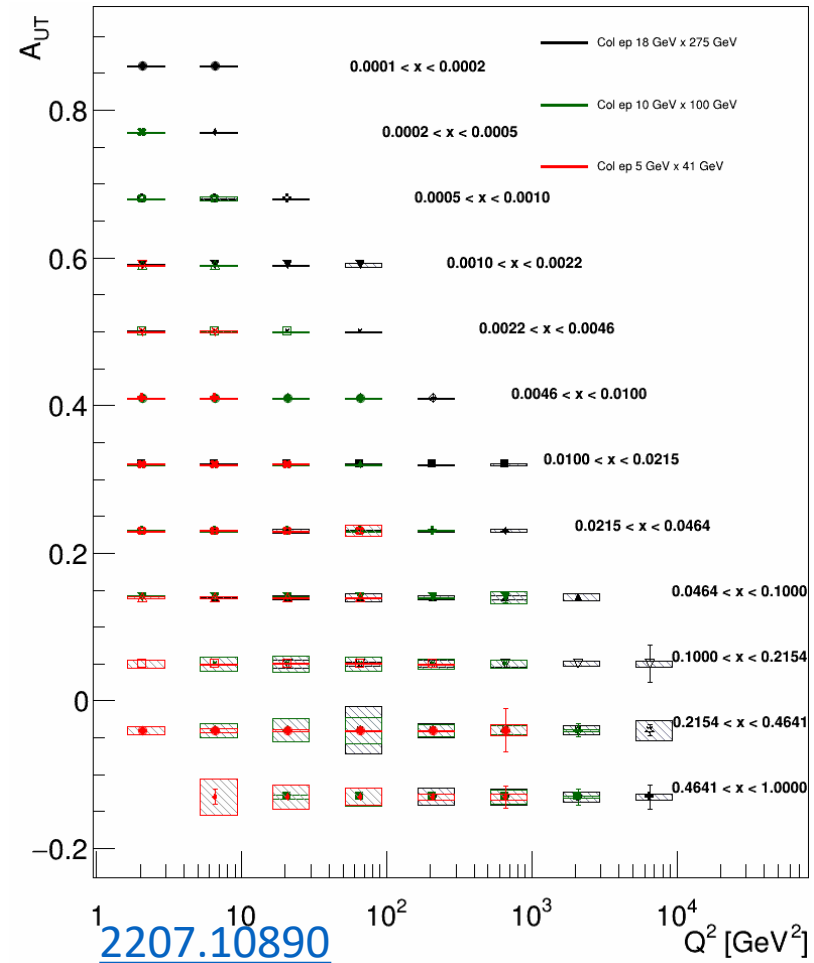
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$$0.2 < z < 0.7, Q^2 > 1.0, y > 0.05, \frac{qT}{Q} < 1.0$$



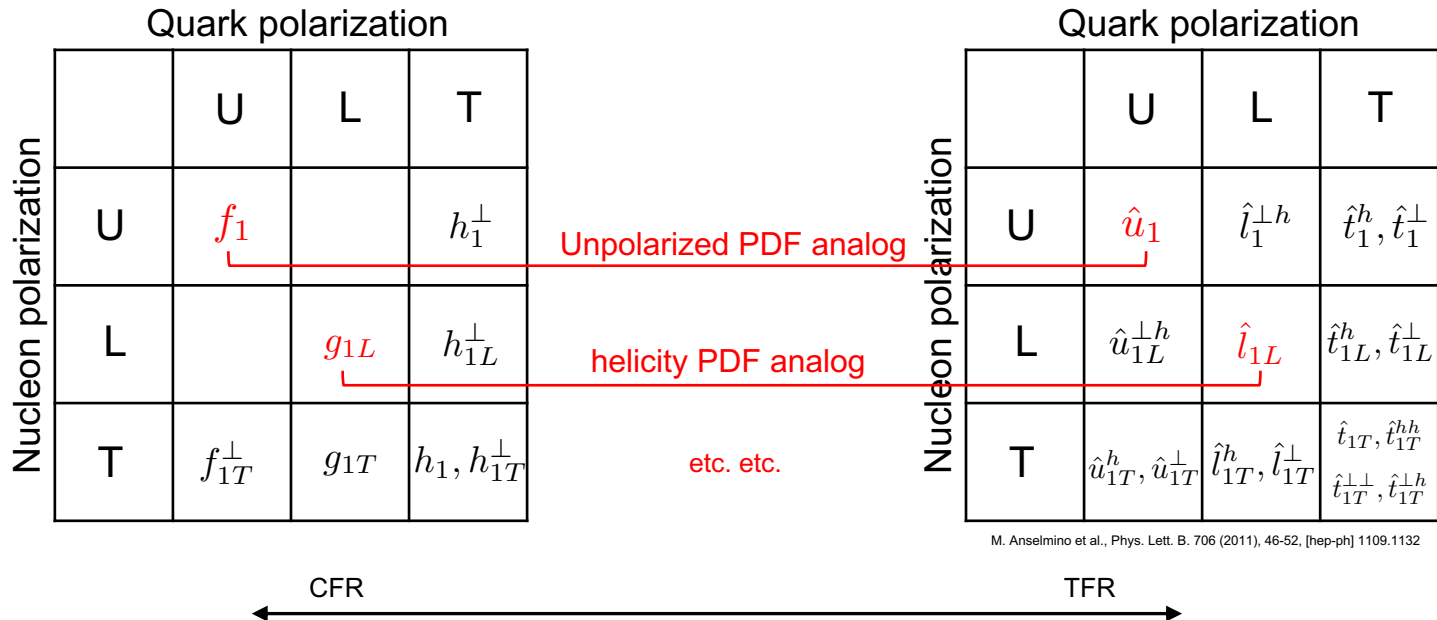
Projected uncertainties on Collins Asymmetries by ECCE

Analog to PDFs; Momentum Sum Rules

- A direct relationship exists to the eight leading twist PDFs after the fracture functions are integrated over the fractional longitudinal nucleon momentum, ζ .

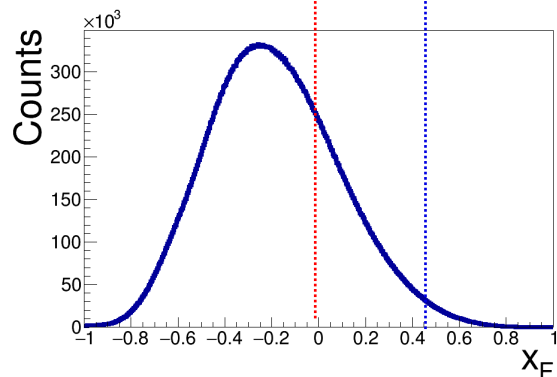
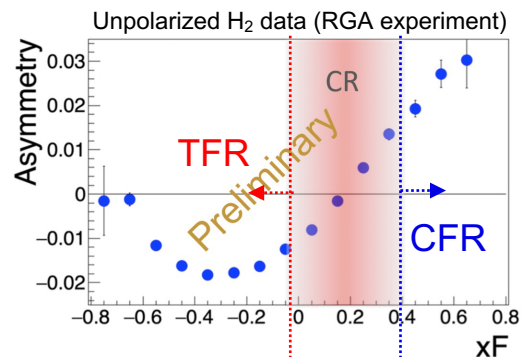
$$\sum_h \int_0^{1-x} d\zeta \zeta M_a(x, \zeta) = (1-x) f_a(x)$$

M. Anselmino et al., Phys. Lett. B. 699 (2011), 108, [hep-ph] 1102.4214



M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132

Can We Separate Target and Current?



Feynman variable

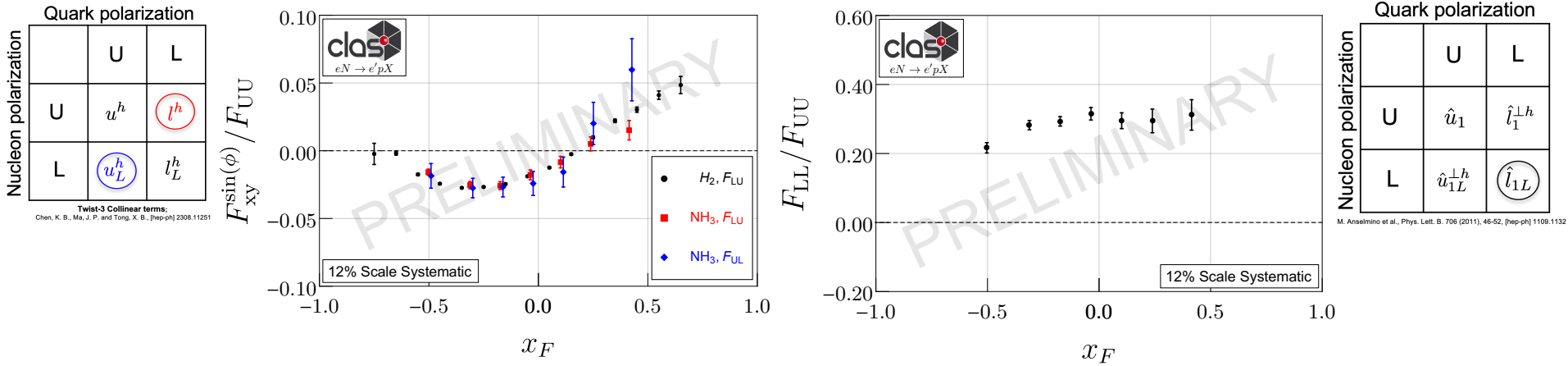
$$x_F = \frac{p_h^z}{p_h^z(\text{max})} \quad \text{in CM frame } \mathbf{p} = -\mathbf{q}, \quad -1 < x_F < 1$$

Rapidity

$$y_h = \frac{1}{2} \log \frac{p_h^+}{p_h^-} = \frac{1}{2} \log \frac{E_h + p_h^z}{E_h - p_h^z}$$

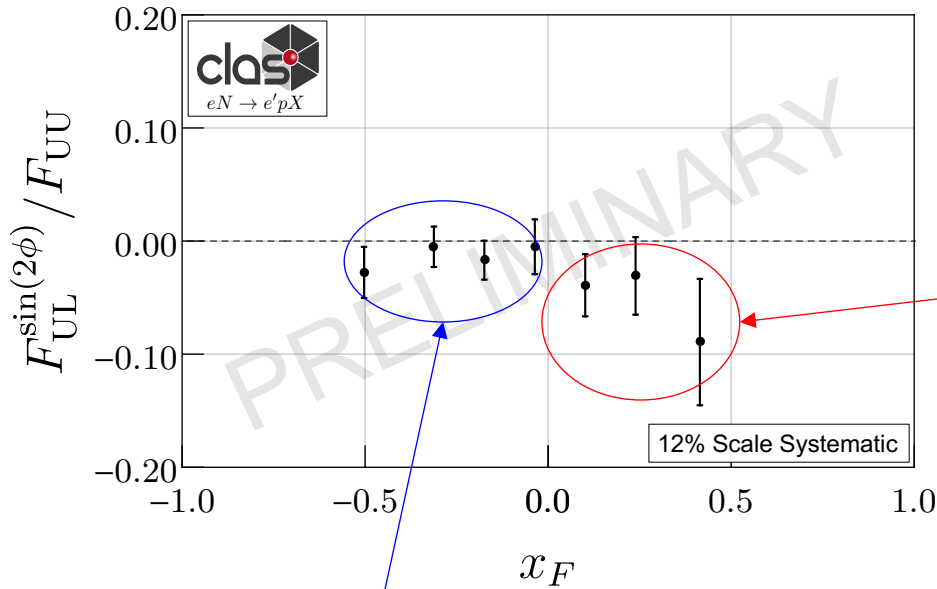
- No clear *experimental* definition of what constitutes current production versus target production.
- Odd structure functions, with different production mechanisms in both regions, give a possible clue.
- Protons (as opposed to mesons) at CLAS12 kinematics give a unique opportunity because they have extensive coverage in both regions.

Current and Target Separation



- Odd-function (sine) modulations exhibit a sign flip around the transition from target to current fragmentation. Interestingly, we observe $F_{LU} \sim F_{UL}$.
- Even-function (cosine) behavior of double-spin asymmetry does not show a sign flip; possible signs decreasing F_{LL} as $x_F \rightarrow \pm 1$ (x_B decreasing but likely not the only cause).
- Consistent beam-spin asymmetries in unpolarized H_2 and polarized NH_3 indicates minimal nuclear medium modification.

Kotzinian-Mulders Asymmetry



$$F_{UL}^{\sin 2\phi_h} = C \left[-\frac{2 (\hat{\mathbf{h}} \cdot \mathbf{k}_T) (\hat{\mathbf{h}} \cdot \mathbf{p}_T) - \mathbf{k}_T \cdot \mathbf{p}_T}{MM_h} h_{1L}^\perp H_1^\perp \right]$$

- The $F_{UL}^{\sin 2\phi}$ asymmetry is purely generated by the **Collins mechanism** – whereby a transversely polarized quark flips orientation during hadronization and produces an asymmetric distribution in the transverse plane.
- Hadronization in the TFR is more isotropic – there is no additional chiral-odd quantity like the **Collins function** to pair with the **Kotzinian-Mulders** TMD because factorization into separate soft and hard scale processes does not hold.

- **No Collins mechanism** in the TFR so $F_{UL}^{\sin 2\phi}$ (and $F_{UU}^{\cos 2\phi}$) are pure twist-4. **We would expect small magnitude at $-x_F$.**

Early signs give a *possible* hint but need more statistics!

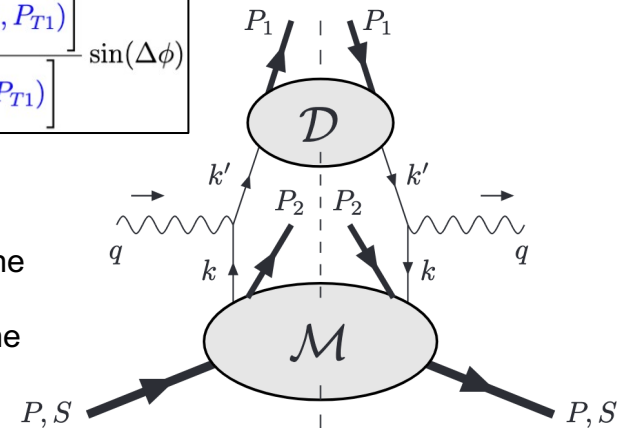
Back-to-back (dSIDIS) Formalism

- When two hadrons are produced “back-to-back”^{1,2} with one in the CFR and one in the TFR the structure function contains a convolution of a **fracture function** and a **fragmentation function**.
- Leading twist beam(target)-spin asymmetry.

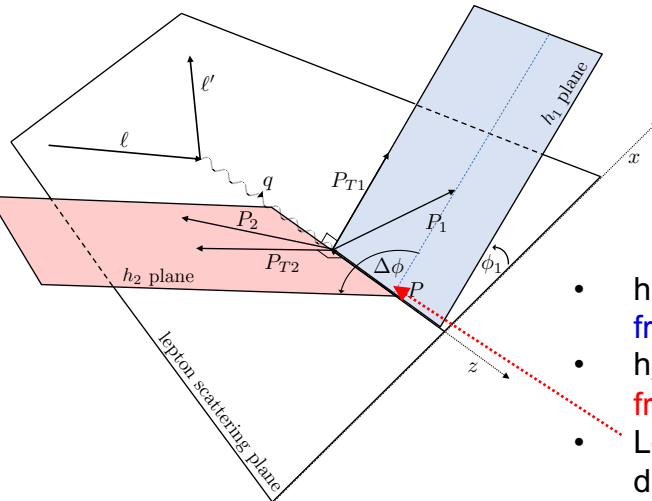
$\hat{l}_1^{\perp h}$ Unique access to longitudinally polarized quarks in unpolarized nucleon... no corresponding PDF!

$$A_{LU} = -k(\epsilon) \frac{P_{T1} P_{T2}}{m_1 m_2} \frac{C \left[w_5 \hat{l}_1^{\perp h}(x, \zeta_2, P_{T2}) D_1(z_1, P_{T1}) \right]}{C \left[\hat{u}_1(x, \zeta_2, P_{T2}) D_1(z_1, P_{T1}) \right]} \sin(\Delta\phi)$$

- h_1 in the CFR with production dictated by the **fragmentation function**
- h_2 in the TFR with production dictated by the **fracture function**
- Long range correlation depends on the difference in azimuthal angles of both hadrons



Handbag diagram for dihadron production; lower blob contains FrFs and upper blob contains the FFs.



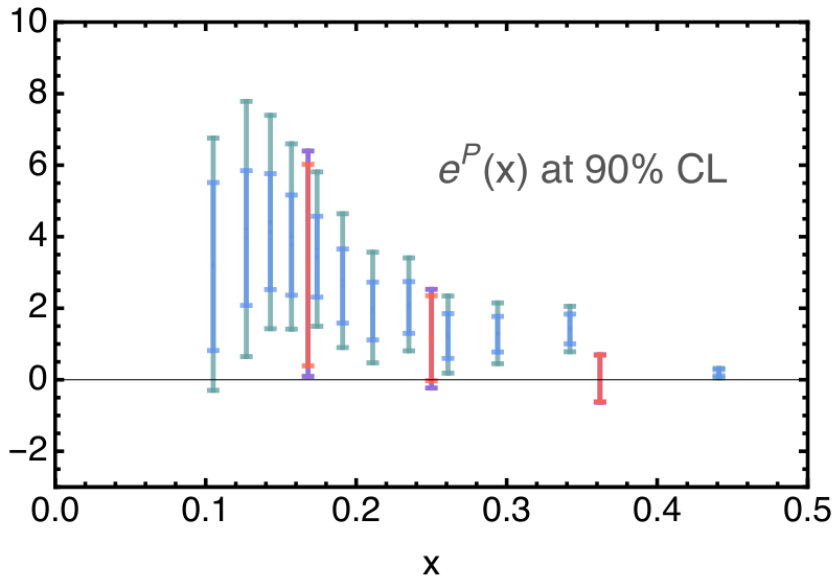
Kinematic plane for b2b dihadron production.

1. M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132
 2. M. Anselmino et al., Phys. Lett. B. 713 (2012), 317-320, [hep-ph] 1112.2604

New $e(x)$ Extraction – Proton Flavor Combination

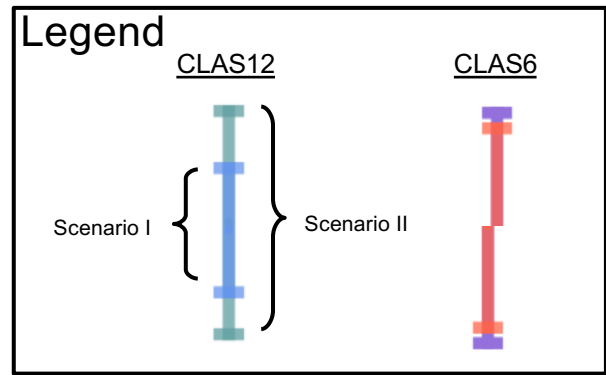
$$A_{LU}^{\sin \phi_R} \propto \frac{M \sum_q e_q^2 \left[x e^q(x) H_{1,sp}^{\Delta,q}(z, m_{\pi\pi}) + \frac{m_{\pi\pi}}{zM} f_1^q(x) \tilde{G}_{sp}^{\Delta,q}(z, m_{\pi\pi}) \right]}{Q \sum_q e_q^2 f_1^q(x) D_{1,ss+pp}^q(z, m_{\pi\pi})}$$

twist-3 DiFF

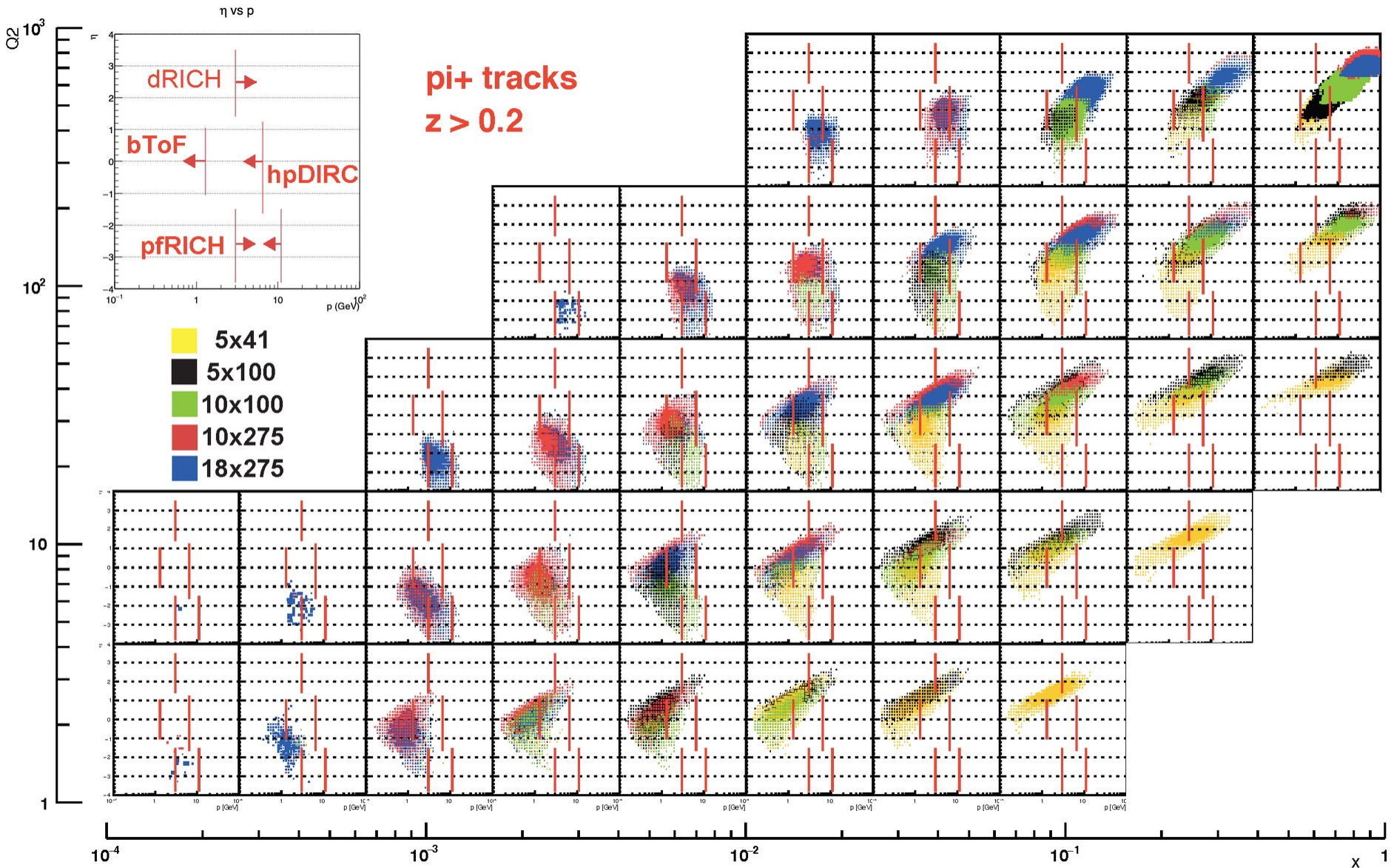


Courtoy, Aurore, et al. *Phys.Rev.D* 106 (2022)
 Courtoy, Aurore – [CPHI 2022](#)

- **Scenario I: Wandzura-Wilczek (WW) Approximation**
 - Drop twist-3 DiFF
- **Scenario II: Beyond WW approximation**
 - Estimate max integrated twist-3 DiFF from COMPASS A_{UL} and A_{LL}



Wide Coverage



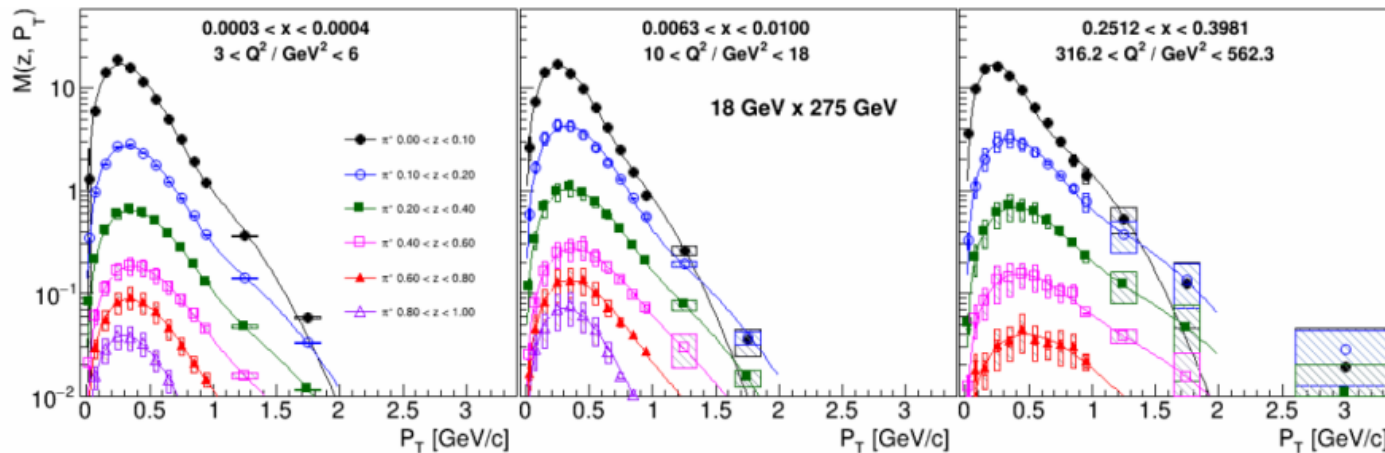
Depolarization Factors

Twist 2

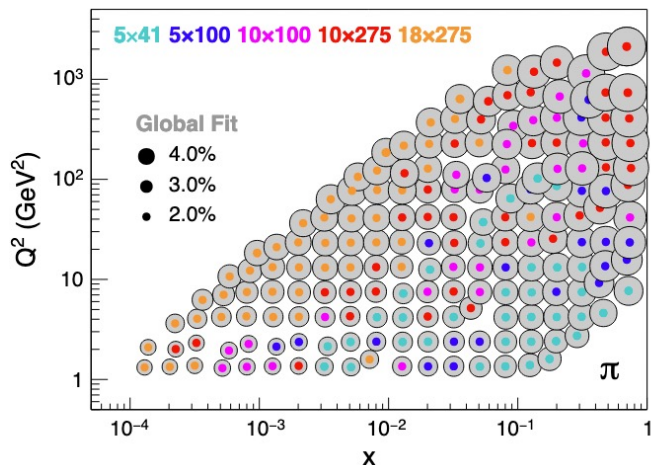
	Polarization	Depolarization
Boer-Mulders	UU	B
Sivers	UT	1
Transversity	UT	B/A
Kotzinian-Mulders	UL	B/A
Wormgear (LT)	LT	C/A
Helicity DiFF G_1^\perp	LU	C/A
	UL	1
<u>Twist 3</u>		
$e(x)$	LU	W/A
$h_L(x)$	UL	V/A
$g_T(x)$	LT	W/A

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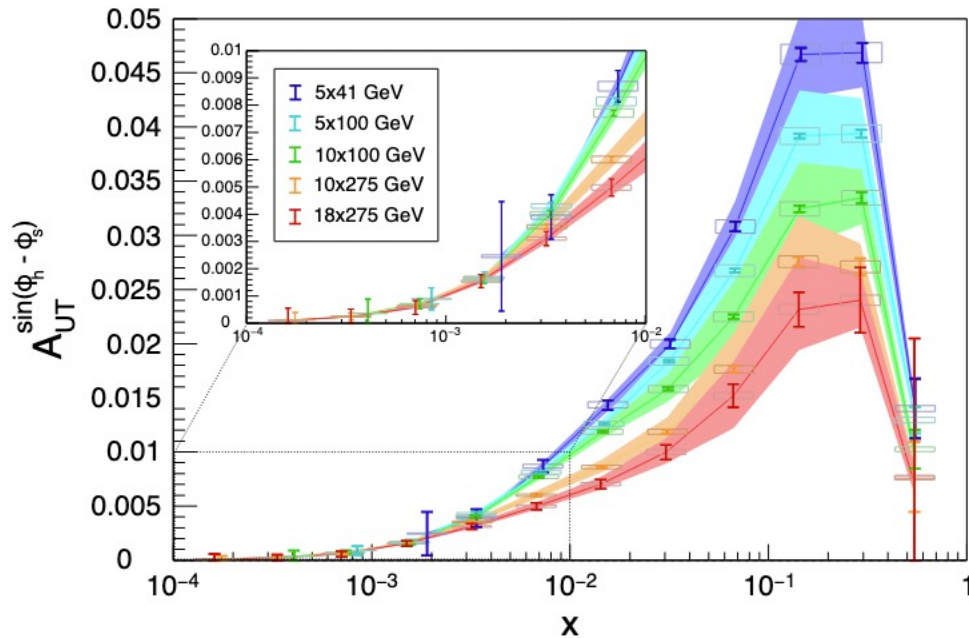


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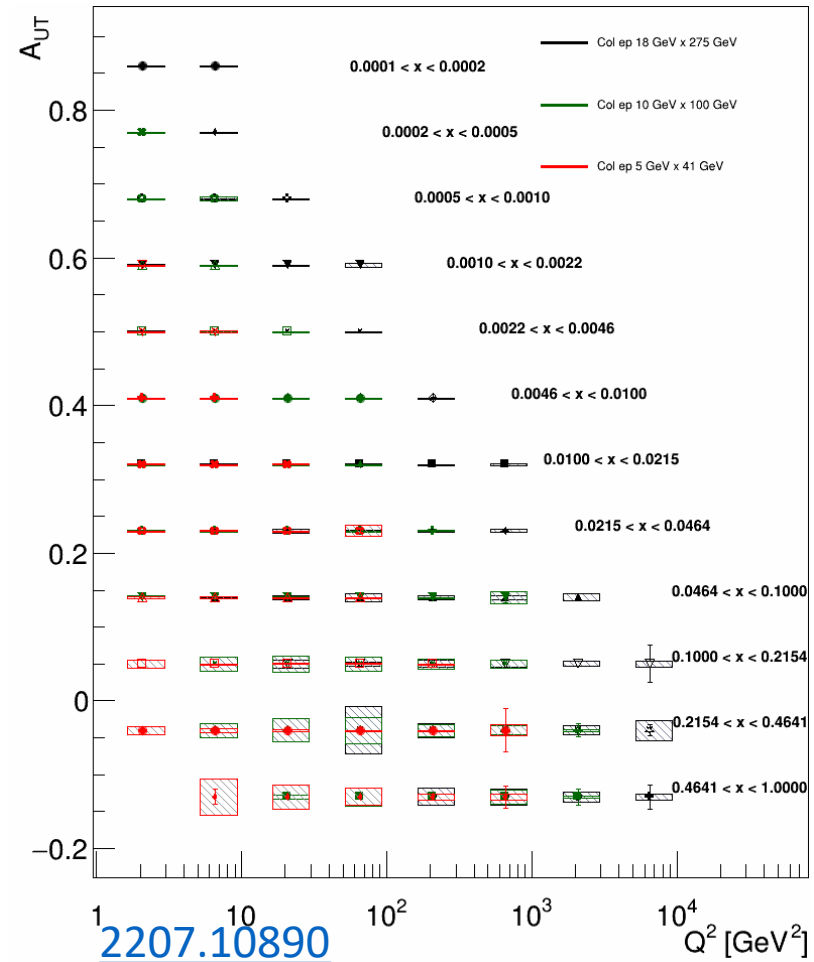
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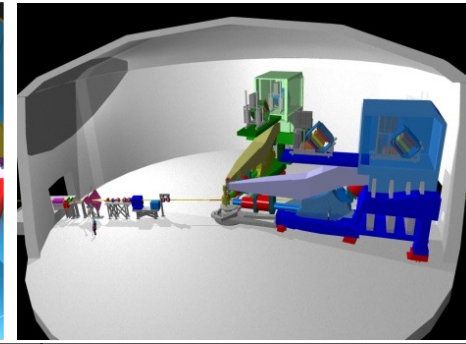
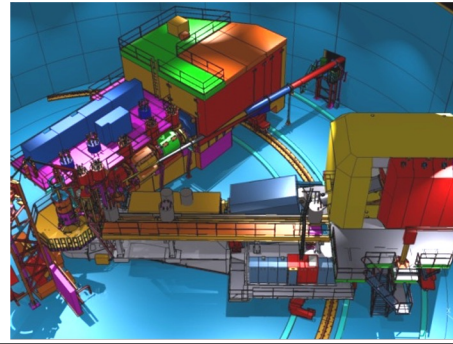
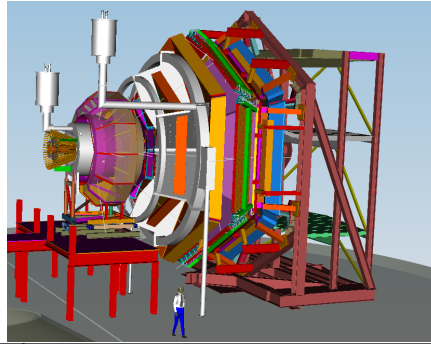
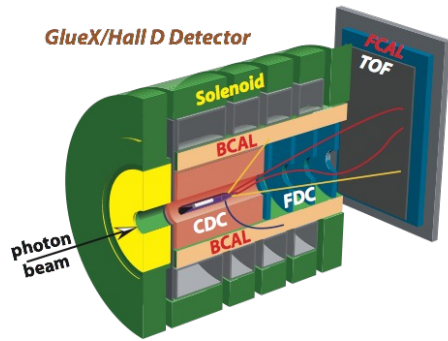
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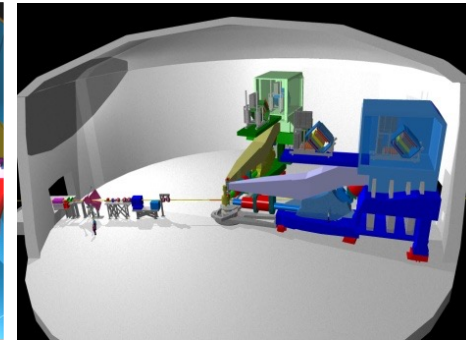
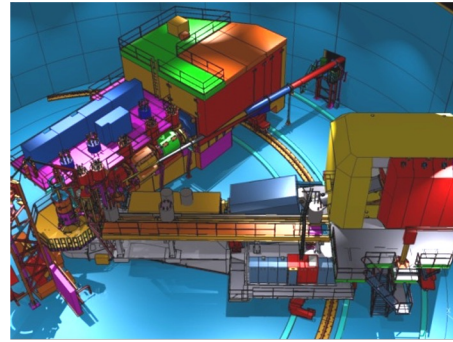
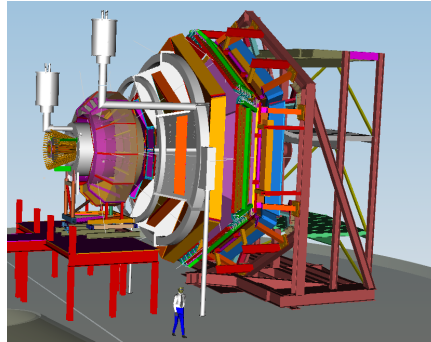
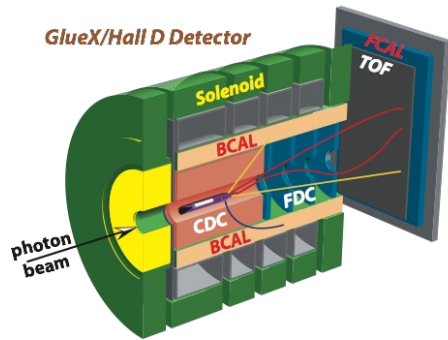
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Detector Requirements: Complementarity



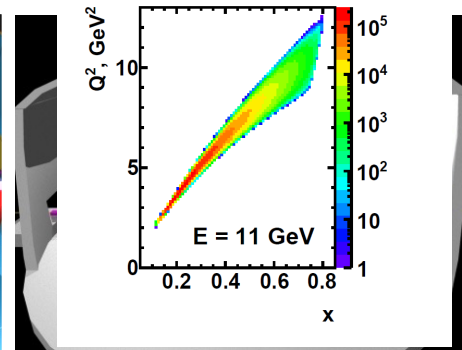
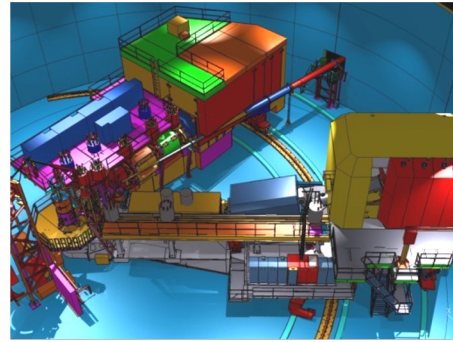
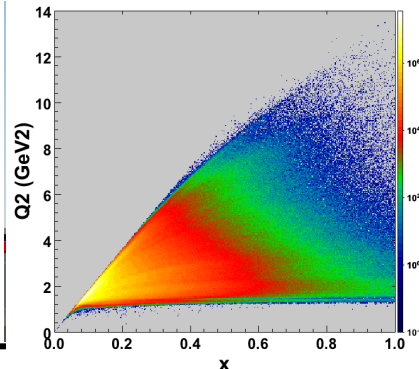
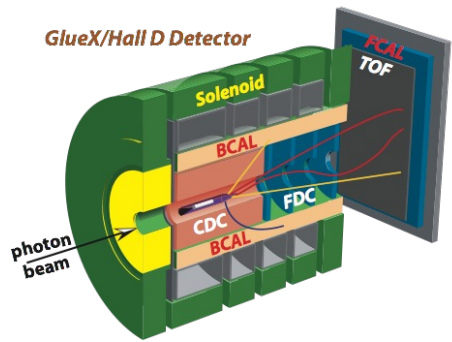
Hall D	Hall B	Hall C	Hall A
excellent hermeticity	luminosity 10^{35}	energy reach	custom installations
polarized photons	hermeticity	precision	
$E_\gamma \sim 8.5-9$ GeV	11 GeV beamline		
10^8 photons/s	target flexibility		
good momentum/angle resolution	excellent momentum resolution		
high multiplicity reconstruction	luminosity up to 10^{38}		
particle ID			

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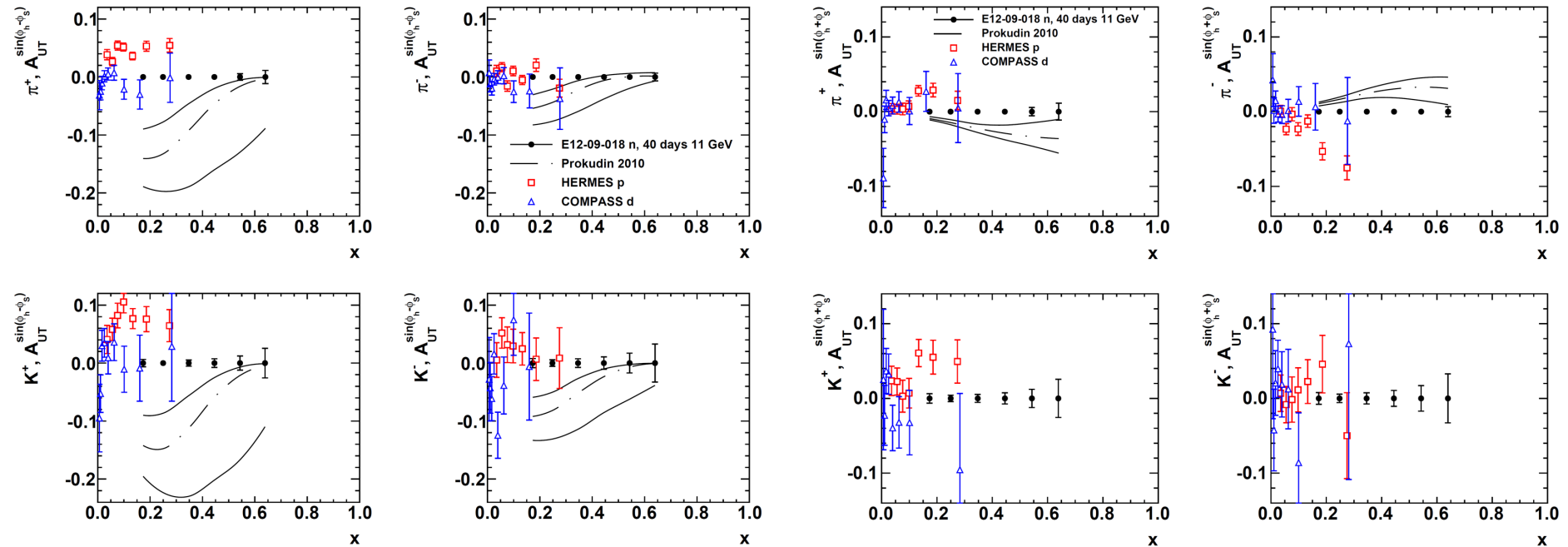
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SBS+BB Projected Results: Collins and Sivers SSAs



Projected A_{UT}^{Sivers} vs. x (11 GeV data only)

Projected A_{UT}^{Collins} vs. x (11 GeV data only)

- E12-09-018 will achieve statistical FOM for the neutron $\sim 100X$ better than HERMES proton data and $\sim 1000X$ better than Hall A E06-010 neutron data. *Near-future more precise COMPASS deuteron data will sharpen expected impacts, urgency of E12-09-018*
- SBS installation starts 2020. E12-09-018 could run as early as 2022; 2023 more likely.